

3/31/81

ME235B Advanced Finite Elements.

Conduct of Course:

HW's (2%); Computing; Etc $\leftarrow$ TA Reader

Midterm 1 hr in class

Final 3 hr in class

IBM 3033 (C.I.T.) can access evening + when
1 1/2 hr constraint per session can relog on the
6 hr/wk

Content: 1. Finish Linear Statics
2. Linear Dynamics

Overview of Final

expect δ fns will reappear again.
on
midterm Barlow stress point problem "

σ^0, ε^0 (last problem) become part of source terms.

Ch 3 §9: Calculation of derivatives of shape fns (since they appear in our bilinear forms).

i.e. $N_{a,x}$: we normally derived $N_a(\xi)$ where ξ are the natural coordinates

it is also impossible to get the inverse map of $\xi(x)$ from $x(\xi)$ especially in
multidim.

let $n_{nd}=2$ for our purpose

$$N_{a,x} = N_{a,\xi} \xi_{,x} + N_{a,\eta} \eta_{,x}$$

$$N_{a,y} = N_{a,\xi} \xi_{,y} + N_{a,\eta} \eta_{,y}$$

$$\begin{bmatrix} N_{a,x} & N_{a,y} \end{bmatrix} = \begin{bmatrix} N_{a,\xi} & N_{a,\eta} \end{bmatrix} \begin{bmatrix} \xi_{,x} & \xi_{,y} \\ \eta_{,x} & \eta_{,y} \end{bmatrix} = N_{a,\xi} \underbrace{\begin{bmatrix} \xi_{,x} \\ \xi_{,y} \end{bmatrix}}_{\text{jacobian of inverse map}}$$

we normally have $\underline{x}(\underline{\xi}) = \sum_{a=1}^{n_{en}} N_a(\underline{\xi}) \underline{x}_a^e$ $\underline{x} = (x, y)$ $\underline{\xi} = (\xi, \eta)$

we can compute $\underline{x}_{,\xi} = \begin{bmatrix} x_{,\xi} & x_{,\eta} \\ y_{,\xi} & y_{,\eta} \end{bmatrix}$ since we have $\underline{x}(\underline{\xi})$.

$$\text{i.e. } x_{,\xi} = \sum_{a=1}^4 N_{a,\xi} x_a^e \quad x_{,\eta} = \sum_{a=1}^4 N_{a,\eta} x_a^e \quad y_{,\xi} = \sum N_{a,\xi} y_a^e \quad y_{,\eta} = \sum N_{a,\eta} y_a^e$$

$$\text{now note } \xi_{,x} = [x_{,\xi}]^{-1} = \begin{bmatrix} y_{,\eta} & -x_{,\eta} \\ -y_{,\xi} & x_{,\xi} \end{bmatrix} \cdot \frac{1}{j}$$

$$\text{where } j = \det x_{,\xi} = x_{,\xi} y_{,\eta} - x_{,\eta} y_{,\xi}$$

These are calculations performed in shape fn subroutine. We will now look at such a subroutine. The one to be handed out will be more involved.

"Shape fn subroutine"

its output will (or can) be $j, N_a, N_{a,x}, N_{a,y}$ evaluated at some point $(\bar{\xi}, \bar{\eta})$ for $1 \leq a \leq n_{en}$ likely to be at quadrature point.

The input would be the point $(\bar{\xi}, \bar{\eta})$ also x_a^e, y_a^e $1 \leq a \leq n_{en}$

in flowchart form:

1. Given: input
2. Calculate: evaluate $N_a(\bar{\xi}, \bar{\eta})$; $N_{a,\xi}(\bar{\xi}, \bar{\eta})$; $N_{a,\eta}(\bar{\xi}, \bar{\eta})$ $1 \leq a \leq n_{en}$ Na's are already in code
3. $x_{,\xi}(\bar{\xi}, \bar{\eta}) = \sum N_{a,\xi}(\bar{\xi}, \bar{\eta}) x_a^e$ also for $x_{,\eta}$; $y_{,\xi}$; $y_{,\eta}$ Na's are " in code.
4. Compute $j(\bar{\xi}, \bar{\eta}) = x_{,\xi} y_{,\eta} - x_{,\eta} y_{,\xi}$ using results of (3).
5. Now we can calculate $\langle N_{a,x}, N_{a,y} \rangle$ using i.e.

$$\left. \begin{aligned} N_{a,x} &= (N_{a,\xi} y_{,\eta} - N_{a,\eta} y_{,\xi}) / j \\ N_{a,y} &= -(N_{a,\xi} x_{,\eta} - N_{a,\eta} x_{,\xi}) / j \end{aligned} \right\} 1 \leq a \leq n_{en}$$
6. return to next step

Exercise do the same for $N_{a,x}$ for $\underline{x} = x, y, z$

$$\text{i.e. } N_{a,\underline{x}} = N_{a,\underline{\xi}} \xi_{,\underline{x}} \quad \text{where } \xi_{,\underline{x}} \text{ is } 3 \times D \text{ of } j \quad \text{where } \xi_{,\underline{x}} = [x_{,\xi}]^{-1}$$

$$\text{where } [x_{,\underline{\xi}}]^{-1} = \frac{[\text{cofactor } x_{,\underline{\xi}}]}{j}$$

$$\text{i.e. cofactor} = y_{,\eta} z_{,\xi} - y_{,\xi} z_{,\eta} \quad \left| \begin{array}{cc} & \\ & \end{array} \right| \text{ cofactor of 11}$$

$$w_{ij} = x_{i5} \text{ cof}_{i1} + x_{i7} \text{ cof}_{i2} + x_{i9} \text{ cof}_{i3}$$

and we follow this through as before.

§10. Programming Aspects in FORTRAN

"Shape fn. subroutine" - these names may appear world wide. we will do this for 4 node bilin quad

Our notation	Fortran	Array Dim	Our notation	Fortran names	Array Dim				
n_{sd}	$NSD(=2)$		N_a	$SH(3, I)$					
n_{en}	$NEN(=4)$		$N_{a,x}$	$SH(1, I)$	$(NSD+1, NEN)$				
ξ_i	S		$N_{a,y}$	$SH(2, I)$					
η_i	T		$\xi_{a/2}$	$SA(I)$	NEN writes in data stream				
x_a^e, y_a^e	$XE(1, I) \ XE(2, I)$	(NSD, NEN)	$\eta_{a/2}$	$TA(I)$	NEN				
j	XE or "XL" <table><tr><td></td><td></td><td></td><td></td></tr></table>						$\begin{bmatrix} x_{i5} & x_{i7} \\ y_{i5} & y_{i7} \end{bmatrix}$	$XS(I, J)$	(NSD, NSD)
	DET								

Coding

sub 7 input output
SUBROUTINE SHAPE (S, T, XE, DET, SH)

ROUTINE TO CALCULATE

DIMENSION SA(4), TA(4), XS(2,2), XE(2,4), SH(3,4)

DATA SA/-0.5, +0.5, .5, -0.5, TA/-0.5, -0.5, .5, .5/

DATA ZERO, P5/0.0, 0.5/

SHAPE FNS & DERIV IN LOCAL COORDS.

DO 10 I=1,4

temporary
usage of SH

$$\begin{aligned} SH(1, I) &= P5 + TA(I) * T = \frac{1}{2} (1 + \eta_a \eta) \\ SH(2, I) &= P5 + SA(I) * S = \frac{1}{2} (1 + \xi_a \xi) \\ SH(3, I) &= SH(1, I) * SH(2, I) = N_{a, \xi} = \xi_{a/2} \cdot \frac{1}{2} (1 + \eta_a \eta) \\ SH(1, I) &= SA(I) * SH(3, I) = N_{a, \eta} = \eta_{a/2} \cdot \frac{1}{2} (1 + \xi_a \xi) \end{aligned}$$

10 SH(2, I) = TA(I) * SH(2, I) = $N_{a, \eta} = \eta_{a/2} \cdot \frac{1}{2} (1 + \xi_a \xi)$

Jacobian matrix $(X_{i,j})$

DO 20 I=1,2

DO 20 J=1,2

XS(I, J) = ZERO

20 $XS(I,J) = XS(I,J) + XE(I,K) * SH(J,K)$

C...

C... JACOBIAN DETERMINANT

C...

DET = XS(1,1) * XS(2,2) - XS(1,2) * XS(2,1)

C...

C... SHAPE FN DERIVS IN GLOBAL COORDS

C...

DO 30 I=1,4

TEMP = (XS(2,2) * SHAPE(1,I) - XS(2,1) * SHAPE(2,I)) / DET This is Na_x

SH(2,I) = (-XS(1,2) * SH(1,I) + X(1,1) * SH(2,I)) / DET This is Na_y

SH(1,I) = TEMP

RETURN

END

$Na_x = Na_x \xi_x + Na_y \eta_x$

Continuation

4/2/81

Now looking at the element stiffness formation for 2D, heat conduction element.

$$K_{ab}^e \quad (a,b \text{ are local nodal indices} = \text{local eqn nos.}) = \int_{\Omega^e} \underline{B}_a^T \underline{D} \underline{B}_b d\Omega$$

Since 1 DOF/node.

using numerical integration we had defined

$$\int_{\Omega^e} d\Omega = \int_{-1}^1 \int_{-1}^1 f(\xi, \eta) d\xi d\eta = \int_{-1}^1 \int_{-1}^1 g(\xi, \eta) d\xi d\eta$$

for $n_{sd} = 2$ $\underline{B}_a = \langle Na_x \quad Na_y \rangle^T$ and $\underline{D} = \underline{K} = \begin{bmatrix} K_{11} & K_{12} \\ K_{21} & K_{22} \end{bmatrix}$ w/ $K_{ij} = K_{ji}$

thus $\int_{-1}^1 \int_{-1}^1 g(\xi, \eta) d\xi d\eta \approx \sum_{l=1}^{n_{int}} (\underline{B}_a^T \underline{D} \underline{B}_b) \big|_{(\xi_l, \eta_l)} W_l$

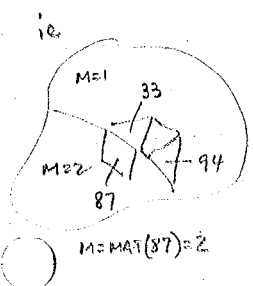
Now we will define $\tilde{\underline{D}}_l = \underline{D} \cdot \underline{J} W_l$ $K_{ab}^e \approx \sum_{l=1}^{n_{int}} (\underline{B}_a^T \tilde{\underline{D}}_l \underline{B}_b)_l$ by writing $\tilde{\underline{D}}$ over $\underline{D}$

we will define our fortran variable names

NOTATION	FORTRAN NAME	ARRAY DIM	NOTATION	FORTRAN NAME	ARRAY DIM
a, b	I, J		η_{en}	NEN	
e	N		material set nos.	M = MAT(N)	NUMEL
			where M = no. of material sets		
n_{el}	NUMEL		no. of material sets	NUMAT	
n_{sd}	NSD (=2)		material property array	C(I, M)	(3, NUMAT)
n_{int}	NINT		ie $K_{11}, K_{12}, K_{21}, K_{22}$		

Notation	Fortran Name	Array Dim	Notation	Fortran Name
k_{ab}	$S(I, J)$	(NEN, NEN)	D_{11}, D_{12}, D_{22}	$D11, D12, D22$
$\tilde{B}_1, \tilde{B}_2$	$SG(L), TG(L)$	$NINT$	$(\tilde{D}B_1)_1, (\tilde{D}B_2)_2$	$DB1, DB2$
$\tilde{L}$	$W(L)$	$NINT$	x_a^e, y_a^e	$XE(1, I) \quad XE(2, I) \quad (NSD, NEN)$
w_L	$B1 = SH(1, I)$	$(NSD+1, NEN)$		
$N_{a,x}$	$B2 = SH(2, I)$			
$N_{a,y}$	$(SH(3, I))$			
N_a				

The following is general as far as coding style. We will not give full code, but we will assume all arrays are dimensioned appropriately, e.g. $NUMAT, NUMEL, NSD, NINT, MAT(NUMEL), C(B, MAT), SG(NINT), TG(NINT), W(NINT)$ are all given. ie gauss quadrature is already picked.



```

C ... ELEMENT LOOP
C
DO 100 N=1, NUMEL
C
C ... OBTAIN PROP. SET. NO.
C
C
M=MAT(N)

C ... INITIALIZE EL STIFF.
C ... CALL CLEAR(S, NEN*NEN)  this sets S(I,J)=0.
C

C ... OBTAIN EL COORDS OF LOCALIZED COORDS
C ... CALL LCOORD( )

C ... BEGIN INTEG LOOP
C ... DO 10 L=1, NINT
C ... OBTAIN SHAPE FNS, DERIV, DET
C ... CALL SHAPE(SG(L), TG(L), XE, DET, SH)

C ... COMPUTE R MATRIX (really  $\tilde{D}$  matrix)
C
DET = DET * W(L)  (if we)
D11 = C(1, M) * DET
D12 = C(2, M) * DET
D22 = C(3, M) * DET
} this assumes that material prop/element is constant.

C ... compute  $\tilde{D} * B$ 
C
DO 10 J=1, NEN  column loop
B1 = SH(1, J) = N_{a,x}
B2 = SH(2, J) = N_{a,y}
DB1 = D11 * B1 + D12 * B2
DB2 = D12 * B1 + D22 * B2
}
C
C ... compute  $B^T D B$ 
C
DO 10 I=1, J  Since we have symmetry if we had no sym then replace J by I

```

$$B1 = SH(1, I)$$

$$B2 = SH(2, I)$$

$$10 \quad S(I, J) = S(I, J) + B1 * DB1 + B2 * DB2$$

... Compute lower triangular part.

CALL SYMM (S, NEN)

... ASSEMBLE S INTO GLOBAL STIFF

100 CALL ADDSTF () pass in part of IM array dealing w/ this particular elem.
this matrix $K = A k_{ab}^e$ (from last quarter).

for elasticity, we ^{would} calculate all elements of the block k_{ab}^e here.

Now we will do same for elasticity in 2D less fortran.

remember that k_a^e was formed in terms of nodal submatrices k_{ab}^e

Element Stiffness Formulation for a 2-D elasticity element.

4/7/81

k_a^e can be written in terms of $n_{en} \times n_{en}$ nodal submatrices k_{ab}^e
ie eqns (6)-(9) in § 2.8

$$k_{ab}^e = \int_{\Omega_e} B_a^T \tilde{D} B_b d\Omega \quad (7)$$

$n_{ed} \times n_{ed}$

now we use the same format as in heat conduction by using gauss quadr

$$k_{ab}^e \approx \sum_{l=1}^{n_{int}} (B_a^T \tilde{D} B_b)_l \quad \text{where } \tilde{D} = \tilde{D} f W_l$$

in 2-D

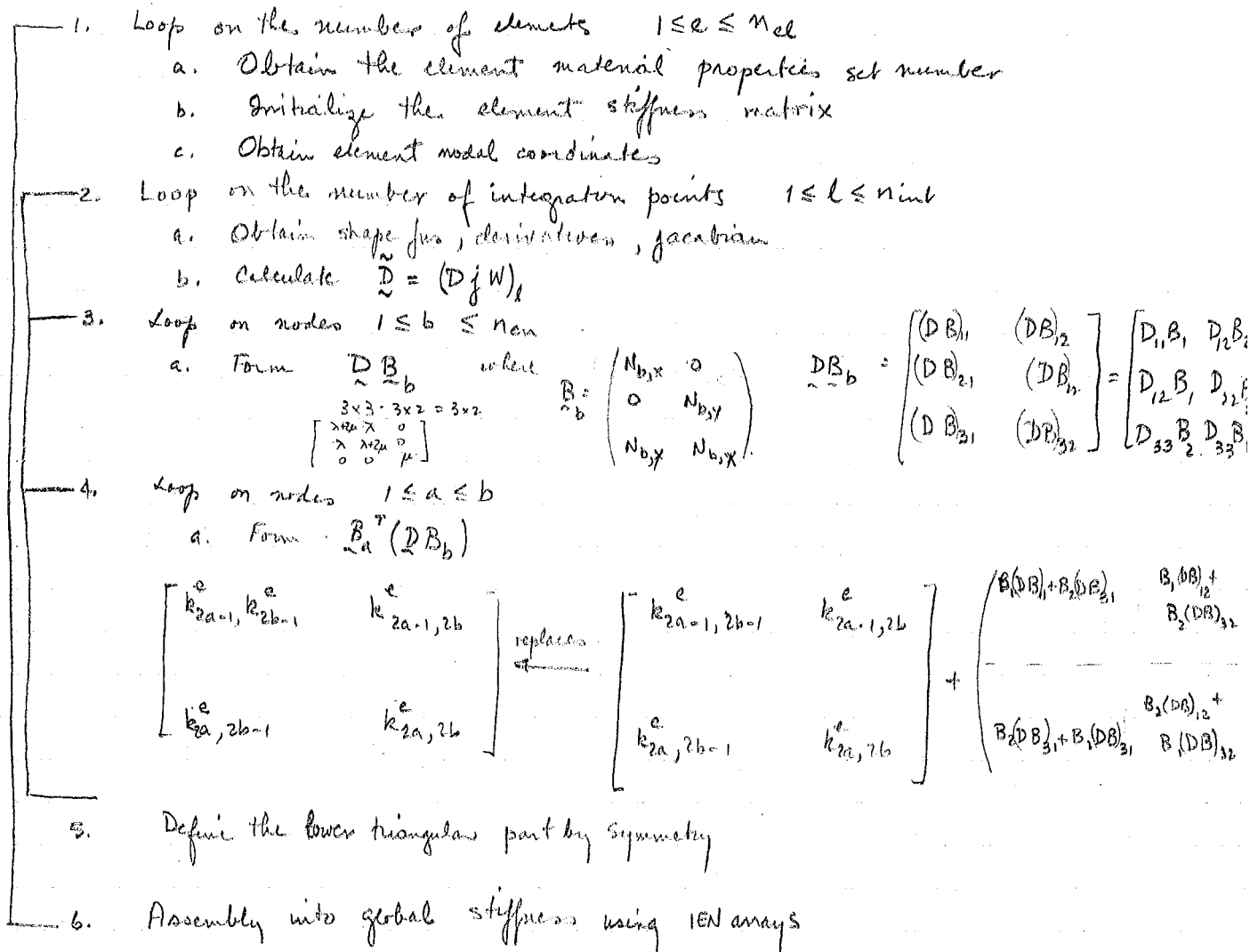
$$k_{ab}^e \text{ is } 2 \times 2 = \begin{bmatrix} k_{2a-1, 2b-1}^e & k_{2a-1, 2b}^e \\ k_{2a, 2b-1}^e & k_{2a, 2b}^e \end{bmatrix}$$

where these elements are referred to the entries of k_a^e

because of the symmetry of k_a^e we need only the nodal submatrices on or above the diagonal.

Before the flowcharting we will assume that $D_{13} = D_{23} = 0$
ie the material properties are constant over the material

Flow chart



Remarks

- forces vectors are formulated outside this sequence except gravity forces and.
 - except for: the part of the element force due to prescribed displacement BC
- $$f_p^e = \int_{\Omega_e} N_i f_i d\Omega + \int_{\Gamma_e} N_i h_i d\Gamma - \sum_{q=1}^{n_{eq}} k_{pq}^e g_q^e$$

References: "Computer Procedures for finite Element Analysis" Chapter 24 in
 O. C. Zienkiewicz The finite Element Method McGraw-Hill, London 1977
 by R. L. Taylor

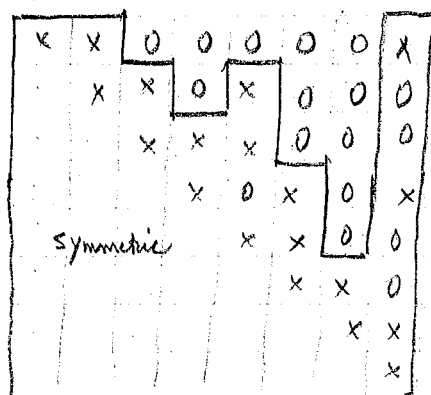
Hinton & Owen - Finite Element Programming, Academic Press, London 1971

LEARN program allows either use of existing or adding changes

Definition The data structure is the mode of storage, location and history of data employed.

- Compact column structure - used to keep only non zero element.
"skyline", "profile", "active column"
- Dynamic Storage Allocation - arrays are set during problem execution
- Element Groups (for each different material set no, inputted simultaneously).

example of $[K^e]$ which is 8×8 let x be nonzero K^e will be stored in 1-D array called A .



- profile or skyline

in BERN - location of diag is stored in IDIAG
dimension of IDIAG is NEQ .

Dimension of A is $NA = \sum$ of column heights
 $1+2+2+2+4+3+2+8$

thus $A(1) = K11$ $A(2) = K12$ $A(3) = K22$ $A(4) = K32$ $A(5) = K33$... $A(24) = K88$

the heights actually depend on D and since we know D we know the structure of K a priori

4/9/81

ARRAY IDIAG gives last element location of column I i.e. $IDIAG(I) = L$; for the above

$IDIAG(1) = 1$ $IDIAG(2) = 3$ $IDIAG(3) = 5$ $IDIAG(4) = 7$ $IDIAG(5) = 11$

$IDIAG(6) = 14$ $IDIAG(7) = 16$ $IDIAG(8) = 24$

this allows us to compact the LHS of the equation $K d = f$

the rhs of the equation is stored in B $f_i = B(I)$

let K be symmetric. Then $\exists$ a nonsingular upper triangular matrix U with unit diagonal entries and a diagonal matrix D $\exists$.

$$K = U^T D U$$

The number of negative, zero & positive eigenvalues of K is equal to the no. of negative, zero, positive diagonal entries of D respectively.

Storage in blank common

Except for element group data, which is stored on a disk file and read in and out of core as needed, all of the main arrays are stored in blank common. This includes both integer and floating point data. In the main program and subroutine LEARN, blank common is denoted by the one-dimensional array A.

The locations and lengths of arrays in blank common are as follows:

Storage in Blank Common - Solution Phase

First Word

Array

example
N1 = 1

ID(NDOF, NUMNP)

N2 = N1 + NDOF*NUMNP

X(NSD, NUMNP)

*takes into account
prescribed body
condition*

N3 = N2 + NSD*NUMNP*IPREC

F(NDOF, NUMNP)
(=D(NDOF, NUMNP))

N4 = N3 + NDOF*NUMNP*IPREC

IDLAG(NEQ)

N5* = N4 + NEQ

B(NEQ)

N6 = N5 + NEQ*IPREC

A(NA)

NA is sum of column heights

N7 = N6 + NA*IPREC

E(MAX)

N8 = N7 + MAX-1 (* total required storage; must be .LT. MTOT
for execution.)

* During input phase, N1 to N5 are the same as above and element group data begins in N5. Individual element routines determine necessary storage for element groups.

where

- $ID(M,N)$ = equation number for degree of freedom M ,
 node N
- $X(M,N)$ = M th coordinate of node N
- $F(M,N)$ } M th direction force/displacement component
 $D(M,N)$ } of node N
- $1 \leq N \leq NUMNP$
 $1 \leq M \leq NDOF$
- $IDIAG(N)$ = location of N th diagonal element in coefficient
 array A .
- $E(N)$ = N th force/displacement component
- $A(N)$ = N th element of coefficient matrix stored in
 compacted column form
- $E(N)$ = N th (single precision) word of element
 group array
- $NUMNP$ = number of nodes
- NSD = number of space dimensions
- $NDOF$ = number of degrees of freedom
- NBE = number of equations
- NA = number of terms in A
- MAX = number of single-precision words in longest
 element group data file
- $IPREC$ = precision needed (1/2 for single/double precision math).

This array name is used for the coefficient matrix in various subroutines and should not be confused with the array A denoting blank common.



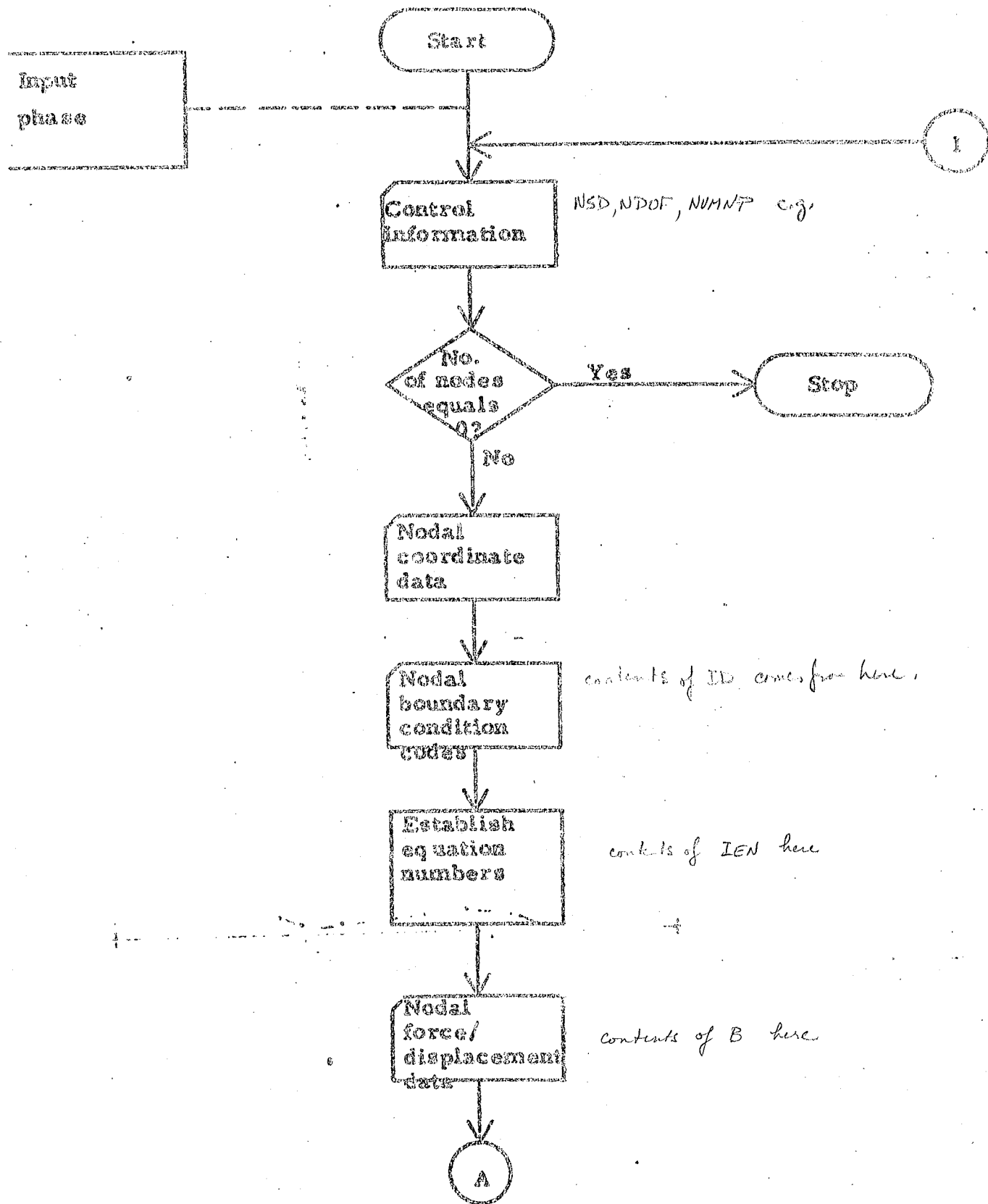
Storage Requirements of Plane Stress/Strain Element (NTYPE=3)

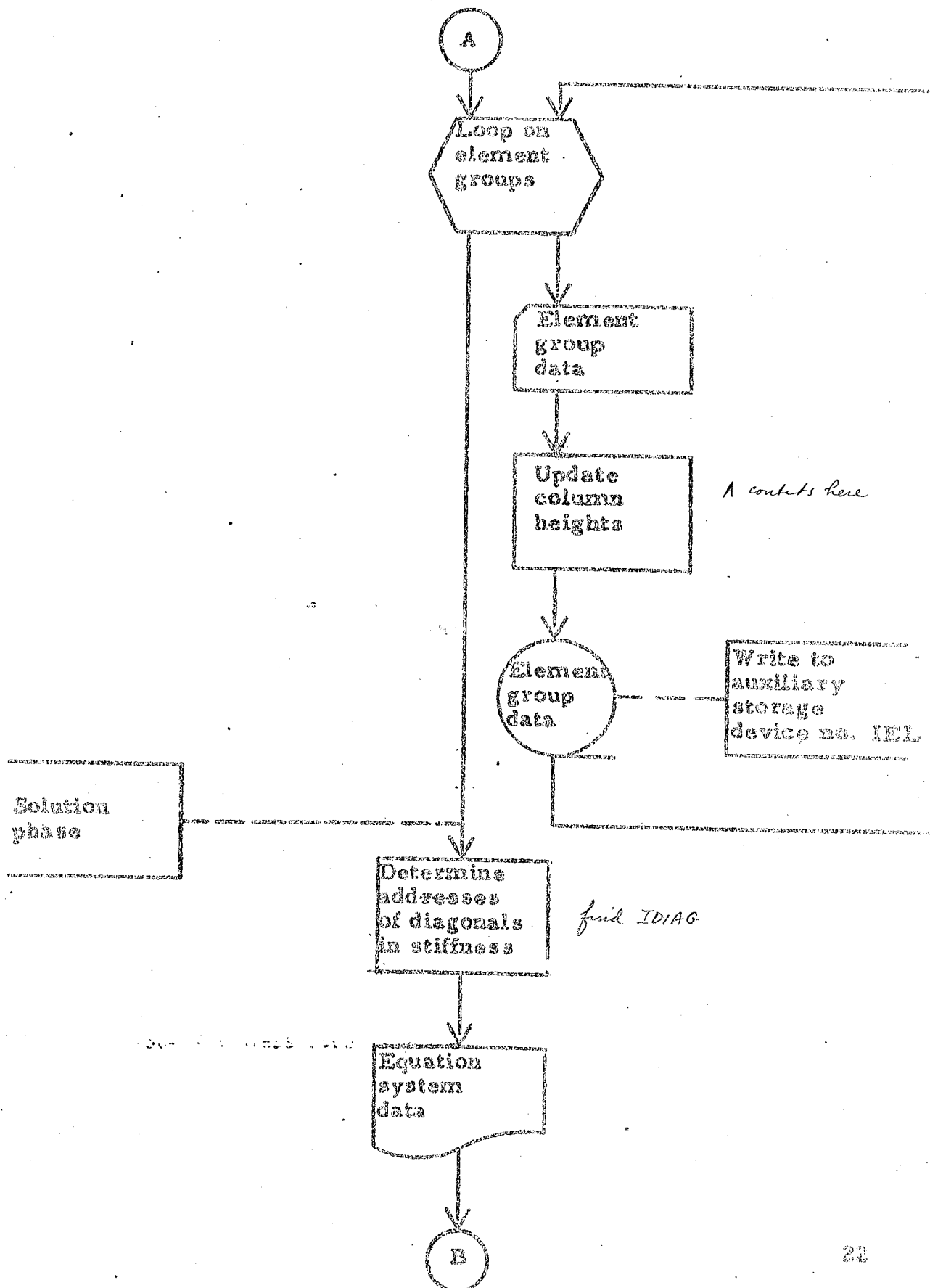
<u>First Word</u>	<u>Array</u>
N101 = NF*	E(NUMAT) <i>Young's mod</i>
N102 = N101 + NUMAT*IPREC	POIS(NUMAT) <i>ν</i>
N103 = N102 + NUMAT*IPREC	WT(NUMAT) <i>weight density for gravity calculations</i>
N104 = N103 + NUMAT*IPREC	TH(NUMAT) <i>thickness</i>
N105 = N104 + NUMAT*IPREC	GRAV(2)
N106 = N105 + 3*IPREC	MAT(NUMEL) <i>material set no.</i>
N107 = N106 + NUMEL	IEN(4, NUMEL) <i>4 node quad.</i>
N108 = N107 + 4*NUMEL	LM(2, 4, NUMEL) <i>i a NSD = 2</i>
N109 = N108 + 8*NUMEL	IELNO(NUMPR)
N110 = N109 + NUMPR	ISIDE(NUMPR)
N111 = N110 + NUMPR	PRES(NUMPR) <i>pressure on side</i>
NL = N111 + NUMPR*IPREC	
LENGTH = NL - NF (=total required storage for element group)	

The meaning of the arrays here may be deduced from the description in the section entitled "User's Manual."

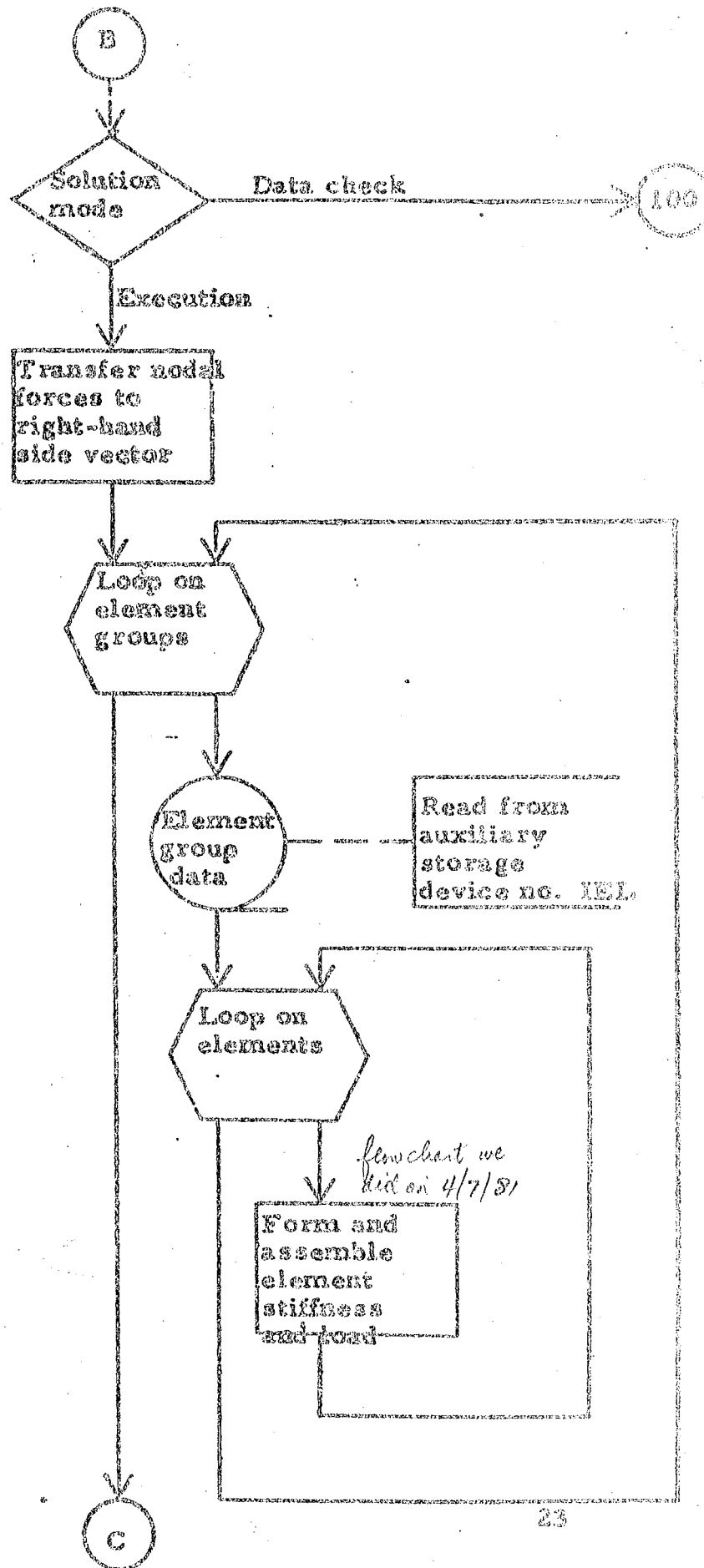
$$MAX = \max_i (NL - NF)_i$$

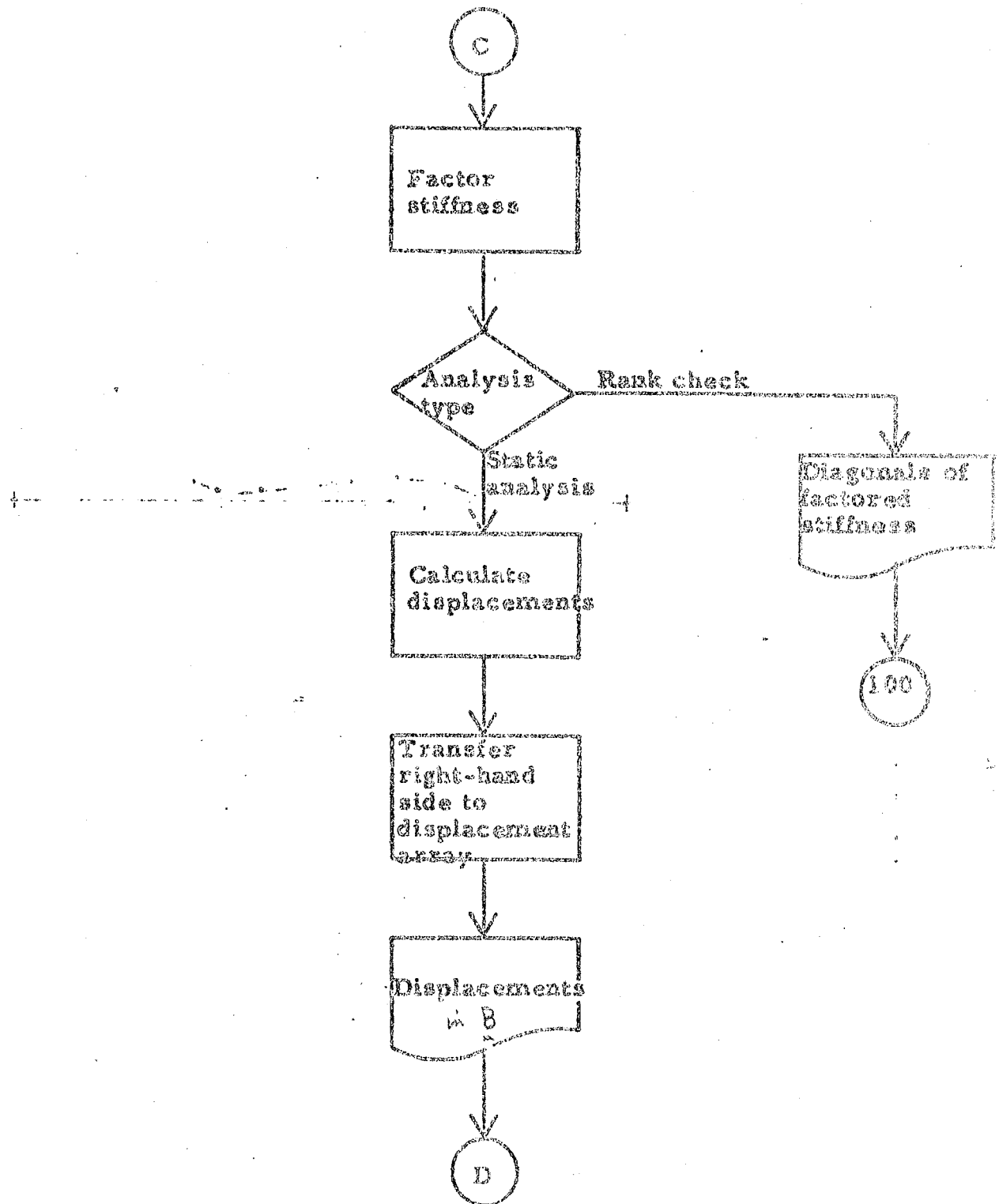
*NF = NS during input phase and NF = N7 during solution phase.



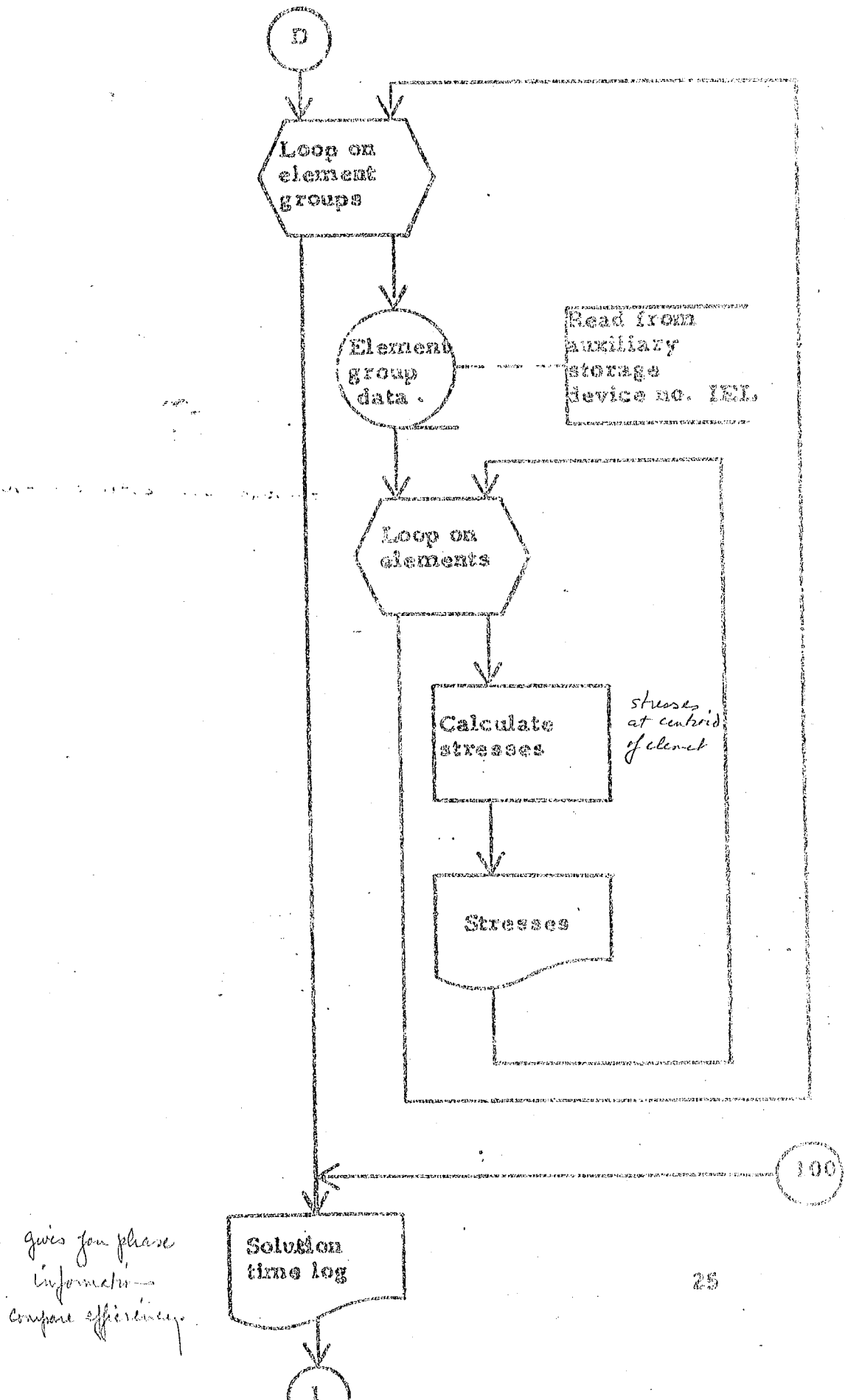












if $\underline{K}$ is + def. then eigenvalues of $\underline{D}$ are > 0
 $\underline{U}$ has same profile of $\underline{K} \Rightarrow$ storage of $\underline{U}$: can use same as $\underline{K}$
 storage of $\underline{D}$: need to provide this

Factoriz is done via Crout's method.
 define intermediate steps

$$\underline{U}^T \underline{x} = \underline{F} \quad (\text{forward reduction}) \quad \text{solve for } \underline{x}$$

$$\underline{D} \underline{y} = \underline{x} \quad \text{solve for } \underline{y}$$

$$\underline{U} \underline{d} = \underline{y} \quad (\text{back substitution})$$

see pg 10-11 of handout - talked above blank common.

suppose we want to pass $\underline{x}$ (NSD, NUBP) to COORD

use CALL COORD (A(NZ), NSD, NUBP) where $\underline{x}$ is defined in blank common array A(I) and $\underline{x}$ begins at A element NZ.

I an element library for each element we pick
 ie SUBROUTINE ELMLIB

```

      GO TO (1,2,...), I must be passed
or be in common.
1    CALL TRUSS
      RETURN
2    CALL PLANE
      RETURN
3

```

To input data

1st CARD TITLE card 80 characters like comment
 control card (I5 for integer F10. for float) This card uses I
 all these are found in LEARN

data card

NODE #	X for DOF=1	Y for DOF=2
21	1	
31		1

data generator ie NG=1 says generate data from 21 to 31 by 1
 using linear interpol between $X_{DOF,1}(21)$ & $X_{DOF,1}(31) = \dots$

4/14/81

LEARN books will be in THURS

Monday 4/20/81 evening will be a short class on computing on classwork at CIT.

Chapter 4. - Advanced Topics in F.E.

we will be dealing in problems w/ near or complete incompressibility, some structural problems.

§ 0. Prelude - why do things work? (sometimes).

Look at heat conduction (or elasticity)

Recall weak form

$$(w) \quad a(w, u) = (w, f) + (w, h)_{\Gamma} \quad \forall w \in U \quad (1)$$

$$(G) \quad a(w^h, u^h) = (w^h, f) + (w^h, h)_{\Gamma} \quad \forall w^h \in U^h \quad (2)$$

we've done this "by the rules"

Now we will look at errors involved.

take (1) and let w^h replace w since $w^h \in U$
then

$$a(w^h, u) = (w^h, f) + (w^h, h)_{\Gamma} \quad (3)$$

subtract (2) from (3)

$$a(w^h, u^h - u) = 0 \quad \text{then } u^h - u = e \Rightarrow e \perp w^h \quad \forall w^h \in U^h$$

$$\text{Consider } a(e + w^h, e + w^h) = a(e, e) + 2a(w^h, e) + a(w^h, w^h) \quad \forall w^h \in U^h$$

Recall that $a(w, w) > 0$ if $w \neq 0$ & $\forall w \in U$ (with right b.c.).

$$\Rightarrow a(e, e) \leq a(e + w^h, e + w^h) \quad \forall w^h \in U^h$$

$$a(e + w^h, \dots) = a(u^h - u + w^h, \dots) = a(U^h - u, U^h - u) \quad \text{where } u \in A$$

where $A = U + \text{b.c. w/ } U^h = u^h + w^h$;

since $u^h \in A^h = U^h + \text{b.c.} \Rightarrow U^h \in A^h$ since $w^h = 0$ as b.c.

now $a(e, e) \leq a(U^h - u, U^h - u) \quad \forall U^h \in A^h \Rightarrow U^h$ can be any fn.

hence any fn picked in A^h will have a measure $\geq$ the error

hence $a(e, e)$ is the best approximation. All problems looked at previously satisfy this.

$$a(e, e) = \int_{\Omega} e_{,i} \kappa_{ij} e_{,j} d\Omega. \quad \text{this is the measure used in heat conduction}$$

$$a(e, e) = \int_{\Omega} e_{(i,j)} c_{ijkl} e_{(k,l)} d\Omega \quad \text{" " " " " " elasticity}$$

hence we are basically "least squares" best fitting the fn. ie we are doing a minimization of the energy. However this result may be misleading

what can possibly go wrong? Suppose $\mathcal{A}^h$ could be a bad approximation.

1. this can happen in incompressible or nearly incompressible problems
2. also the measure $a(\cdot, \cdot)$ might be inappropriate ie best approx wrt wrong measure. (ie incomp problems)
3. Combination of 1 and 2.

Note that $a(e, e) \leq a(U^h - u, U^h - u)$ is not a statement of convergence must let $h \rightarrow 0$ to show convergence ie $\mathcal{A} = \lim_{h \rightarrow 0} \{\mathcal{A}^h\} \Rightarrow u = \lim_{h \rightarrow 0} \{u^h\}$

Convergence (use best approx + $\mathcal{A}^h \approx \mathcal{A}$).

Hypotheses

$$(1) \exists \text{ constants } c_1, c_2 > 0 \text{ (not fns of } h) \quad c_1 \|v\|_m^2 \leq a(v, v) \leq c_2 \|v\|_m^2$$

this is an equivalence between $\|v\|_m^2$ & $a(v, v)$

where $\|\cdot\|_m = m \text{ norm}$ $\|\cdot\|_0 = \int_{\Omega} (\cdot)^2 d\Omega$; $\|\cdot\|_1^2 = \int_{\Omega} ((\cdot)^2 + [\nabla(\cdot)]^2) d\Omega$
(Sobolev norm).

$$(2) \mathcal{A}^h \approx \mathcal{A} \text{ given any } U \in \mathcal{A}, \exists U^h \in \mathcal{A}^h \Rightarrow \|U^h - U\|_m \leq c_3 h^\alpha \|U\|_s$$

where c_3 is a constant (not a fn of h); where $\alpha > 0$ constant; as $h \rightarrow 0$.

$s \geq m$

m = order of differentiability of $\mathcal{A}$

for most of our work $m=1$ (for bernoulli-euler $m=2$).

Proof of convergence

$$c_1 \|e\|_m^2 \leq a(e, e)$$

using (1)

by best approx
by (1)

$$a(e, e) \leq a(U^h - u, U^h - u) \quad \neq U^h \in \mathcal{A}^h$$

$$\leq c_2 \|U^h - u\|_m^2$$

now pick $U^h \Rightarrow (2)$ holds.
 $\therefore c_2 \|U^h u\|_m^2 \leq c_2 c_3^2 h^{2\alpha} \|u\|_s^2$

now from $c_1 \|e\|_m^2 \leq c_2 c_3^2 h^{2\alpha} \|u\|_s^2$

$$\|e\|_m \leq \left(\frac{c_2 c_3^2}{c_1} \right)^{1/2} h^\alpha \|u\|_s$$

now $\left(\frac{c_2 c_3^2}{c_1} \right)^{1/2} \& \|u\|_s$ are not fun of h & are fixed.
 thus as $h \rightarrow 0$ $\|e\|_m \sim h^\alpha \rightarrow 0$

if c_1 is a fun of h (c_1 is a measure of positive definiteness). As $h \rightarrow 0$ $c_1 \rightarrow 0$
 i.e. lowest eigenvalue goes thru 0; thus the matrix is rank deficient, or rigid body modes are not taken into account.

if c_2, c_3 are very huge then $\Rightarrow h \ll 1$ for the effect of c_2, c_3 to be minimized;
 this type of effect occurs in structural modes.

Hence for $\|e\|_m \leq \left(\frac{c_2 c_3^2}{c_1} \right)^{1/2} h^\alpha \|u\|_s$ h^α means something if c_2, c_3, c_1 are well behaved

... Prof Hughes will give us the results of all this for some of the elements we have looked at.

We now look at problems where things go wrong.

§ 1. Incompressible elasticity & Stokes flow.

(i.e. rubber) Incompressibility = elasticity w/ modification $u_{i,i} = 0$
 stiff bulk mod wrt shear mod. (stress deviator)

Now $\sigma_{ij} = -p \delta_{ij} + c'_{ijkl} \epsilon_{kl}$. we will assume c'_{ijkl} have same properties as c_{ijkl} of last quarter.

p is unknown & must add one more equation ($\nabla \cdot u = 0$)

Now (5) $\sigma_{ij,j} + f_i = 0$ w/ $\nabla \cdot u = 0$ in Ω

and $g_i = u_i$ on Γ_{g_i}
 $\sigma_{ij} n_j = h_i$ on Γ_{h_i}

p appears in the bdy & goes jag.

Stokes flow - steady creeping viscous flow.

if we interpret u_i to be velocity, then eqns. ^{are same as} isotropic incompressible elasticity
 ie $c_{ijkl} = \mu (\delta_{ik}\delta_{jl} + \delta_{ij}\delta_{kl})$. where μ is dynamic viscosity.

then $\int_{\Omega} u_{i,j} d\Omega = \int_{\Gamma} u_i n_j d\Gamma = \int_{\Gamma} g_j n_j d\Gamma$ if $\Gamma_{g_i} = \Gamma$

since $u_{i,j} = 0 \Rightarrow \int_{\Gamma} g_j n_j d\Gamma = 0$

16 April 1991

Orientation lecture of CIT Computing Monday 7:30-9:30 PM III Polya
 Jim Wilczak

Examples of Sobolev norms for our problems. ie $\|u\|_1^2 = \int_0^1 (u^2 + u_{xx}^2) dx$; $\|u\|_2^2 = \int_0^1 (u^2 + u_{xx}^2 + u_{xxx}^2) dx$
 $\xrightarrow{h}$ (2 node elements)

1. Linear element 1-D

$$\|u - U^h\|_1 \leq ch^{\frac{1}{2}} \|u\|_2$$

$$\|u - U^h\|_0 \leq ch^2 \|u\|_2$$

if we do not make any restrictions of the deriv then we approximate u better. (

2. Quad elements

$$\|u - U^h\|_1 \leq ch^2 \|u\|_3$$

$$\|u - U^h\|_0 \leq ch^3 \|u\|_3$$

U^h is quadratic 3 node element.

All these results are independent of finite elements. The Sobolev norms tell you that if a U^h that approximates u very well but doesn't tell you how to get it. The FE method is then the mechanism to find this U^h .

for higher node elements = n

$$\|u - U^h\|_1 \leq ch^{n-1} \|u\|_n$$

$$\|u - U^h\|_0 \leq ch^n \|u\|_n$$

3. for Hermite cubics (displ + 1st deriv at nodes needed)



$Q(1, \cdot)$ involves 2nd deriv (ie beam problem)

$$\|u - U^h\|_2 \leq ch^2 \|u\|_4$$

$$\|u - U^h\|_1 \leq ch^3 \|u\|_4$$

$$\|u - U^h\|_0 \leq ch^4 \|u\|_4$$

then that is equiv. to $\| \cdot \|_2$

4. Isoparametric elements

(since we have slope discont. then m is for 0 only

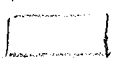
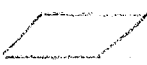
Always $\|u - U^h\|_1 \leq ch^{\frac{1}{2}} \|u\|_2$

$$\|u - U^h\|_0 \leq ch^2 \|u\|_2$$

parent domain satisfy completeness of linear polyn.

if $x(\xi)$ is linear then $\|u - U^h\|_1 \leq ch^k \|u\|_{k+1}$ where k is the degree of complete.

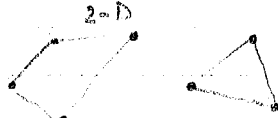
$\| \cdot \|_1, \| \cdot \|_0, \dots, \| \cdot \|_{k+1}, \dots$ } both

linear mapping is  or  from  stretching or rotation

if $X(\xi)$ is a "perturbation" of a linear mapping then the above is still true
 "perturb" $\Rightarrow$ not too distorted

for $k=1$
 (complete linear)

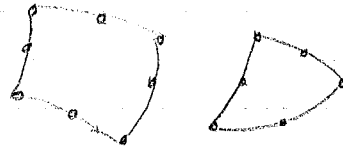
1-D



3-D



for $k=2$
 (complete quadratic poly)



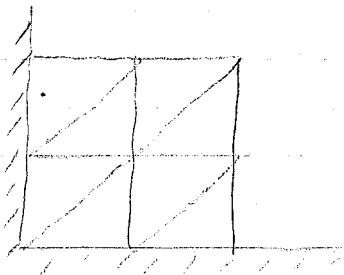
20 node brick



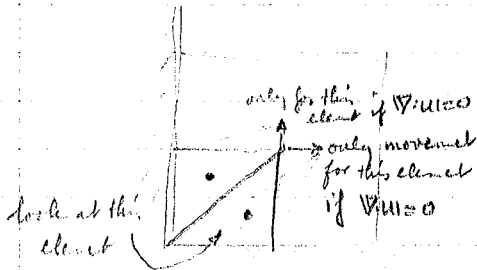
Now back to incompressibility and Stokes Flow.

from last time we had unknowns u_i , and new one p ($\frac{1}{2} \sigma_{ii}$)
 and the normal eqns + the new one $\nabla \cdot u = 0$.

Suppose we have the following problem.

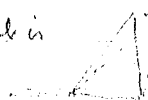


linear triangles.
 and zero displ prescribed on bdy.



only for this element if $\nabla \cdot u = 0$
 only movement for this element if $\nabla \cdot u = 0$
 then this means the point is fixed $u = 0$
 if we go to the next set - the same is true $\Rightarrow u = 0$
 $\Rightarrow \forall$ nodes $u = 0$

The phenomenon is known as locking

if only movement, then $\nabla \cdot u \neq 0$
 since $\nabla \cdot u$ is the area of triangle, which increases by this movement.
 $\Rightarrow$ only allowable movement is ; but this is incompatible

or adjacent element $\Rightarrow u = 0$ at node.

Suppose we let $u_i^h \approx u$ and we put this into our variational principle
 $p^h \approx p$ we get

$$a(w, u) = \left(\int_{\Omega} w_{(i,j)} C_{ijkl} u_{(k,l)} d\Omega \right) - (V \cdot w, p) = \text{usual RHS}$$

for the $V \cdot u = 0$ eq. the variational principle leads to $(q, V \cdot u) = 0$
 where q is the weighting fun. corresponding to p .

also $(V \cdot w, p) = \int_{\Omega} w_{i,i} p d\Omega$

in Galerkin method // replace w by w^h ; u by u^h ; p by p^h ; q by q^h where
 $q^h, p^h \in P^h$ (space of fns).

Now after the matrix construction we get

$$\begin{bmatrix} \underline{K}_1 & \underline{K}_2 \\ \underline{K}_2^T & \underline{Q} \end{bmatrix} \begin{bmatrix} \underline{d} \\ \underline{p} \end{bmatrix} = \begin{bmatrix} \underline{F} \\ \underline{0} \end{bmatrix} \quad \underline{F} \text{ is the RHS.}$$

where $\underline{K}_1$ corresponds to C_{ijkl} and $\underline{K}_2^T \underline{d} = \underline{0}$ represents $V \cdot u = 0$
 now we must ask what is P^h ? that is the problem.

1. Generally we get disastrous results if $P^h = A^h$ (ie same type interpol as A^h).
2. This has led to P^h having lower order than A^h ie continuous fns.
 example: 9 node elements w/ bilinear pressures (biquadratic displacements)
 7 node triangle w/ linear pressures.

x - press node
 . - u node



3. Since $V \cdot u$ is discont. from element to element we can take p discontinuous also
 example 9 node displ. + constant p .



u biquad $u = \dots$
 p - const. $p = x$

P - linear 9 node displ.
 $p = \alpha_0 + \alpha_1 x + \alpha_2 y$



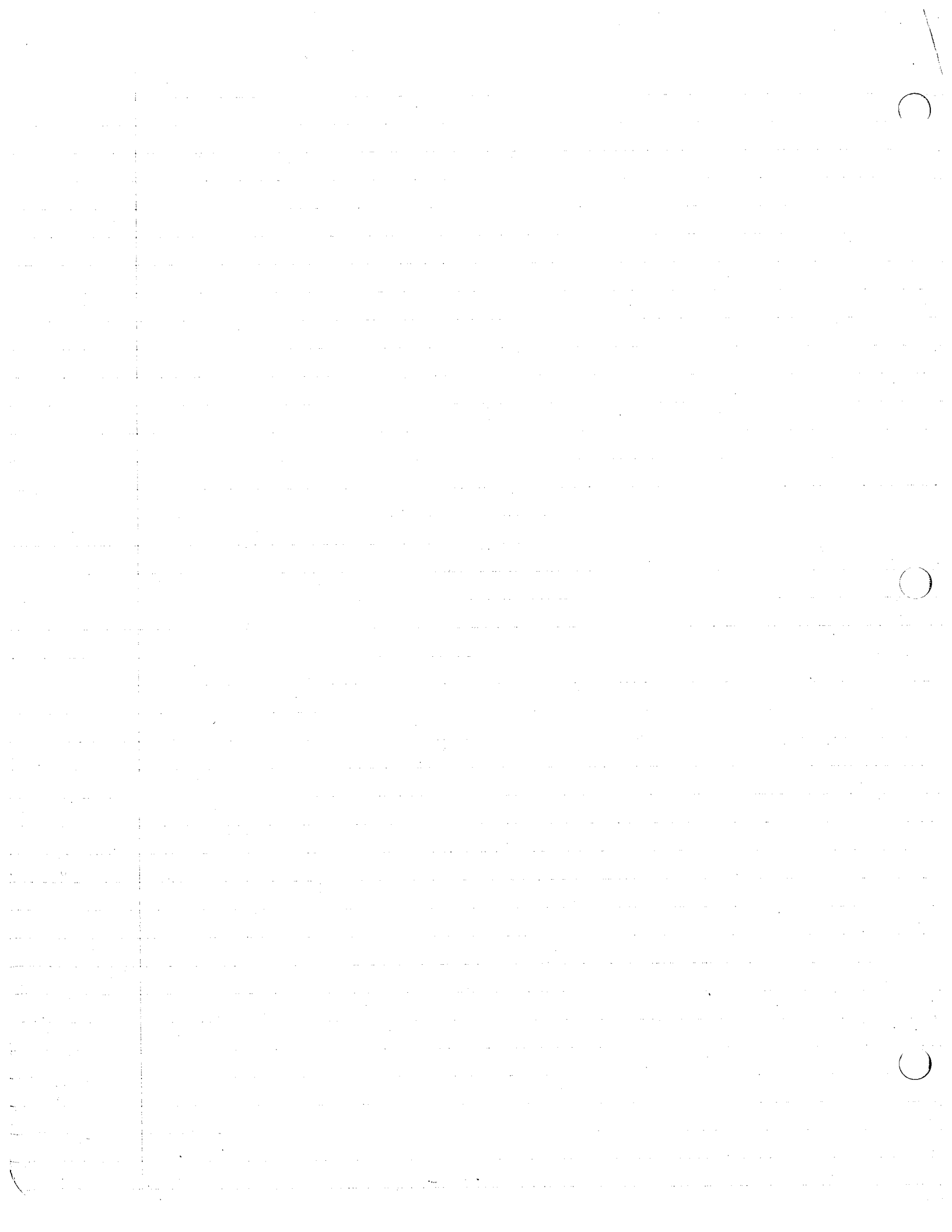
u quadratic $u = \dots$
 p linear this is better for optimality
 d discont. from elements

optimality in sense of 2 u_i eqns + 1 p eqn $\therefore$ ratio 2/1 in 2-D

if we pick constant p and bilinear u



but has pressure oscill w/ some problems ie "checkerboard pressure"



PENALTY FUNCTION FORMULATION

4/21/81

This will provide an approximation for incompressible elasticity / Stokes flow which can be used in practice

(1) Replace $p = -\frac{1}{\epsilon} \text{div } u_i^{(\epsilon)}$ (ϵ is a very small parameter)
 $= -\frac{1}{\epsilon} \nabla \cdot u$

(2) Drop $u_{,i} = 0$

This leads to the following

Problem: given f_i, g_i, h_i , etc,

$$(S^{(\epsilon)}) : \sigma_{ij,j}^{(\epsilon)} + f_i = 0$$

$$u_i^{(\epsilon)} = g_i \quad \text{on } \Gamma_{g_i}$$

$$\sigma_{ij}^{(\epsilon)} n_j = h_i \quad \text{on } \Gamma_{h_i}$$

$$\sigma_{ij}^{(\epsilon)} = -\frac{1}{\epsilon} u_{k,k}^{(\epsilon)} \delta_{ij} + C'_{ijkl} u_{k,l}^{(\epsilon)} = p \delta_{ij} + C_{ijkl} \epsilon_{kl}$$

$$= C_{ijkl} u_{k,l}^{(\epsilon)}$$

$$\text{where } C_{ijkl} = C'_{ijkl} + \frac{1}{\epsilon} \delta_{ij} \delta_{kl}$$

Note as $\epsilon \rightarrow 0$, $(H^1) u_i^{(\epsilon)} \rightarrow u_i = \text{soln. of incomp. problem}$

$(L_2) p^{(\epsilon)} \rightarrow p = \text{soln. of incomp. problem}$

We have changed the incompressible problem to a slightly compressible problem

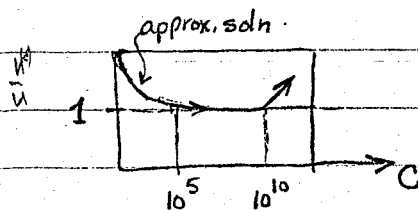
If we assume the isotropic case, note that $\frac{1}{\epsilon}$ corresponds to λ . λ acts like a bulk modulus in the incompressible limit

Rule of thumb for size of ϵ :

$$\frac{1}{\epsilon} = \lambda = \frac{C_{\max} C'_{ijkl}}{2\mu}$$

where C is a constant

so problem is relatively insensitive to C in the range $10^5 - 10^{10}$



"1" = actual soln.

Double prec. IBM - 64 bit

Single prec. CDC - 60 bit

computer info - why CDC is better

The definitions of λ & μ are

$$\lambda = \frac{\nu E}{(1+\nu)(1-2\nu)}$$

$$\mu = \frac{E}{2(1+\nu)}$$

As $\nu \rightarrow \frac{1}{2}$, $\lambda \rightarrow \infty$

For Plane stress - there is no incompressible problem

$$\begin{aligned} \bar{E} &= E/(1-\nu^2) \\ \bar{\nu} &= \nu/(1-\nu) \end{aligned} \quad \left. \begin{array}{l} \text{As } \nu \rightarrow \frac{1}{2}, \bar{E} \text{ \& } \bar{\nu} \text{ are} \\ \text{finite} \end{array} \right\}$$

For Plane strain - problem still blows up

Recall the stiffness calculation

$$k_{ab}^e = \int_{\Omega^e} \underline{B}_a^T \underline{D} \underline{B}_b d\Omega$$

$$\underline{D} = \underline{D}_\lambda + \underline{D}_\mu$$

$\underline{D}_\lambda$ is locking up the Kinematic in the incompressible mode

$$= \underbrace{\int_{\Omega^e} \underline{B}_a^T \underline{D}_\lambda \underline{B}_b d\Omega}_{\text{must do something special here to avoid disaster}} + \underbrace{\int_{\Omega^e} \underline{B}_a^T \underline{D}_\mu \underline{B}_b d\Omega}_{\text{no problem encountered by this calculation}}$$

For $n_{sd} = 2$ (plane strain): $\underline{D}_\lambda = \lambda \begin{bmatrix} 1 & 1 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix}$

$$\underline{D}_\mu = \mu \begin{bmatrix} 2 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

For $n_{sd} = 3$:

$$\underline{D}_\lambda = \begin{bmatrix} 1 & 1 & 1 & | & 0 \\ 1 & 1 & 1 & | & 0 \\ 1 & 1 & 1 & | & 0 \\ \hline 0 & 0 & 0 & | & 0 \end{bmatrix}$$

$$\underline{D}_\mu = \begin{bmatrix} 2 & & & | & 0 \\ & 2 & & | & 0 \\ & & 2 & | & 0 \\ \hline 0 & 0 & 0 & | & 1 \end{bmatrix}$$

In the non-isotropic case,
 $(D_{\mu})_{IJ} = C'_{ijkl}$

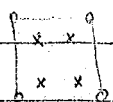
One possible way to develop an improved formulation:
REDUCED/SELECTIVE INTEGRATION

- use different rules on different terms

$$k_{ab}^e \approx \sum_{l^{(1)}=1}^{n_{int}^{(1)}} (\underbrace{B_a^T}_{\sim \lambda} \underbrace{D_{\mu}}_{\sim \mu} \underbrace{B_b}_{\sim \mu})_{l^{(1)}} W_{l^{(1)}}^{(1)} + \sum_{l^{(2)}=1}^{n_{int}^{(2)}} (\underbrace{B_a^T}_{\sim \lambda} \underbrace{D_{\mu}}_{\sim \mu} \underbrace{B_b}_{\sim \mu})_{l^{(2)}} W_{l^{(2)}}^{(2)}$$

- rule 1 is lower than rule 2

- see PLANE1, pp. 67-68 of LEARN manual



There is a normal rule that goes with an element,
 i.e. 2x2 for 4-node bilinear el.

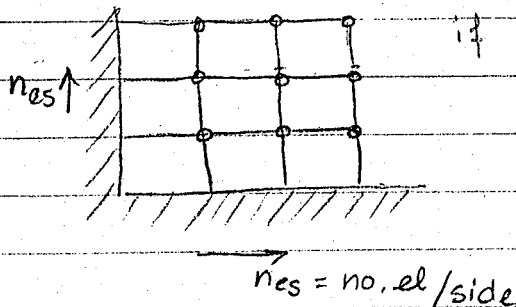


A reduced rule means a lower order,
 i.e. 1 pt. rule for 4-node quad.

Uniform reduced integration involves a 1-pt.
 rule on λ & μ terms (in 4-node example)

Selective reduced integration would involve a
 1-pt. rule on λ and 2x2 on μ

Constraint counting



n_{eq} = no. equations

n_c = no. constraints imposed
 by incompressibility

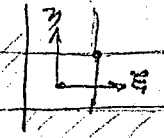
$$\text{Constraint index} = \frac{n_{eq} - n_c}{n_{es}^2} \quad (CI)$$

$$\text{Constraint ratio} = \frac{n_{eq}}{n_c} \quad (CR)$$

All successful elements

have constraint ratio > 2

(= 2 for 2-D
 continuum)



$n_{eq} = 2$

$x = x(\xi, \eta) = x(\xi)$

$y = y(\xi, \eta) = y(\eta)$

eg. $\frac{1}{2}(1-\eta)y_1 + \frac{1}{2}(1+\eta)y_2$
 a linear function
 physically, just from stretching

$u_i^h = u_i^h(\xi, \eta)$
 bilin.
 $= u_i^h(x, y)$
 also bilin.

Take $0 \cong u_{i,i}^h$

$u_i^h = \alpha_{0i} + \alpha_{1i}x + \alpha_{2i}y + \alpha_{3i}xy$

$\alpha_{ji} = \text{constants}$
 = linear combs. of nodal values d_{ji}

$u_{i,i}^h = u_{1,1}^h + u_{2,2}^h = \alpha_{11} + \alpha_{22} + \alpha_{31}y + \alpha_{32}x \cong 0$

$= 0 \Rightarrow \alpha_{11} + \alpha_{22} = 0$

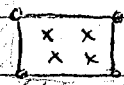
$\alpha_{31} = 0$

$\alpha_{32} = 0$

} 3 constraints on nodal values at each point

Thus $n_c/n_d = 3$

4/23/81



CI = -1 { 2 eq/ed
3 constr.

opt. "CR = 2" + 1 node/into
 $2 \times 2 \Rightarrow$ bilin u_i , bilin $u_{ji} \Rightarrow$ 3 const.

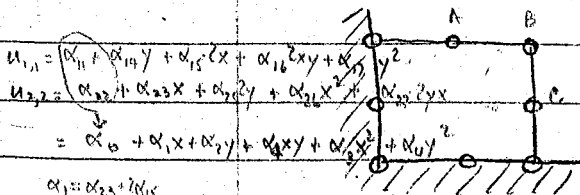


CI = 2 - 1 = 1
 CR = $\frac{2}{1} = 2$

(lambda-term) 1 pt $\Rightarrow u_{iii}^h = \text{const} \Rightarrow$ 1 const.

EXAMPLES: (WE ALWAYS LOOK AT THE NODE IN THE CORNER SINCE THE OTHERS WILL BE THE SAME)

(1) 8-node - standard quadrature is 3×3



6 eqns. (2 at each A/B/C). This is a serendipity element
 terms involved are $1, x, y, xy, x^2, y^2, x^2y, y^2x$
 $0 \cong u_{i,i}^h = \alpha_0 + \alpha_1x + \alpha_2y + \alpha_3x^2 + \alpha_4xy + \alpha_5y^2$
 we need to require all $\alpha_i = 0$ for $u_{i,i}^h = 0$

CI = 6 - 6 = 0

CR = $\frac{6}{6} = 1$

Now try reduced/selective integration. Use 2×2 on lambda-term

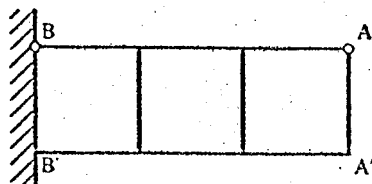
$n_c = 4$ (at most)

Since 2×2 is bilin and has 1 unknown at each of the quad pts

CT = 1 - 1 - 2

$\frac{6}{2} = 3$ better but not up to 2 yet

Probably the most serious economic problem of complex curvilinear elements is the computer time necessary for performing the numerical integrations. Here some economic limit on the accuracy required in this integration must obviously be imposed.



Type of element	Vertical Load of A		Couple at AA'	
	Max. defl. at AA'	Max. stress BB'	Max. defl. at AA'	Max. stress BB'
	0.26	0.19	0.22	0.22
	0.65	0.56	0.67	0.67
	0.53	0.51	0.52	0.55
	0.99	0.99	1.00	1.00
	1.00	1.00	1.00	1.00
EXACT	1.00	1.00	1.00	1.00

Fig. 9.1 A cantilever in plane stress analysed by various elements. Accuracy improvement with higher order elements

9.2 Required Accuracy of Numerical Integration

In the previous chapter it has been indicated how the element matrices can be formed using numerical integration in terms of n Gauss points. The effort of this integration over plane area is roughly proportional to n^2 —the number of points at which the function has to be found—while in a three-dimensional situation it is proportional to n^3 . The determination of an adequate minimum number of Gauss points is thus of some importance.

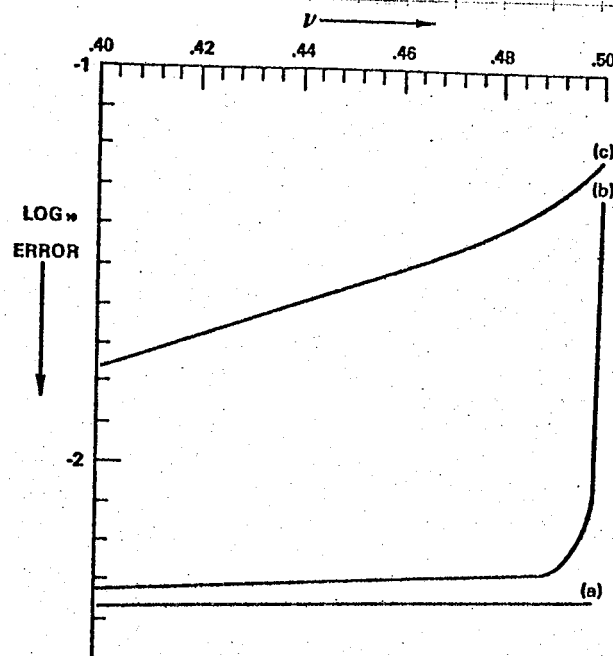
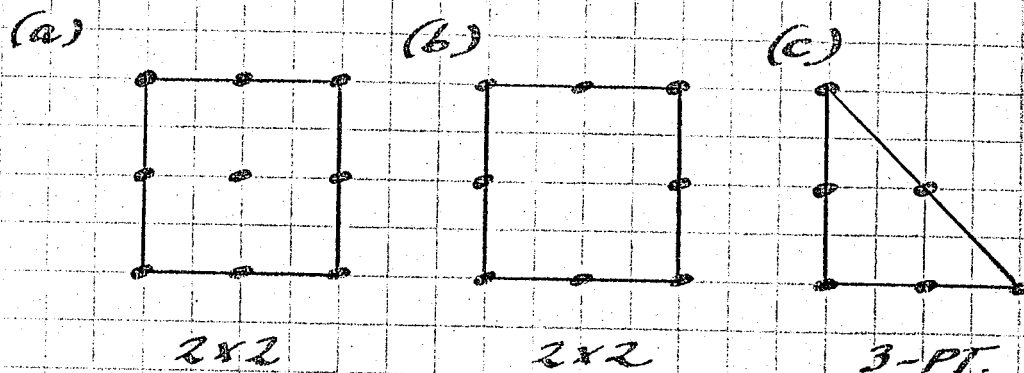


Fig. 3. $\|u^h - u\|$ —error in approximating the exact solution with the artificial ν as this $\nu \rightarrow \frac{1}{2}$, for three types of quadratic elements, integrated with the formulas of Table 1: (a) Lagrange, (b) serendipity, (c) triangle.



From D.S. Malkus, 'A Finite Element Displacement Model Valid for any value of the Compressibility', Int. J. Solids Structures, Vol. 12, 731-738 (1976).

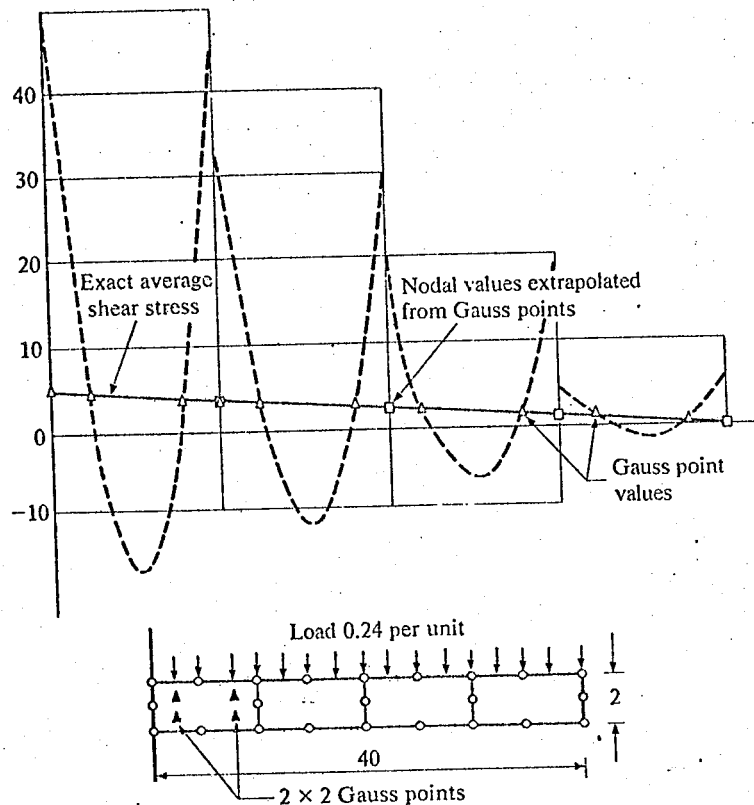


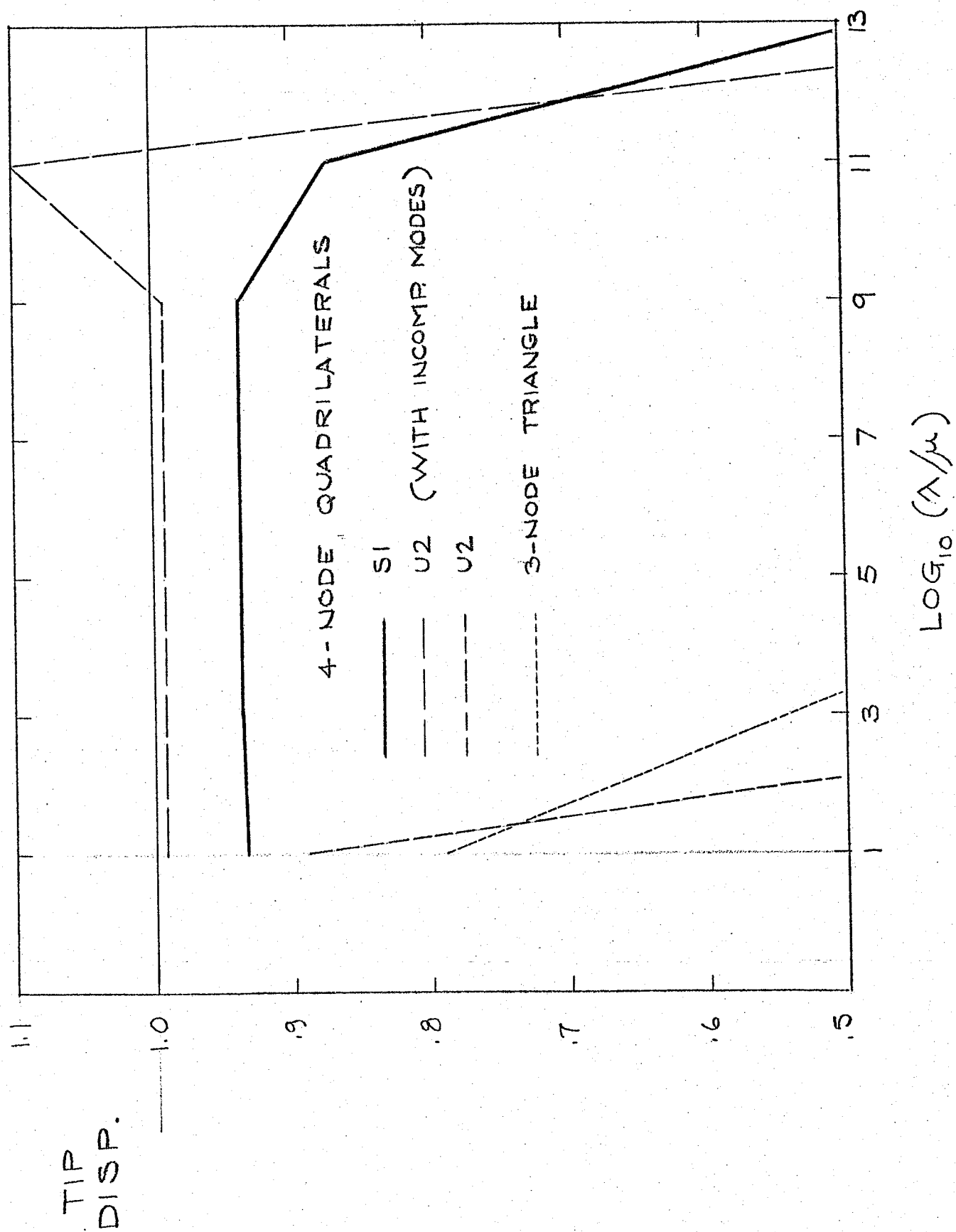
Fig. 11.9 A cantilever beam with four parabolic elements. Stress sampling at Gauss points and linear extrapolation to nodes

In Fig. 11.9 we show, for instance an analysis of a cantilever beam using four quadratic, 'serendipity' type, elements. Whilst the results for deflections and axial stresses are excellent, the shear stresses show a parabolic 'variation' in each element which provides an extremely poor representation of the actual stresses. However, the values sampled at the nodes are an excellent representation of the correct mean shear stresses.

Similar improvements can be shown in the context of other elements and problems, although (fortunately) the discrepancies are not always so large.

The example just quoted suggests that in quadratic C_0 elements, whether two- or three-dimensional, the stresses (or similar quantities) *should never be calculated at nodes*. If nodal values are desired, then a simple bilinear extrapolation from Gauss points should be made. Such values are again shown to be excellent in Fig. 11.9. Further examples of





32 ELEMENTS

Q4

λ	EXACT	2X2 INCOMP		2X2 2X2		2X2 1X1	
		DISP	DISP/EXACT	DISP	DISP/EXACT	DISP	DISP/EXACT
10	74.27	73.77	0.993	66.02	0.889	69.13	0.931
10^3	68.56	68.03	0.992	15.34	0.224	64.21	0.937
10^5	68.50	67.96	0.992	4.14	0.06	64.16	0.937
10^7	68.50	67.96	0.992	4.00	0.058	64.16	0.937
10^9	68.50	67.94	0.992	4.00	0.058	64.14	0.937
10^{11}	68.50	75.22	1.098	5.09	0.074	60.01	0.876
10^{13}	68.50	10.87	0.158	17.09	0.25	5.95	0.087
10^{15}	68.50	0.57	0.008	0.39	0.006	0.787	0.011

λ	E	ν
10	2.9090909090909090	0.4545454545454545
10^3	2.999000999000999000	0.4995004995004995
10^5	2.999990000099999	0.4999950000499995
10^7	2.9999999000000001	0.4999999500000005
10^9	2.999999999000000	0.499999999500000
10^{11}	2.999999999990000	0.499999999995000
10^{13}	2.999999999999900	0.499999999999950
10^{15}	2.999999999999999	0.4999999999999995

64 ELEMENTS

TRIANGLES

λ	EXACT	2X2 INCOMP		2X2 2X2		2X2 1X1	
		DISP	DISP/EXACT	DISP	DISP/EXACT	DISP	DISP/EXACT
10	74.27	58.555	0.788	58.555	0.788	58.555	0.788
10^3	68.56	36.905	0.538	36.905	0.538	36.905	0.538

shape functions	bilinear	biquadratic	bicubic
λ -term	1 point	2×2	3×3
μ -term	2×2	3×3	4×4

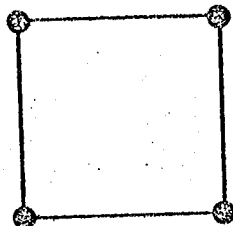
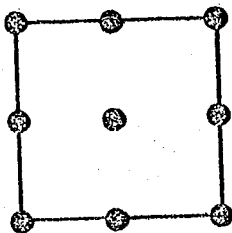
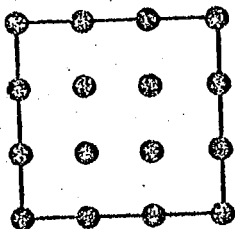
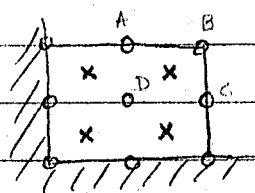


Figure 3. Selective Gauss-Legendre integration rules for 2-dimensional isoparametric Lagrange elements.

(2) 9-node element ("everybody's darling")



8 eqns. (2 at each A, B, C, D)

8 fns. from serendipity element +

 x^2y^2 from bubble

$$0 \cong u_{i,i}^h = \underbrace{\text{QUADRATIC}}_{6 \text{ param's}} + \alpha_6 xy^2 + \alpha_7 y^2 x$$

$$CI = 8 - 8 = 0$$

$$CR = 8/8 = 1$$

RED/SEL INT: 2×2 on λ

$$n_c = 4$$

linear in p is $u_{i,i}$ involves $1, x, y, xy$ fns

$$CI = 8 - 4 = 4$$

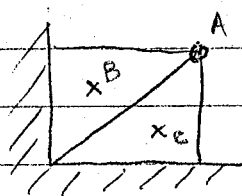
$$CR = \frac{8}{4} = 2$$

(3) TRIANGLE (LINEAR)

4 (2 from each triangle) less 2 (constraint from pt p same) = 2

why

(a)

 $n_{eq} = 2$ (2 each at A)

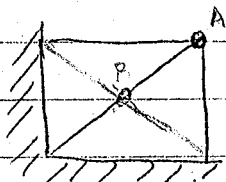
$$CI = 2 - 2 = 0$$

 $n_c = 2$ (1 each at B, C)

$$CR = 1? = 2/2 \text{ wrt rectangular}$$

 $u_{i,i}$ is constant in each element - "lock-up"

(b)



why

 $n_{eq} = 4$ (2 each at A, B) $n_c = 3$ { because of linear dependence of the original 4

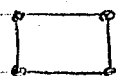
$$CI = 1$$

$$CR = \frac{4}{3}$$

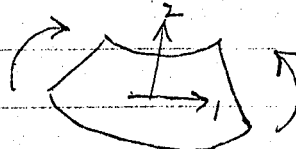
this "more-or-less" works

(4) INCOMPATIBLE MODES

(a)



does not bend well - usual isoparametric element



$$u_1 = cxy$$

$$u_2 = \frac{c_1}{2}(a^2 - x^2)$$

$$+ c_2(b^2 - y^2)$$

is this the actual solution?



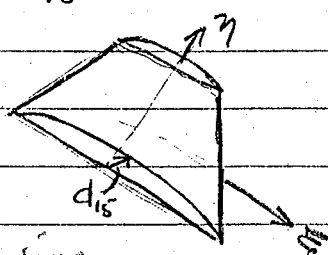
bilinear element response

Wilson et al. (Berkeley) have suggested that you add the ϵ^2 and γ^2 deformation patterns

$$u_i^h = \sum_{a=1}^4 N_a(\epsilon, \gamma) d_{ia}^e + \sum_{a=5}^6 N_a \hat{d}_{ia}^e$$

$N_5 = 1 - \epsilon^2$
 $N_6 = 1 - \gamma^2$
 $\nwarrow$ not nodal parameters

$$u_i^h = \hat{d}_{i5} (1 - \epsilon^2)$$



This element is not continuous (that's what incompatible means)

this changes the following

$$\underline{\underline{k}}^e = \int_{\Omega^e} \underline{\underline{B}}^T \underline{\underline{D}} \underline{\underline{B}} d\Omega \quad \text{where} \quad \underline{\underline{B}} = [\underline{\underline{B}}_1 \dots \underline{\underline{B}}_4 \mid \underline{\underline{B}}_5 \quad \underline{\underline{B}}_6]$$

$\underline{\underline{k}}^e$ is 12×12
 $\underline{\underline{B}}$ is 3×2

$$\underline{\underline{k}}^e = \begin{bmatrix} \underline{\underline{k}} & \hat{\underline{\underline{k}}} \\ \hat{\underline{\underline{k}}}^T & \hat{\underline{\underline{k}}} \end{bmatrix}$$

$\underline{\underline{k}}$ is 8×8
 $\hat{\underline{\underline{k}}}$ is 4×4

$\hat{\underline{\underline{k}}}$ is whatever is in lower box

Assume that the generalized forces corresponding to the incompressible modes ($\hat{d}_{i5}, \hat{d}_{i6}$) are zero

$$\hat{\underline{\underline{k}}}^T \underline{\underline{d}}^e + \hat{\underline{\underline{k}}} \hat{\underline{\underline{d}}} = 0$$

$\hat{\underline{\underline{k}}}^T$ is 4×8
 $\underline{\underline{d}}^e$ is 8×1
 $\hat{\underline{\underline{k}}}$ is 4×4
 $\hat{\underline{\underline{d}}}$ is 4×1

Solve the eqn. for $\hat{\underline{\underline{d}}}$:

$$\hat{\underline{\underline{d}}} = -\hat{\underline{\underline{k}}}^{-1} \hat{\underline{\underline{k}}}^T \underline{\underline{d}}^e$$

Substitute in: $\underline{\underline{k}} \underline{\underline{d}}^e + \hat{\underline{\underline{k}}} \hat{\underline{\underline{d}}} = \underline{\underline{f}}^e$

$$(\underline{\underline{k}} - \hat{\underline{\underline{k}}} \hat{\underline{\underline{k}}}^{-1} \hat{\underline{\underline{k}}}^T) \underline{\underline{d}}^e = \underline{\underline{f}}^e$$

Stiffness with incompressible modes

This proc. is called STATCON (static condensation)

And is included in LEARN

1. Coding - add N_5, N_6 to shape fn. routine
2. Increase nodal loops from 4 to 6

looking to do the above.

??! $\rightarrow$ 3. Static condensation

$$\sigma_{ij,j} + f_i = 0 \Rightarrow \sigma_{ii,i} + \sigma'_{ij,j} + f_i = 0$$

$$\sigma_{ii,i} = c_{iiii} u_{i,i}$$

$$\int \omega_i (\sigma_{ij,j}) d\Omega = \int \omega_i f_i d\Omega = 0$$

$$\nabla \cdot (w \sigma) = \sigma \cdot \nabla w$$

$$\int \omega_i (\sigma_{ii,i}) + \int \omega_i (\sigma'_{ij,j})$$

$$\int [(\sigma_{ii} \omega_i)_{,i} - \sigma_{ii} \omega_{i,i}] d\Omega = \int \omega_i (\sigma'_{ij,j}) d\Omega$$

$$\int \sigma_{ii} \omega_i n_i d\Gamma - \int \sigma_{ii} \nabla \cdot w$$

$$\int \omega_i \sigma_{ij,j} + \int \omega_i f_i d\Omega = 0$$

$$\int (\omega_i \sigma_{ij})_{,j} - \int (\omega_{i,j} \sigma_{ij}) + \int \omega_i f_i d\Omega = 0 \quad \text{using } -p \delta_{ij} + C_{ijkl} \varepsilon_{kl} = \sigma_{ij}$$

$$\downarrow + \int \omega_{i,j} p \delta_{ij} d\Omega - \int \omega_{i,j} C_{ijkl} u_{(k,l)} d\Omega + \dots = 0$$

$$\int (\omega_i \sigma_{ij} n_j) d\Gamma + \int \omega_{i,j} p d\Omega - \int \omega_{(i,j)} C_{ijkl} u_{(k,l)} d\Omega + \dots = 0$$

$\Rightarrow$

$$\int_{\Omega} \omega_{(i,j)} C_{ijkl} u_{(k,l)} d\Omega - \int_{\Omega} \nabla \cdot w p d\Omega = \int_{\Omega} \omega_i f_i d\Omega + \int_{\Gamma_{hi}} \omega_i h_i d\Gamma \quad h_i = \sigma_{ij} n_j \text{ on } \Gamma_{hi}$$

$$\text{if } \nabla \cdot u = 0 \Rightarrow \int_{\Omega} q u_{i,i} d\Omega = 0 \Rightarrow \int [(q u_i)_{,i} - (q_{,i} u_i)] d\Omega = 0 \Rightarrow (q, \nabla \cdot u) = 0$$

$$\Rightarrow \int_{\Gamma_{hi}} q u_i n_i d\Gamma = \int_{\Omega} u \cdot \nabla q d\Omega$$

$$\int_{\Gamma_{hi}} q q_{,i} n_i d\Gamma = \int_{\Omega} u \cdot \nabla q d\Omega$$

$$\text{Note that } \int_{\Omega} \omega_{(i,j)} C_{ijkl} u_{(k,l)} d\Omega = \int \underline{B}_a \underline{D}_{\mu} \underline{B}_b d\Omega$$

$$-\int_{\Omega} \nabla \cdot w \left(-\frac{1}{\varepsilon} \nabla \cdot u \right) d\Omega = \int \underline{B}_a \underline{D}_{\lambda} \underline{B}_b d\Omega \quad \text{where } \frac{1}{\varepsilon} = \lambda$$

$$\nabla \cdot w = \begin{pmatrix} N_{a,x} & 0 \\ 0 & N_{a,y} \\ 0 & 0 \end{pmatrix} \begin{pmatrix} \delta_{11} \\ \delta_{12} \end{pmatrix} \quad \text{since } \nabla \cdot w = w_{1,1} + w_{1,2}$$

1. E.L. Wilson, R.L. Taylor, W.P. Doherty and J. Chakraborti, "Incompatible Displacement Models," pp. 43-57 in Numerical and Computer Methods in Structural Methods, Edited by J. Fanning, W. Parnes, A.R. Robinson and W.C. Schnobrich, Academic Press, 1973.
2. P. Lesaint, "On the Convergence of Wilson's Non-conforming Elements for Solving the Elastic Problems," Computer Methods in Applied Mechanics and Engineering, Vol. 7, 1-16 (1976).
3. R.L. Taylor, P.J. Benford and E.L. Wilson, "A Nonconforming Element for Stress Analysis," International Journal for Numerical Methods in Engineering, Vol. 10, No. 6, 1211-1220 (1976).
4. D. Kooloff and G.A. Frazier, "Treatment of Hourglass Patterns in Low Order Finite Element Codes," Numerical and Analytical Methods in Geomechanics, Vol. 2, 57-72 (1978).

see handout of next lesson

4/28/81 Comments on incompatible modes

- 1) The element fns. are no longer continuous
- 2) Elements converge for linear transformations
- 3) Taylor's method - change evaluation of shape fns. only for incompatible modes

$$\tilde{B}_a = \begin{bmatrix} N_{a,1} & 0 \\ 0 & N_{a,2} \\ N_{a,2} & N_{a,1} \end{bmatrix}$$

$$\tilde{B}_5 = \frac{j}{2} \begin{bmatrix} -\xi(y,\eta) & 0 \\ 0 & \xi(x,\eta) \\ \xi(x,\eta) & -\xi(y,\eta) \end{bmatrix}$$

$$B_6 = \frac{j}{2} \begin{bmatrix} \eta(y,\xi) & 0 \\ 0 & -\eta(x,\xi) \\ -\eta(x,\xi) & \eta(y,\xi) \end{bmatrix}$$

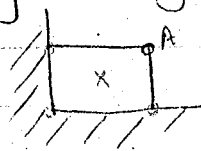
normal method

Taylor - substitute
in $\bigcirc$ their values
at $\xi = \eta = 0$

This is contained in
LEARN in SHAP2φ

Remarks

1. This method works well for bending
2. It is a bit experimental
3. It can be generalized to 3-D with
 $1 - \xi^2, 1 - \eta^2, 1 - \zeta^2$
4. The axisymmetric case does not converge
5. We get a good constraint index



trial $w^h = 1, x, y, x^2, xy, y^2$

$0 = u_{,x} \Rightarrow 3$ constraints

C.I. = $6 - 3 = 3$

C.R. = $6/3 = 2$

I can get 2 (from 1)
where do they come
from? from $\hat{d}$ in
Taylor formulae
one pt quad on p

ENGINEERING CRITERIA - how good are the results?

1. Invariance - results must be independent of reference configuration
 - this is usually satisfied (we will not go into when it is not)
2. Patch test - implies that the element is complete
 - satisfied by most standard integration rules

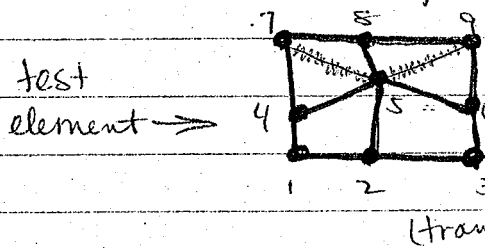
a) Engr. Version - (Irons 1965)

done on computer

implies completeness

if element can represent deformation states

1, x, y in an arbitrary patch of elements, then it passes



Set boundary nodes

Test# d_{1a} d_{2a}

1 { 1 1 0 a: 1 → 4, 2 0 1 b → 9

set these values for the boundary nodes & get the analysis to give you the displacement of node 5.

x { 3 x_a 0 4 0 x_a y { 5 y_a 0 6 0 y_a

A higher order patch test implies higher order completeness

(also a mathematical version due to Strang)

[illegible]

LINEAR TRANSIENT ANALYSIS

Look at the parabolic case - Heat eqn.

$$\begin{aligned}
 &\text{Given: } f: \Omega \times]0, T[\rightarrow \mathbb{R} \quad f: (\cdot, \cdot) \\
 &\quad g: \Gamma_g \times]0, T[\rightarrow \mathbb{R} \\
 &\quad u_i: \Omega \rightarrow \mathbb{R} \\
 &\text{Find: } u: \bar{\Omega} \times [0, T] \rightarrow \mathbb{R} \\
 &\quad \rho c u_t + q_{i,i} = f \quad \text{on } \Omega \times]0, T[\\
 &\quad u = g \quad \text{on } \Gamma_g \times]0, T[\\
 &\quad q_i n_i = -\bar{h} \quad \text{on } \Gamma_h \times]0, T[\\
 &\quad u(x, 0) = u_0 \quad x \in \Omega
 \end{aligned}
 \tag{S}$$

where $q_i = -K_{ij} u_{j,i}$ if $K_{ij} = -K \delta_{ij} \Rightarrow q_i = -K u_{i,i}$
 This converts to the following

$$\begin{aligned}
 &\text{Given: } f, g, \bar{h}, u_0 \text{ as in (S)} \\
 &\text{Find: } u: [0, T] \rightarrow S_{\text{gen}} \quad \forall w \in V \\
 &\quad (\rho c u, w) + Q(u, w) = (f, w) + (\bar{h}, w)_\Gamma \quad (*) \\
 &\quad (u(0) - u_0, w) = 0 \quad (**) \quad \text{see next page}
 \end{aligned}$$

Exercise: Verify the equivalence of (S) & (w)

Note: (*) + usual hypothesis $\Rightarrow$ P.D.E. + 2 b.c.
 (**) $\Rightarrow$ i.c.

$$\begin{aligned}
 &\text{Given: } f, g, \bar{h}, u_0 \text{ as in (S)} \\
 &\text{Find: } u: [0, T] \rightarrow S_{\text{gen}}^h \quad \forall w \in V^h \\
 &\quad (\rho c u^h, w^h) + Q(u^h, w^h) = (f, w^h) + (\bar{h}, w^h)_\Gamma \\
 &\quad (u^h(0) - u_0, w^h) = 0
 \end{aligned}
 \tag{G1}$$

G1 & G2 are the
 semi-discrete
 Galerkin eqns
 (it is still
 continuous)

$$\begin{aligned}
 &v^h: [0, T] \rightarrow V^h \\
 &u^h = v^h + g^h \quad g^h: [0, T] \rightarrow S^h \\
 &\quad (\rho c v^h, w^h) + Q(v^h, w^h) = (f, w^h) + (\bar{h}, w^h)_\Gamma - (\rho c g^h, w^h) \\
 &\quad (u^h(0) - u_0, w^h) = 0
 \end{aligned}
 \tag{G2}$$

$$\int (\underbrace{pc u_t - [\kappa_{ij} u_{,j}]_{,i}}_{q_{i,i}}) w = \int w f d\Omega$$

$$\int (pc u_t w) d\Omega - \underbrace{\int_{\Gamma} w (\kappa_{ij} u_{,j}) d\Gamma n_i}_{\text{Green's Thm.}} = \int w f d\Omega$$

$$(pc u, w) - \int_{\Gamma} w \bar{h} d\Gamma + Q(w, u) = (w, f) + (w, \bar{h})_{\Gamma_{\bar{h}}}$$

move

$$\int w (u(x, 0) - u_0) d\Omega = 0 \Rightarrow (w, u(0) - u_0) = 0$$

now $q_i n_i = \bar{h}$ on $\Gamma_{\bar{h}}$

what about Γ_g
it is assumed that $w=0$ on Γ_g

and if $u = v + \tilde{g} \Rightarrow$

$$(pc v, w) + Q(w, v) = (w, f) + (w, \bar{h})_{\Gamma_{\bar{h}}} - (pc \tilde{g}, w) - Q(w, \tilde{g})$$

$$\text{and } (w, v(0) - v_0) = (w, \tilde{g}_0 - \tilde{g}(0))_{\Gamma_{\bar{h}}}$$

$$(w, v(0)) = (w, u_0) - (w, \tilde{g}(0))$$

$$(pc \sum N_a \dot{a}_a, \sum N_b \dot{c}_b) + Q(\sum N_b \dot{c}_b, \sum N_a \dot{a}_a) = (\sum N_b \dot{c}_b, f) + (\sum N_b \dot{c}_b, \bar{h})_{\Gamma_{\bar{h}}} - (\sum N_b \dot{c}_b, pc \sum N_a \dot{a}_a) - Q(\sum N_b \dot{c}_b, \sum N_a \dot{a}_a)$$

$$\text{then for any } \sum c_b \{ (pc \sum N_a \dot{a}_a, N_b) + Q(N_b, \sum N_a \dot{a}_a) = (N_b, f) + (N_b, \bar{h})_{\Gamma_{\bar{h}}} - (N_b, pc \sum N_a \dot{a}_a) - Q(N_b, \sum N_a \dot{a}_a) \}$$

$$\text{let } \int pc N_a \dot{a}_a d\Omega = [m_{ab}] \quad \int N_b N_a d\Omega = [k_{ab}]$$

$$\underline{M} \underline{\dot{d}} + \underline{K} \underline{d} = \underline{F} \quad \text{where } \underline{F} \text{ is the RHS}$$

$$- \sum m_{ab} \dot{g}_a - \sum k_{ab} g_a$$

and since $u(0) = u_0$

$$\text{then if } u(0) = v(0) + g(0) = \sum N_a \dot{a}_a(0) = u_0 = \sum N_a \dot{a}_{a_0}$$

$$\underline{\dot{d}}(0) = \underline{\dot{d}}_0$$

$$\text{from } (w, v(0)) = (w, u_0) - (w, \tilde{g}(0))$$

$$\Rightarrow (N_b, \sum N_a \dot{a}_{a_0}) = (N_b, u_0) - (N_b, \sum N_a \tilde{g}_a(0))$$

$$\underline{\bar{M}} \underline{\dot{d}}_0 = \int N_b u_0 d\Omega - \underline{\bar{M}} \underline{\tilde{g}}(0) = \underline{\dot{d}}$$

$$\underline{\dot{d}}_0 = \underline{\bar{M}}^{-1} \left[\underbrace{\int N_b u_0 d\Omega}_{\underline{\dot{d}}} - \underline{\bar{M}} \underline{\tilde{g}}(0) \right]$$

$$\underline{\bar{M}} \underline{\tilde{g}}(0) = \sum_{j=1}^{n_{\tilde{g}}} \bar{m}_{pj}^{\tilde{g}} \tilde{g}_j(0)$$

1992, 1993, 1994, 1995, 1996, 1997, 1998, 1999, 2000, 2001, 2002, 2003, 2004, 2005, 2006, 2007, 2008, 2009, 2010, 2011, 2012, 2013, 2014, 2015, 2016, 2017, 2018, 2019, 2020, 2021, 2022, 2023, 2024, 2025, 2026, 2027, 2028, 2029, 2030, 2031, 2032, 2033, 2034, 2035, 2036, 2037, 2038, 2039, 2040, 2041, 2042, 2043, 2044, 2045, 2046, 2047, 2048, 2049, 2050, 2051, 2052, 2053, 2054, 2055, 2056, 2057, 2058, 2059, 2060, 2061, 2062, 2063, 2064, 2065, 2066, 2067, 2068, 2069, 2070, 2071, 2072, 2073, 2074, 2075, 2076, 2077, 2078, 2079, 2080, 2081, 2082, 2083, 2084, 2085, 2086, 2087, 2088, 2089, 2090, 2091, 2092, 2093, 2094, 2095, 2096, 2097, 2098, 2099, 2100, 2101, 2102, 2103, 2104, 2105, 2106, 2107, 2108, 2109, 2110, 2111, 2112, 2113, 2114, 2115, 2116, 2117, 2118, 2119, 2120, 2121, 2122, 2123, 2124, 2125, 2126, 2127, 2128, 2129, 2130, 2131, 2132, 2133, 2134, 2135, 2136, 2137, 2138, 2139, 2140, 2141, 2142, 2143, 2144, 2145, 2146, 2147, 2148, 2149, 2150, 2151, 2152, 2153, 2154, 2155, 2156, 2157, 2158, 2159, 2160, 2161, 2162, 2163, 2164, 2165, 2166, 2167, 2168, 2169, 2170, 2171, 2172, 2173, 2174, 2175, 2176, 2177, 2178, 2179, 2180, 2181, 2182, 2183, 2184, 2185, 2186, 2187, 2188, 2189, 2190, 2191, 2192, 2193, 2194, 2195, 2196, 2197, 2198, 2199, 2200, 2201, 2202, 2203, 2204, 2205, 2206, 2207, 2208, 2209, 2210, 2211, 2212, 2213, 2214, 2215, 2216, 2217, 2218, 2219, 2220, 2221, 2222, 2223, 2224, 2225, 2226, 2227, 2228, 2229, 2230, 2231, 2232, 2233, 2234, 2235, 2236, 2237, 2238, 2239, 2240, 2241, 2242, 2243, 2244, 2245, 2246, 2247, 2248, 2249, 2250, 2251, 2252, 2253, 2254, 2255, 2256, 2257, 2258, 2259, 2260, 2261, 2262, 2263, 2264, 2265, 2266, 2267, 2268, 2269, 2270, 2271, 2272, 2273, 2274, 2275, 2276, 2277, 2278, 2279, 2280, 2281, 2282, 2283, 2284, 2285, 2286, 2287, 2288, 2289, 2290, 2291, 2292, 2293, 2294, 2295, 2296, 2297, 2298, 2299, 2300, 2301, 2302, 2303, 2304, 2305, 2306, 2307, 2308, 2309, 2310, 2311, 2312, 2313, 2314, 2315, 2316, 2317, 2318, 2319, 2320, 2321, 2322, 2323, 2324, 2325, 2326, 2327, 2328, 2329, 2330, 2331, 2332, 2333, 2334, 2335, 2336, 2337, 2338, 2339, 2340, 2341, 2342, 2343, 2344, 2345, 2346, 2347, 2348, 2349, 2350, 2351, 2352, 2353, 2354, 2355, 2356, 2357, 2358, 2359, 2360, 2361, 2362, 2363, 2364, 2365, 2366, 2367, 2368, 2369, 2370, 2371, 2372, 2373, 2374, 2375, 2376, 2377, 2378, 2379, 2380, 2381, 2382, 2383, 2384, 2385, 2386, 2387, 2388, 2389, 2390, 2391, 2392, 2393, 2394, 2395, 2396, 2397, 2398, 2399, 2400, 2401, 2402, 2403, 2404, 2405, 2406, 2407, 2408, 2409, 2410, 2411, 2412, 2413, 2414, 2415, 2416, 2417, 2418, 2419, 2420, 2421, 2422, 2423, 2424, 2425, 2426, 2427, 2428, 2429, 2430, 2431, 2432, 2433, 2434, 2435, 2436, 2437, 2438, 2439, 2440, 2441, 2442, 2443, 2444, 2445, 2446, 2447, 2448, 2449, 2450, 2451, 2452, 2453, 2454, 2455, 2456, 2457, 2458, 2459, 2460, 2461, 2462, 2463, 2464, 2465, 2466, 2467, 2468, 2469, 2470, 2471, 2472, 2473, 2474, 2475, 2476, 2477, 2478, 2479, 2480, 2481, 2482, 2483, 2484, 2485, 2486, 2487, 2488, 2489, 2490, 2491, 2492, 2493, 2494, 2495, 2496, 2497, 2498, 2499, 2500, 2501, 2502, 2503, 2504, 2505, 2506, 2507, 2508, 2509, 2510, 2511, 2512, 2513, 2514, 2515, 2516, 2517, 2518, 2519, 2520, 2521, 2522, 2523, 2524, 2525, 2526, 2527, 2528, 2529, 2530, 2531, 2532, 2533, 2534, 2535, 2536, 2537, 2538, 2539, 2540, 2541, 2542, 2543, 2544, 2545, 2546, 2547, 2548, 2549, 2550, 2551, 2552, 2553, 2554, 2555, 2556, 2557, 2558, 2559, 2560, 2561, 2562, 2563, 2564, 2565, 2566, 2567, 2568, 2569, 2570, 2571, 2572, 2573, 2574, 2575, 2576, 2577, 2578, 2579, 2580, 2581, 2582, 2583, 2584, 2585, 2586, 2587, 2588, 2589, 2590, 2591, 2592, 2593, 2594, 2595, 2596, 2597, 2598, 2599, 2600, 2601, 2602, 2603, 2604, 2605, 2606, 2607, 2608, 2609, 2610, 2611, 2612, 2613, 2614, 2615, 2616, 2617, 2618, 2619, 2620, 2621, 2622, 2623, 2624, 2625, 2626, 2627, 2628, 2629, 2630, 2631, 2632, 2633, 2634, 2635, 2636, 2637, 2638, 2639, 2640, 2641, 2642, 2643, 2644, 2645, 2646, 2647, 2648, 2649, 2650, 2651, 2652, 2653, 2654, 2655, 2656, 2657, 2658, 2659, 2660, 2661, 2662, 2663, 2664, 2665, 2666, 2667, 2668, 2669, 2670, 2671, 2672, 2673, 26

FINITE ELEMENTS

LAST TIME: INCOMPATIBLE MODES

REMARK: THE ASSUMPTION $f_{i\alpha}^e = 0 \quad \alpha=5,6$ APPEARS REASONABLE

SINCE WE WOULD LIKE $\sum_{\alpha=1}^6 f_{i\alpha}^e = \int_{\Omega^e} f_i d\Omega$

I.E. THE SUM OF THE NODAL FORCES SHOULD EQUAL THE TOTAL FORCE ON THE ELEMENT

BUT IF WE INCLUDE N_5, N_6 IN THE FORCE CALC,

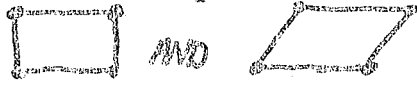
WE END UP OVERLOADING THE ELEMENT

SINCE $\sum_{\alpha=1}^6 N_{\alpha} > 1 \Rightarrow \sum_{\alpha=1}^6 f_{i\alpha}^e > \int_{\Omega^e} f_i d\Omega$

COMMENTS:

1.) EL. FUNCTIONS ARE NO LONGER CONTINUOUS.

2.) THE EL.S CONVERGE FOR LINEAR TRANSFORMATIONS ONLY

I.E.  BUT NOT 

3.) NEED ONE FURTHER MODIFICATION DUE TO TAYLOR (REF. 3.) TO OBTAIN CONVERGENCE IN ARBITRARY CONFIGURATIONS

$$B_{\alpha} = \begin{bmatrix} N_{\alpha,x} & 0 \\ 0 & N_{\alpha,y} \\ N_{\alpha,y} & N_{\alpha,x} \end{bmatrix} \quad B_5 = \frac{2}{J} \begin{bmatrix} -\xi \eta \eta & 0 \\ 0 & \xi \eta \eta \\ \xi \eta \eta & -\xi \eta \eta \end{bmatrix}$$

$$B_6 = \frac{2}{J} \begin{bmatrix} \eta \xi \xi & 0 \\ 0 & -\eta \xi \xi \\ -\eta \xi \xi & \eta \xi \xi \end{bmatrix}$$

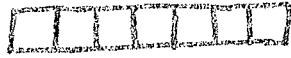
TAYLOR'S MODIFICATION

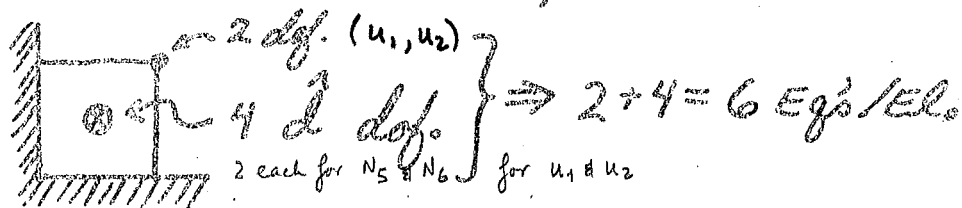
REPLACES () QUANTITIES BY THEIR VALUES AT $\xi=\eta=0$

NOTE: THIS MODIFICATION IS DONE IN SHAP20

REMARKS:

Incompatible Modes

- 1.) WORKS WELL FOR BENDING: 
WITH JUST ONE ELEMENT THROUGH THE THICKNESS.
- 2.) ITS A BIT EXPENSIVE
- 3.) THE 3-D GENERALIZATION IS OBVIOUS,
ADD $1-x^2$, $1-y^2$, $1-z^2$
- 4.) THE AXISYMMETRIC CASE DOESNT CONVERGE.
- 5.) GOOD CONSTRAINT INDEX,



CONSTRAINTS: $w^h = 1, x, y, x^2, xy, y^2$? trial fns

$0 \equiv u_{i,j} \Rightarrow 3 \text{ CONSTRAINTS at center 1 pt quad}$

$$CI = 6 - 3 = +3$$

$$CR = 6/3 = 2 \quad \left. \vphantom{\begin{matrix} CI \\ CR \end{matrix}} \right\} \text{ VERY GOOD !}$$

HOURLASS STIFFNESS :

uniform, reduced integ 1 pt.

OBSERVATIONS — 1 PT. QUAD. (U1) ON BILIN. EL.

IS 1.) OPT. FAST

2.) GIVES FAVORABLE $CI = 2 - 1 = 1$

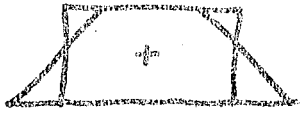
3.) EL. IS SINGULAR I.E. THERE ARE TWO ^{additional} SPURIOUS ZERO ENERGY MODES OF k^e .

k^e USUALLY HAS THREE ZERO ENERGY MODES CORRESPONDING TO

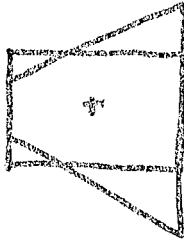
a) TWO RIGID BODY TRANSLATIONS.

b) ONE RIGID BODY ROTATION.

THE TWO ADDITIONAL MODES ARE TERMED
"HOURGLASS MODES"



$$u_1 = C \xi \eta \zeta$$



$$u_2 = C \xi \eta \zeta$$

EQUIVALENT
WITH PROPER EL.
ORIENTATION

$$0 = \underline{d}^T \underline{k}^e \underline{d} = \int_{\Omega^e} c_{ijkl} u_{i,j} u_{k,l} d\Omega$$

$$\left. \begin{array}{l} u_{i,x} = u_{i,y} = C_1 \eta \\ u_{i,y} = u_{i,x} = C_2 \xi \end{array} \right\} \Rightarrow u_{i,x}(0,0) = 0$$

KOSLOFF: ADDS TO 1 PT QUAD. STIFFNESS A STIFFNESS
ASSOCIATED WITH THESE MODES.

REMARKS:

- 1.) IT TURNS OUT THIS EL. IS EXACTLY THE INCOMPATIBLE MODES EL. IN A RECTANGULAR CONFIGURATION.
- 2.) APPEARS TO BEHAVE BETTER THAN INCOMP. MODES IN ARBITRARY CONFIGURATIONS.
- 3.) MUCH CHEAPER TO FORMULATE THAN INCOMP. MODES.
- 4.) ERROR IN KOSLOFFS PAPER
- 5.) CODED IN LERRN SUBROUTINE PLANE1 LINES 165-180
USES FORMS1 AND FORMS2

NOTE: COMPARE ELEMENT ~~FORMS1~~ ~~FORMS2~~ IN HW PROB ?
formation time

ENGINEERING CRITERIA FOR ELEMENTS

RULES OF THUMB FOR REDUCED/SELECTIVE INTEGRATION
AND INCOMPATIBLE MODES (SO CALLED NON-CONFORMING
ELEMENTS)

- 1) INVARIANCE CONDITION - { RESULTS MUST BE INDP.
OF REFERENCE FRAME.
- 2) PATCH TEST - { VERIFICATION OF COMPLETENESS

THE REDUCED/SELECTIVE INTEGRATION EL. AND INCOMPATIBLE
MODES (NON-CONFORMING) EL. BREAK THE "GALERKIN
RULES". IN FACT SO DO ALL NUMERICALLY INTEGRATED
ELEMENTS IF THE INTEGRATION RULE IS NOT EXACT.

FOR ISOPARAMETRIC ELS THE INTEGRATION RULE, NO
MATTER HOW HIGH, IS NEVER EXACT FOR EL. DOMAINS
WHICH ARE NOT $\square$ $\square$ (I.E. ONES FOR WHICH

$x = N_a x_a$, $y = N_a y_a$, $z = N_a z_a$ ARE NOT CONSTANTS) BECAUSE
CHAIN RULES PUT J 's IN THE DENOMINATOR.

THEREFORE WE NEED TO INTEGRATE RATIONAL
FUNCTIONS, NOT POLYNOMIALS.

HOW CAN WE ASCERTAIN THE VALIDITY, OR LACK OF SAME,
FOR THESE ELEMENTS?

THE PATCH TEST

<u>VERSION</u>	<u>ORIGINATOR</u>
ENGINEERING	B. IRONS '65
MATHEMATICS	G. STRANG '73 (BASED ON IRONS)

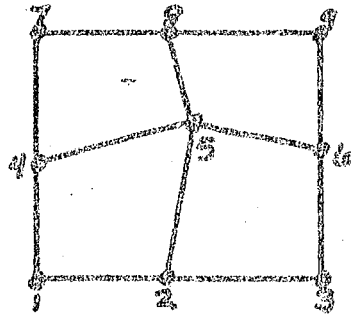
ENGINEERING VERSION OF PATCH TEST

1.) DONE ON THE COMPUTER. [∴ IT CAN BE USED TO CHECK EL. PROGRAM]

2.) IT IMPLIES THE COMPLETENESS CONDITION AND TELLS US WHATEVER "TRICKS" WE EMPLOYED TO DERIVE THIS ELEMENT (INCOMP. MODES/ QUID "TRICKS" ETC.) ARE IN FACT OK.

3.) THE PATCH TEST IS PASSED FOR AN ELEMENT IF THE DEFORMATION STATES $1, \kappa, \gamma$ (γ) ARE EXACTLY REPRESENTABLE BY AN ARBITRARY "PATCH" OF ELEMENTS. (WHENEVER THE EXACT SOLUTION BEHAVES ACCORDINGLY)

EXAMPLE:



"ARBITRARY PATCH"

SET BOUNDARY NODES ($\alpha = 1, 2, 3, 4, 7, 8, 9$) TO CORRECT VALUES
LOOK AT COMPUTED SOLUTION $\alpha = 5$

TEST #		u_α^h	u_α^h
1 {	1	1	0
	2	0	1
κ {	3	κ_α	0
	4	0	κ_α
γ {	5	γ_α	0
	6	0	γ_α

ALL $1, \kappa, \gamma$ STATES
CORRECTLY REPRESENTED
⇒ COMPLETENESS

THE SOLUTION AT $\alpha=5$ MUST BE EXACT AND SO MUST BE THE DISPLACEMENT GRADIENT (STRESSES AND STRAINS) WITHIN THE ELEMENTS.
HIGHER ORDER PATCH TEST ⇒ HIGHER ORDER COMPLETENESS
($\kappa^2, \kappa \cdot \gamma, \gamma^2$ ETC.)

LINEAR TRANSIENT ANALYSIS

PARABOLIC CASE : HEAT EQUATION

GIVEN:

$$f: \Omega \times]0, T[\rightarrow \mathbb{R}$$

$$g: \Gamma_g \times]0, T[\rightarrow \mathbb{R}$$

$$h: \Gamma_h \times]0, T[\rightarrow \mathbb{R}$$

$$u_0: \Omega \rightarrow \mathbb{R}$$

$$T \neq \infty \text{ (FINITE)}$$

(S) FIND:

$$u: \bar{\Omega} \times [0, T] \rightarrow \mathbb{R} \exists$$

$$\forall \epsilon u_{,t} + f_{L,i} = f \text{ on } \Omega \times]0, T[$$

$$\left. \begin{aligned} u &= g \text{ on } \Gamma_g \times]0, T[\\ f_{L,i} n_i &= -h \text{ on } \Gamma_h \times]0, T[\end{aligned} \right\} \text{BC's}$$

$$u(x, 0) = u_0(x) \quad x \in \Omega \quad \} \text{IC}$$

$$f_{L,i} = -K_{i,j} u_{,j}$$

GIVEN: f, g, h, u_0 AS IN (S)

FIND:

$$u: [0, T] \rightarrow S_{g,w} \exists \forall w \in V \quad (u = v + g)$$

$$(\forall \epsilon u, w) + \alpha(u, w) = (f, w) + (h, w)_\Gamma \quad (*)$$

$$(u(0) - u_0, w) = 0 \quad (**)$$

NOTE: $(\cdot, \cdot)$ $\alpha(\cdot, \cdot)$ SAME AS FOR STATIC CASE.

EXERCISE: VERIFY EQUIVALENCE (W) $\Leftrightarrow$ (S)

NOTE: (*) & USUAL HYP. $\Rightarrow$ P.D.E. + 2 BC's

(**) $\Rightarrow$ IC

$$(G_1) \left\{ \begin{array}{l} \text{GIVEN: } f, g, \bar{h}, u_0 \text{ AS IN (5)} \\ \text{FIND: } u^h: [0, T] \rightarrow S^h, \exists \forall w^h \in V^h \\ (f, u^h, w^h) + Q(u^h, w^h) = (f, w^h) + (\bar{h}, w^h)_T \\ (u^h(0) - u_0, w^h) = 0 \end{array} \right.$$

$$(G_2) \left\{ \begin{array}{l} \text{GIVEN: } f, g, \bar{h}, u_0 \text{ AS IN (5)} \\ \text{FIND: } v^h: [0, T] \rightarrow V^h \\ u^h = v^h + g^h, \quad g^h: [0, T] \rightarrow S^h \\ \exists \forall w^h \in V^h \\ (f, v^h, w^h) + Q(v^h, w^h) = (f, w^h) + (\bar{h}, w^h)_T \\ (v^h(0), w^h) = (u_0, w^h) - (g^h(0), w^h) - (f, g^h, w^h) - Q(g^h, w^h) \end{array} \right.$$

(G₁) & (G₂) ARE CALLED SEMI-DECRETE GALERKIN EQUATIONS (t IS STILL CONT.)

ASSUME:

$$v^h(x, t) = \sum_{n=1}^N N_n(x) d_n(t) \in V^h$$

$$g^h(x, t) = \sum_{n=1}^N N_n(x) g_n(t) \in S^h$$

WHERE $u^h(x, t) = v^h(x, t) + g^h(x, t)$

$$\sum_{n=1}^N N_n d_n + \sum_{n=1}^N N_n g_n$$

NOTE: THE ABOVE IS JUST A FORM OF APPROXIMATION,
NOT SEPARATION OF VARIABLES!

THE N_n 's ARE THE USUAL SHAPE FUNCTIONS, THEY HAVE NO TIME DEPENDENCE.

FIND : $\underline{d} : [0, T] \rightarrow \mathbb{R}^{N_d} \ni$

$$\underline{M} \underline{\ddot{d}} + \underline{K} \underline{d} = \underline{F}(t)$$

(MATRIX SYSTEM OF LINEAR O.D.E.'s)

$$\underline{d}(0) = \underline{d}_0$$

WHERE :

$$\underline{M} = \overset{N_d}{\underset{e=1}{\overset{N_d}{A}}} (\underline{M}^e) \dots \text{MASS OR CAPACITY MATRIX}$$

$$\underline{M}^e = [m_{pq}^e]$$

$$m_{pq}^e = \int_{\Omega^e} \rho_c N_p N_q d\Omega$$

$$\underline{K} = \overset{N_d}{\underset{e=1}{\overset{N_d}{A}}} (\underline{k}^e) \dots \text{CONDUCTIVITY MATRIX}$$

$$\underline{k}^e = [k_{pq}^e]$$

$$k_{pq}^e = \int_{\Omega^e} \underline{B}_p^T \underline{D} \underline{B}_q d\Omega$$

$$\underline{F}(t) = \overset{?}{\underset{\text{load}}{E}}(t) + \overset{N_d}{\underset{e=1}{\overset{N_d}{A}}} (\underline{f}^e(t)) \dots \text{FORCE VECTOR}$$

$$\underline{f}^e = \{f_p^e\}$$

$$f_p^e = \int_{\Omega^e} N_p f d\Omega + \int_{\Gamma_h^e} N_p h d\Gamma$$

$$- \sum_{q=1}^{N_{en}} (k_{pq}^e q_q^e + m_{pq}^e \ddot{q}_q^e)$$

AND THE CONSISTENT $\underline{d}_0$ (INITIAL VALUE) IS GIVEN BY:

$$\underline{d}_0 = \underline{\bar{M}}^{-1} \{ \overset{N_d}{\underset{e=1}{\overset{N_d}{A}}} (\underline{\bar{d}}^e) \}$$

$$\underline{\bar{M}} = \overset{N_d}{\underset{e=1}{\overset{N_d}{A}}} (\underline{\bar{M}}^e), \quad \bar{m}_{pq}^e = \int_{\Omega^e} N_p N_q d\Omega$$

$$\bar{f}_p^e = \int_{\Omega^e} u_0 N_p d\Omega - \sum_{q=1}^{N_{en}} \bar{m}_{pq}^e \ddot{q}_q^e(0)$$

comes from $\int_{\Omega} w(u(0) - u_0) d\Omega = 0$

$$\Rightarrow \int_{\Omega} N_B N_A d_0 d\Omega + \int_{\Omega} N_B \bar{B}_A \ddot{q}_A^e(0) d\Omega - \frac{1}{2} \int_{\Omega} N_B u_0 d\Omega = 0$$

$$0 = \underline{\bar{M}} \underline{d}_0 + \underline{\bar{K}} \underline{\bar{M}} \underline{\bar{d}}(0) - \int_{\Omega} N_B u_0 d\Omega$$

$$\therefore \underline{d}_0 = \underline{\bar{M}}^{-1} \{ \int_{\Omega} N_B u_0 d\Omega - \underline{\bar{M}} \underline{\bar{d}}(0) \}$$



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REMARKS:

- 1) $\underline{M}$ (AND $\bar{\underline{M}}$) IS SYMMETRIC POSITIVE DEFINITE AND HAS THE SAME BAND PROFILE AS $\underline{K}$
- 2) IT IS COMMON PRACTICE TO DIRECTLY SPECIFY $\underline{\dot{d}}$, THUS AVOIDING AN ADDITIONAL MATRIX SOLUTION.
- 3) RECALL $p = N_{nd}(q-1) + \overset{i=1}{\underset{1}{c}}$, FOR HEAT $N_{nd} = 1$ THUS
$$p = (q-1) + 1 = q$$

ONE-STEP ALGORITHMS FOR THE SEMI-DISCRETE HEAT EQUATION: GENERALIZED TRAPEZOIDAL METHOD

GIVEN: (THE FOLLOWING EQ'S DERIVED IN (M))

F.D.

$$\underline{M} \underline{\dot{d}} + \underline{K} \underline{d} = \underline{E} \quad \text{if } \underline{\dot{d}} \Rightarrow d_{n+1} = d_n + \Delta t \underline{\dot{d}}_n \quad (1)$$

$$\underline{d}(0) = \underline{D}_0 \quad (2)$$

WHERE:

$\underline{M}$... CAPACITY MATRIX; SYMM. POS. DEF.

$\underline{K}$... CONDUCTIVITY MATRIX, SYMM. POS. SEMI-DEF.

$\underline{E} = \underline{E}(t) \quad t \in [0, T]$... HEAT SUPPLY VECTOR (PREX.)

$\underline{d}$... TEMPERATURE VECTOR.

$\underline{\dot{d}}$... TIME (t) DERIVATIVE OF $\underline{d}$.

$\underline{D}_0$... GIVEN INITIAL TEMPERATURE.

APPLYING THE GENERALIZED TRAPEZOIDAL METHOD FOR FIRST ORDER SYSTEMS TO (1), WE GET

$$\underline{M} \underline{V}_{n+1} + \underline{K} \underline{d}_{n+1} = \underline{E}_{n+1} \quad (3)$$

$$\underline{d}_{n+1} = \underline{d}_n + \Delta t \underline{V}_{n+\alpha} + \underline{\dot{d}}_{\text{eqn}} \quad (4)$$

$$\underline{V}_{n+\alpha} = (1-\alpha) \underline{V}_n + \alpha \underline{V}_{n+1} \quad (5)$$

$$\underline{d}_{n+1} = \underline{d}_n + \Delta t [1-\alpha] \underline{\dot{d}}_n + \Delta t \alpha \underline{\dot{d}}_{n+1}$$

NOTATION:

$n_{ts} \dots$ NUMBER OF TIME STEPS

$\Delta t = T/n_{ts} \dots$ TIME STEP (ASSUMED CONSTANT)

$t_n = n\Delta t \dots$ TIME n

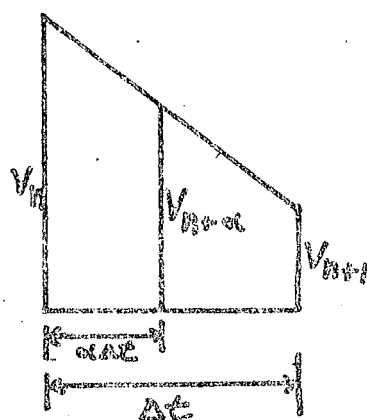
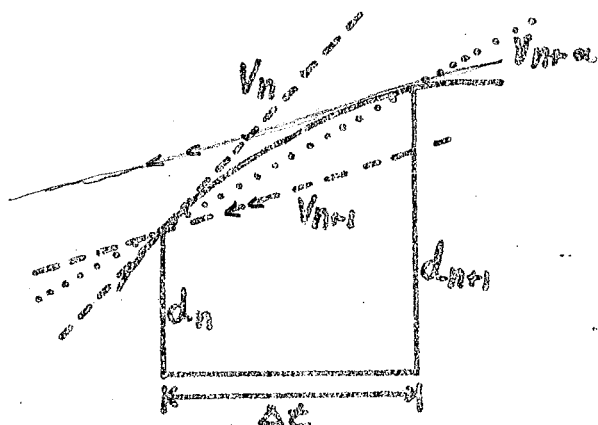
$n \dots$ INTEGER STEP NUMBER ($0 \leq n \leq n_{ts}$)

$d_n \approx d(t_n) \dots$ APPROX. TEMP.

$\dot{v}_n \approx \dot{d}(t_n) \dots$ APPROX. TIME DERIVATIVE OF TEMP.

$\alpha \in [0, 1] \dots$ PARAMETER

ILLUSTRATION OF $V_{n+\alpha}$ (1 dof)



SOME WELL KNOWN MEMBERS OF THE GENERALIZED TRAPEZOIDAL FAMILY,

α	METHOD
0	FORWARD DIFFERENCES; FORWARD EULER
1/2	TRAPEZOIDAL RULE; MIDPOINT RULE; CRANK-NICOLSON
2/3	GALERKIN
1	BACKWARD DIFFERENCES; BACKWARD EULER

REMARKS:

1) FOR $\alpha = 0$ THE METHOD IS SAID TO BE EXPLICIT.

IF $\alpha = 0$ AND M IS "LUMPED" (DIAGONAL) THEN IT CAN BE SEEN (3)-(5) THAT THE SOLUTION CAN BE ADVANCED WITHOUT THE NECESSITY OF SOLVING MATRIX EQUATIONS.

2) FOR $\alpha \neq 0$ THE METHOD IS SAID TO BE IMPLICIT.

IF $\alpha \neq 0$ A SYSTEM OF EQUATIONS WITH COEFFICIENT MATRIX $M + \alpha \Delta t B$ MUST BE SOLVED AT EACH STEP k for y_k for x_k then put into two formulas of pg 9 to get d_k

DEFINITION: CONVERGENCE

WE CALL AN ALGORITHM CONVERGENT IF FOR EACH FIXED t_n , $\Delta t = t_n / n$

$$d_n \rightarrow d(t_n) \text{ AS } \Delta t \rightarrow 0 \text{ (} n \rightarrow \infty \text{)}$$

TO BE SHOWN: STABILITY + CONSISTENCY $\Rightarrow$ CONVERGENCE

IN ORDER TO STUDY THE CHARACTERISTICS OF A TIME INTEGRATION ALGORITHM (i.e ITS STABILITY, CONSISTENCY, ACCURACY AND CONVERGENCE) WE WILL EMPLOY THE MODAL APPROACH (ALSO CALLED SPECTRAL OR FOURIER ANALYSIS)

MODAL ANALYSIS: REDUCTION TO A SINGLE DEGREE-OF-FREEDOM (SDOF) MODEL EQUATION

WE WILL NOW DECOMPOSE (1) INTO n uncoupled SCALAR EQUATIONS. AND THUS THE ANALYSIS OF THE ALGORITHM WILL REDUCE TO THE ANALYSIS OF A SIMPLE SINGLE-DEGREE-OF-FREEDOM CASE.

FOR Transient Analysis

$$\text{let } \underline{d} = \underline{\psi} e^{-\lambda t} \quad \lambda \geq 0$$

$$\underline{\dot{d}} = -\lambda \underline{\psi} e^{-\lambda t}$$

$$\therefore \underline{M} \underline{\dot{d}} + \underline{K} \underline{d} = (-\lambda \underline{M} + \underline{K}) \underline{\psi} e^{-\lambda t}$$

look at homogeneous system, first

$$\therefore \underline{\psi} e^{-\lambda t} \neq 0 \Rightarrow (\underline{K} - \lambda \underline{M}) \underline{\psi} = 0$$

EIGENVALUE PROBLEM:

GIVEN: $\underline{M} \dots$ SYMM. POS. DEF.

$\underline{K} \dots$ SYMM. POS. SEMI-DEF.

FIND: $\lambda_l, \underline{\psi}_l \in \mathbb{R}^{n_{eq}} \quad 1 \leq l \leq n_{eq} \quad \exists$

$$(\underline{K} - \lambda_l \underline{M}) \underline{\psi}_l = 0 \quad l \in \{1, 2, \dots, n_{eq}\} \quad (6)$$

WHERE $0 \leq \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_{n_{eq}}$ (7)

AND $\underline{\psi}_l^T \underline{M} \underline{\psi}_m = \delta_{lm} \dots$ ORTHONORMALITY (8)

RECALL THAT $\{\underline{\psi}_l\}_1^{n_{eq}}$ FORMS A BASIS FOR $\mathbb{R}^{n_{eq}}$ AND THAT $\delta_{lm} = \begin{cases} 0 & l \neq m \\ 1 & l = m \end{cases}$ (KRONCKER DELTA)

FROM (8) (ORTHONORMALITY)

$$\underline{\psi}_l^T \underline{K} \underline{\psi}_m = \lambda_l \delta_{lm} \quad (\text{NO SUM}) \quad (9)$$

USING THE ABOVE WE WILL DECOMPOSE FIRST THE ORIGINAL SYSTEM OF O.D.E.S (1), AND SECOND THE GENERALIZED TRAPEZOIDAL ALGORITHM (3)-(5)

FIRST, LET

$$\underline{d}(t) = \sum_{m=1}^{n_{eq}} d_{lm}(t) \underline{\psi}_m \quad \text{now } \underline{\psi}_l^T \underline{M} \underline{d}(t) = \sum_{m=1}^{n_{eq}} d_{lm}(t) \delta_{lm} = d_l(t) \quad (10)$$

THUS

$$\underline{\dot{d}}(t) = \sum_{m=1}^{n_{eq}} \dot{d}_{lm}(t) \underline{\psi}_m \quad \text{and } \underline{\psi}_l^T \underline{M} \underline{\dot{d}} = \dot{d}_l(t) \quad \text{similarly} \quad (11)$$

WHERE $d_{lm}(t)$ ARE CALLED FOURIER COEFFICIENTS,

$$d_{lm}(t) = \underline{\psi}_l^T \underline{M} \underline{d}(t) \quad (12)$$

$$\dot{d}_{lm}(t) = \underline{\psi}_l^T \underline{M} \underline{\dot{d}}(t) \quad (13)$$

SUBSTITUTION OF (10) & (11) INTO (1) YIELDS

$$\sum_{m=1}^{N_{eq}} (\dot{d}_{(m)} \psi_m^T M \psi_m) + \underbrace{d_{(m)} \psi_m^T K \psi_m}_{\lambda_{(m)} d_{(m)}} = \psi_m^T E \quad (14)$$

$$\text{THUS } \dot{d}_{(1)} + \lambda_1 d_{(1)} = F_{(1)} \quad (15)$$

$$\text{WHERE } F_{(1)}(t) = \psi_1^T E(t)$$

THE INITIAL CONDITION REDUCES TO

$$\begin{aligned} d_{(1)}(0) &= \psi_1^T M d(0) \\ &= \psi_1^T M D_0 \\ &\stackrel{\text{def.}}{=} D_{(1)} \end{aligned} \quad (16)$$

TO SIMPLIFY SUBSEQUENT WRITING DROP 1 SUBSCRIPT.
THUS OUR TYPICAL MODAL INITIAL-VALUE PROBLEM
BECOMES

$$\dot{d} + \lambda d = F, \quad t \in [0, T] \quad (17)$$

$$d(0) = D_0 \quad (18)$$

THIS IS SOMETIMES CALLED THE SINGLE-DEGREE-
OF-FREEDOM MODEL PROBLEM.

A SIMILAR DECOMPOSITION OF THE GENERALIZED
TRAPAZOIDAL METHOD (3)-(5) YIELDS

$$d_n = \sum_{m=1}^{N_{eq}} d_{n(m)} \psi_m \quad (19)$$

$$d_{n+1} = \sum_{m=1}^{N_{eq}} d_{n+1(m)} \psi_m \quad (20)$$

$$\text{WHERE } d_n(t) = \psi_1^T M d_n \quad (21)$$

$$d_{n+1}(t) = \psi_1^T M d_{n+1} \quad (22)$$

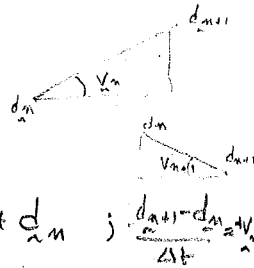


let $\dot{\underline{d}} = \underline{v}$

$$\underline{M} \underline{v} + \underline{K} \underline{d} = \underline{F}$$

then $\underline{M} \underline{v}_{n+1} + \underline{K} \underline{d}_{n+1} = \underline{F}_{n+1}$

for $\underline{d}_{n+1} = \underline{d}_n + \Delta t \underline{v}_n$ forward diff $\alpha=0$ $\frac{\underline{d}_{n+1} - \underline{d}_n}{\Delta t} = \underline{v}_n$



$$\underline{v}_{n+\alpha} = (1-\alpha) \underline{v}_n + \alpha \underline{v}_{n+1} \Rightarrow \underline{v}_n = 1 \cdot \underline{v}_n + 0 \cdot \underline{v}_{n+1}$$

for $\underline{d}_{n+1} = \underline{d}_{n+1} - \Delta t \underline{v}_{n+1}$ backward diff $\alpha=1 \Rightarrow \underline{d}_{n+1} = \Delta t \underline{v}_{n+1} + \underline{d}_n$; $\frac{\underline{d}_{n+1} - \underline{d}_n}{\Delta t} = \underline{v}_{n+1}$

$$\underline{v}_{n+\alpha} = (1-\alpha) \underline{v}_n + \alpha \underline{v}_{n+1} \Rightarrow \underline{v}_{n+1} = 0 \cdot \underline{v}_n + 1 \cdot \underline{v}_{n+1}$$

for $\underline{v}_{n+1/2} = \frac{\underline{v}_{n+1} + \underline{v}_n}{2} = \alpha \underline{v}_{n+1} + (1-\alpha) \underline{v}_n$ centered diff

$$\underline{d}_{n+1} = \underline{d}_n + \Delta t \underline{v}_{n+1/2} \quad \underline{v}_{n+1/2} = \frac{\underline{d}_{n+1} - \underline{d}_n}{2 \cdot \Delta t/2}$$

proof $\left. \begin{aligned} \underline{d}_{n+1} &= \underline{d}_{n+1/2} + \frac{\Delta t}{2} \underline{v}_{n+1/2} \\ \underline{d}_n &= \underline{d}_{n+1/2} - \frac{\Delta t}{2} \underline{v}_{n+1/2} \end{aligned} \right\} \underline{d}_{n+1} - \underline{d}_n = \Delta t \underline{v}_{n+1/2}$

thus $\underline{d}_{n+1} = \underline{d}_n + \Delta t \underline{v}_{n+\alpha} = \underline{d}_n + \Delta t \{ (1-\alpha) \underline{v}_n + \alpha \underline{v}_{n+1} \}$

if $\alpha=0$ explicit & $\underline{M}$ is diag then from the above $\underline{d}_{n+1} = \underline{d}_n + \Delta t \underline{v}_n$

$$\underline{M} \underline{v}_{n+1} = \underline{F}_{n+1} - \underline{K} (\underline{d}_n + \Delta t \underline{v}_n) = \text{known rhs} \quad \therefore (\underline{v}_i)_{n+1} = (\underline{F}_{n+1} - \underline{K} \underline{d}_{n+1})_i / [\underline{M}]_{ii}$$

if $\alpha \neq 0$ then since $\underline{K}$ is not diag then $\{ \underline{M} + \alpha \Delta t \underline{K} \} \underline{v}_{n+1} = \underline{F}_{n+1} - \underline{K} \{ \underline{d}_n + \Delta t (1-\alpha) \underline{v}_n \}$
 This is implicit not diag known rhs

solve for $\underline{v}_{n+1}$; put into $\underline{v}_{n+\alpha} = (1-\alpha) \underline{v}_n + \alpha \underline{v}_{n+1}$ to get $\underline{v}_{n+\alpha}$; put into $\underline{d}_{n+1} = \underline{d}_n + \Delta t \underline{v}_{n+\alpha}$ to get $\underline{d}_{n+1}$

Now $\alpha \underline{M} \underline{v}_{n+1} + \alpha \underline{K} \underline{d}_{n+1} = \alpha \underline{F}_{n+1}$

$$(1-\alpha) \underline{M} \underline{v}_n + \underline{K} (1-\alpha) \underline{d}_n = (1-\alpha) \underline{F}_n$$

$$\underline{M} \left\{ \frac{\underline{d}_{n+1} - \underline{d}_n}{\Delta t} \right\} + \underline{K} \{ (1-\alpha) \underline{d}_n + \alpha \underline{d}_{n+1} \} = \alpha \underline{F}_{n+1} + (1-\alpha) \underline{F}_n$$

$$\therefore (\underline{M} + \alpha \underline{K} \Delta t) \underline{d}_{n+1} = \underline{M} \underline{d}_n - \Delta t \underline{K} (1-\alpha) \underline{d}_n + \Delta t \underline{F}_{n+\alpha}$$

$$\underline{F}_{n+\alpha} = \alpha \underline{F}_{n+1} + (1-\alpha) \underline{F}_n$$

Now let $\underline{\psi}_L^T$ be applied to both sides of eqn w/ $\underline{d} = \underline{d}_m^T \underline{\psi}_m$ & $\underline{\psi}_L^T \underline{M} \underline{\psi}_m = \delta_{Lm}$ $\underline{\psi}_L^T \underline{K} \underline{\psi}_m = \lambda_L$

then $[1 + \alpha \lambda_L \Delta t] \underline{d}_{L,n+1} = [1 - \Delta t \lambda_L (1-\alpha)] \underline{d}_{L,n} + \Delta t \underline{\psi}_L^T \underline{F}_{n+\alpha}$

w/ $\underline{d}(0) = \underline{\psi}_L^T \underline{M} \underline{d}(0)$

Let T be a bounded linear operator on a Hilbert space H . Then T is called a *Hilbert-Schmidt operator* if

$$\|T\|_2^2 = \sum_{n=1}^{\infty} \|Te_n\|^2 < \infty$$

$$\|T\|_2 = \sqrt{\|T\|_2^2}$$

$$\sum_{m=1}^{n_{eq}} [d_{n+1}(m) \psi_2^T (M + \alpha \Delta t K) \psi_m - d_{n+1}(m) \psi_2^T (M - (1-\alpha) \Delta t K) \psi_m] = \Delta t \psi_2^T E_{n+\alpha} \quad (23)$$

$$(1 + \alpha \Delta t \lambda_2) d_{n+1}(0) = (1 - (1-\alpha) \Delta t \lambda_2) d_n(0) + \Delta t F_{n+\alpha} \quad (24)$$

$$d_0(0) = D_0(0) \quad (25)$$

WHERE $E_{n+\alpha} = (1-\alpha) E_n + \alpha F_{n+1}$

$$F_{n+\alpha}(0) = \psi_2^T E_{n+\alpha}$$

DROP THE l^{TH} MODAL SUBSCRIPT

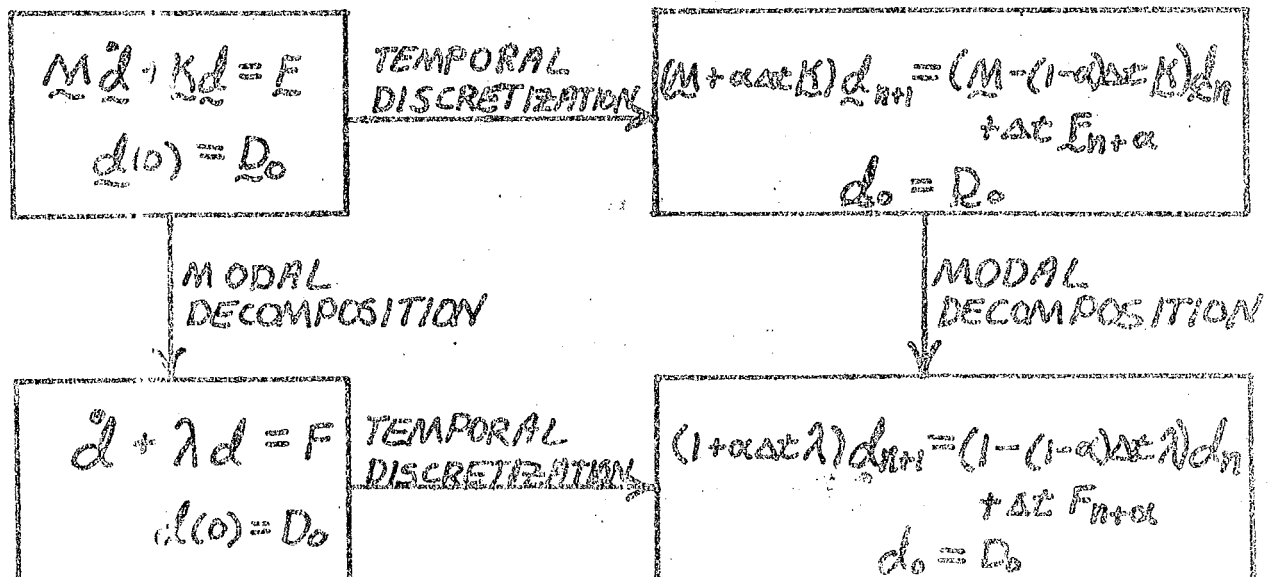
$$(1 + \alpha \Delta t \lambda) d_{n+1} = (1 - (1-\alpha) \Delta t \lambda) d_n + \Delta t F_{n+\alpha} \quad (26)$$

$$d_0 = D_0 \quad (27)$$

THIS IS SOMETIMES CALLED THE TEMPORALLY DISCRETIZED SDOF MODEL PROBLEM.

REMARKS:

- 1.) NOTE THAT DIRECTLY APPLYING THE GEN. TRAP METHOD TO (17) ALSO RESULTS IN (24). THIS FACT IS DEPICTED IN THE COMMUTATIVE DIAGRAM BELOW.



2.) THE CONVERGENCE OF d_n TO $d(t_n)$ CAN BE ESTABLISHED BY SHOWING THE FOURIER COEFFICIENTS CONVERGE (i.e. $d_{n(l)}$ CONVERGES TO $d_{(l)}(t_n)$ FOR EACH $l \in \{1, 2, \dots, \text{Neg}\}$).

LET $e(t_n) = d_n - d(t_n) \dots$ ERROR IN d_n

$e_{(l)}(t_n) = d_{n(l)} - d_{(l)}(t_n) \dots l^{\text{th}}$ FOURIER COEFF.

USING M (SYMM. POS. DEF.) TO DEFINE A NORM.

$$\begin{aligned} e(t_n)^T M e(t_n) &= \sum_{l,m=1}^{\text{Neg}} \{e_{(l)}(t_n) \psi_l\}^T M \{e_{(m)}(t_n) \psi_m\} \\ e(t_n) &= \sum_{l=1}^{\text{Neg}} \{e_{(l)}(t_n) \psi_l\} \\ &= \sum_{l,m=1}^{\text{Neg}} e_{(l)}(t_n) e_{(m)}(t_n) \psi_l^T M \psi_m \\ &= \sum_{l,m=1}^{\text{Neg}} e_{(l)}(t_n) e_{(m)}(t_n) \delta_{lm} \\ &= \sum_{l=1}^{\text{Neg}} (e_{(l)}(t_n))^2 \end{aligned}$$

SO $e(t_n)^T M e(t_n) \rightarrow 0$ IF AND ONLY IF $e_{(l)}(t_n) \rightarrow 0$

FOR EACH $l \in \{1, 2, \dots, \text{Neg}\}$. SINCE M IS POSITIVE

DEFINITE $e(t_n)^T M e(t_n) \rightarrow 0$ IF AND ONLY IF

$e(t_n) \rightarrow 0$.

OR $e_{n(l)} \rightarrow 0 \Leftrightarrow e_{(l)}(t_n) \rightarrow 0 \Leftrightarrow d_n \rightarrow d(t_n) \Rightarrow$ CONVERGENCE

THUS

CONVERGENCE OF SDOF PROB. $\Leftrightarrow$ CONVERGENCE OF SYSTEM

Hand it in (HW) a week from exam.

In a dynamic system the nodal values become fns of time. However we are not assuming separation of variables.

Most common is trapezoidal method. Solve problem by working w/ the EV, $\vec{EV}$ to help decouple time/space coord.

① now we get $\underline{M} \dot{\underline{d}} + \underline{K} \underline{d} = \underline{F}$ w/ $\underline{d}(0) = \underline{D}_0$ from discretiz of FE problem

② Under temporal discretiz we get $(\underline{M} + \alpha \Delta t \underline{K}) \underline{d}_{n+1} = (\underline{M} - (1-\alpha) \Delta t \underline{K}) \underline{d}_n + \Delta t \underline{F}_{n+\alpha}$
where $\underline{F}_{n+\alpha} = (1-\alpha) \underline{F}_{n+1} + \alpha \underline{F}_n$

③ Step n+1 based on results of step n

④ Using modal decomp on the original problem.

$$\dot{d} + \lambda d = F \quad \forall \lambda \text{ of the system.}$$

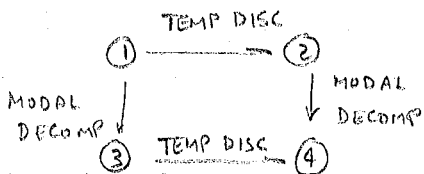
④ TEMP DISC

$$d_0 = D_0$$

If we use generalized Trapez algorithm of the temporal discretiz we get

$$(1 + \alpha \Delta t \lambda) d_{n+1} = (1 - (1-\alpha) \Delta t \lambda) d_n + \Delta t F_{n+\alpha}$$

w/ $d_0 = D_0$



non linear problems will not necessarily satisfy this diagram.

if $1 \rightarrow 2 \rightarrow 4 \neq 1 \rightarrow 3 \rightarrow 4$ then numerical problem may not converge to initial problem.

Convergence let $\underline{d}_n$ be soln to matrix problem & $\underline{d}(t)$ be soln to actual problem.
and define $\underline{e}(t_n) = \underline{d}_n - \underline{d}(t_n)$ and look at temporal convergence.
We had shown that static problems are convergent, we will now look at discretized time and see whether the "quasi static" results are now convergent in time

now $\underline{e}(t_n) \rightarrow 0$ if $\underline{e}(e)(t_n) = \underline{d}_n(e) - \underline{d}(e)(t_n) \rightarrow 0$

Proof: Recall $\underline{M}$ is $\underline{M}^T$ & $\underline{e}^T \underline{M} \underline{e} > 0$ for $\underline{e} \neq 0$

Now take $\underline{e}_n^T \underline{M} \underline{e}(t_n) = \sum_{n=1}^{n_{eq}} (\underline{e}(t_n) \underline{\psi}_{(n)})^T \underline{M} (\underline{e}_{(m)}(t_n) \underline{\psi}_{(n)})$

$$= \sum_{l,n=1}^{n_{eq}} e_l(t_n) e_m(t_n) \psi_l^T M \psi_m = \sum_{l,m=1}^{n_{eq}} e_l e_m \delta_{lm} = \sum_{l=1}^{n_{eq}} e_l^2 > 0$$

$$\text{if } \sum e_l^2 = 0 \Rightarrow e_l = 0 \therefore \underline{e} = 0 \quad \& \quad \underline{d}_m = \underline{d}(t_n) \quad QED$$

Stability of algorithm + consistency $\Rightarrow$ convergence.

Consistency is equivalent to accuracy. Stability will now be looked at

Look at homogeneous system $\dot{d} + \lambda d = 0 \Rightarrow d = C e^{-\lambda t}$ w/ $d = d(t_n)$ when $t = t_n$
 $d(t_n) = C e^{-\lambda t_n} \quad C = d(t_n) e^{\lambda t_n} \therefore d = d(t_n) e^{-\lambda(t-t_n)}$

$$\text{EXACT sol.} \quad d(t_{n+1}) = \exp \left\{ -\lambda \frac{(t_{n+1} - t_n)}{\Delta t} \right\} d(t_n) \quad \text{where } \lambda \geq 0.$$

$$\text{if } \Delta t > 0 \quad |d(t_{n+1})| < |d(t_n)| \quad \text{if } \lambda > 0$$

$$\text{if } \lambda = 0 \quad d(t_{n+1}) = |d(t_n)| \quad \text{since eqn becomes } \dot{d} = 0 \quad d = \text{const.}$$

look at the homogeneous version of time discrete results

$$(1 + \alpha \Delta t \lambda) d_{n+1} = (1 - (1 - \alpha) \Delta t \lambda) d_n$$

$$\lambda \geq 0 \quad \Delta t > 0 \quad \alpha \in [0, 1] \quad \therefore \quad 1 + \alpha \Delta t \lambda > 0$$

$$\therefore d_{n+1} = A d_n \quad A = \frac{(1 - (1 - \alpha) \Delta t \lambda)}{1 + \alpha \Delta t \lambda} \text{ is the amplif factor}$$

A approx $\exp(-\lambda \Delta t)$

$$\text{we now ask that } |d_{n+1}| < |d_n| \quad \text{if } \lambda > 0 \quad \text{as analog of contin. syst.}$$

$$|d_{n+1}| = |d_n| \quad \lambda = 0$$

$$\text{if } \lambda = 0 \quad A = 1 \quad \therefore d_{n+1} = d_n$$

$$\text{if } \lambda > 0 \quad |A| < 1 \quad \text{but } A = 1 - \frac{\Delta t \lambda}{1 + \alpha \Delta t \lambda} \quad \therefore \quad \frac{\Delta t \lambda}{1 + \alpha \Delta t \lambda} > 0$$

$$-1 < A < 1 \quad \text{leads to} \quad 1 > \frac{\Delta t \lambda}{2(1 + \alpha \Delta t \lambda)} > 0$$

$$\therefore 2 + 2\alpha \Delta t \lambda > \Delta t \lambda$$

$$2 > \Delta t \lambda (1 - 2\alpha)$$

$$\frac{2}{\Delta t \lambda} > (1 - 2\alpha)$$

$$\therefore \frac{1}{2} \left(1 - \frac{2}{\Delta t \lambda} \right) < \alpha$$

$$\boxed{\text{for stability } \alpha \geq \frac{1}{2} \quad \text{if } \Delta t, \lambda > 0}$$

$$\text{if } \alpha < \frac{1}{2} \quad \Delta t < \frac{2}{\lambda(1 - 2\alpha)}$$

Δt depends on largest λ of system i.e. λ_{max} .

Unconditional stability NO Δt limit (EG $\alpha \geq 1/2$)
 Conditional stability $\exists!$ a Δt restrict EG $\alpha < 1/2$

which is $\Delta t < \frac{2}{\lambda(1-2\alpha)}$

if $d_{n+1} = A d_n$ $d_n = A d_{n-1}$ $\therefore d_{n+1} = A^2 d_{n-1} = \dots A^{n+1} d_0$

if $|A| < 1 \Rightarrow A^n \rightarrow 0$; if $|A| > 1$ unstable.

where A^n grows

A^n	100	1000
.99	.37	4.32×10^{-5}
1.01	2.7	2.09×10^4
.9	2.66×10^{-5}	1.75×10^{-46}
1.1	1.39×10^{11}	2.47×10^{41}

5/7/81

Midterm changed to Thursday 14 May

To summarize - stability for general trapezoidal methods is dependent on the amplification factor

$$A = \frac{1 - (1-\alpha)\lambda\Delta t}{1 + \alpha\lambda\Delta t} \quad \text{for } d_{n+1} = A d_n$$

The system is stable when $|A| < 1$ and $\lambda = \lambda_{\text{neg}}$

it is unconditionally stable for $\alpha \geq 1/2$

" " conditionally " " $\alpha < 1/2$ ie $\Delta t < 2 / [(1-2\alpha)\lambda_{\text{neg}}]$

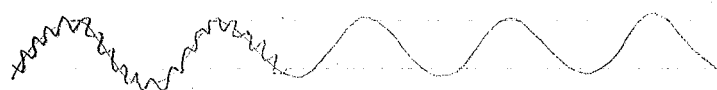
See handout:

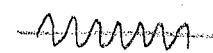
1. if $\alpha=1$ and let $\Delta t \rightarrow \infty$ this is backward difference

$$K d_{n+1} = F_{n+1} \quad (\text{ie we tend to equilib})$$

2. if $\lambda\Delta t \gg 1$ for $\alpha=1/2$ (trapezoid) high mode $A_{\infty} = -1$ since $\frac{1 - 1/2\lambda\Delta t}{1 + 1/2\lambda\Delta t} \approx -1$
 $\therefore d_{n+1} = -d_n \quad \therefore d_n = (-1)^n d_0$

non-suppressed due to this
 flip-flop of the high mode



to filter the oscill out report only the average $\frac{d_{n+1} + d_n}{2} = \frac{(-1)^{n+1} + (-1)^n}{2} = 0$
 this technique removes high mode problems.
 the midpoint defines the oscillation limits 

Consistency and convergence

we had looked at homog system. Now must include load

$$\therefore d_{n+1} = A d_n + L_n \quad (1)$$

$$L_n = F_{n+\alpha} \Delta t / (1 + \alpha \Delta t \lambda)$$

this is the algorithm

$$\text{the exact can be represented by } d(t_{n+1}) = A d(t_n) + L_n + \tau(t_n) \quad (2)$$

where we are fitting the ^{exact} algorithm to the algorithm and looking at the difference between (1) & (2), which defines the truncation error.

$$\text{Now if } |\tau(t)| \leq C \Delta t^{k+1} \quad t \in [0, T] \quad \text{with } C \neq C(\Delta t) \quad k \geq 0$$

then the algorithm is said to be consistent and k is defined as the order of accuracy of algorithm, i.e. rate of convergence.

$$\left(\begin{array}{l} \text{Facts: for generalized trapez family} \\ \text{for } d + \lambda d = F \end{array} \right. \quad \begin{array}{ll} k=1 & \alpha \in [0, 1] \\ k=2 & \alpha = 1/2 \text{ (trapez rule)} \end{array} \left. \right)$$

To show this take

$$\left. \begin{array}{l} d(t_{n+1}) \\ d(t_n) \\ F(t_{n+1}) \end{array} \right\} \begin{array}{l} \text{expand about } t_{n+\alpha} \\ \text{using Taylor exp} \end{array}$$

and use exact i.e. $(d + \lambda d = F)$ at $t_{n+\alpha}$ also use time deriv of it.
eqn.

EXERCISE do this (might be on exam).

Theorem

Assume t_n is fixed ($n, \Delta t$ vary). i.e. $n \rightarrow \infty \quad \Delta t \rightarrow 0$.

Consider (1) and (2) and assume (i) stability $|A| \leq 1$

$$(ii) \text{ consistency } |\tau(t)| \leq C \Delta t^{k+1} \quad k \geq 0$$

then $e(t_n) = d_n - d(t_n) \rightarrow 0$ as $\Delta t \rightarrow 0$.

Subtract (2) from (1) then

$$e(t_{n+1}) = A e(t_n) - \tau(t_n) \quad \text{this is the error eqn.}$$

$$\text{now } e(t_n) = A e(t_{n-1}) - \tau(t_{n-1})$$

$$\therefore e(t_{n+1}) = A^2 e(t_{n-1}) - A e(t_{n-1}) - \tau(t_n)$$

$$e(t_{n+1}) = A^3 e(t_{n-2}) - A^2 e(t_{n-2}) - A e(t_{n-1}) - \tau(t_n)$$

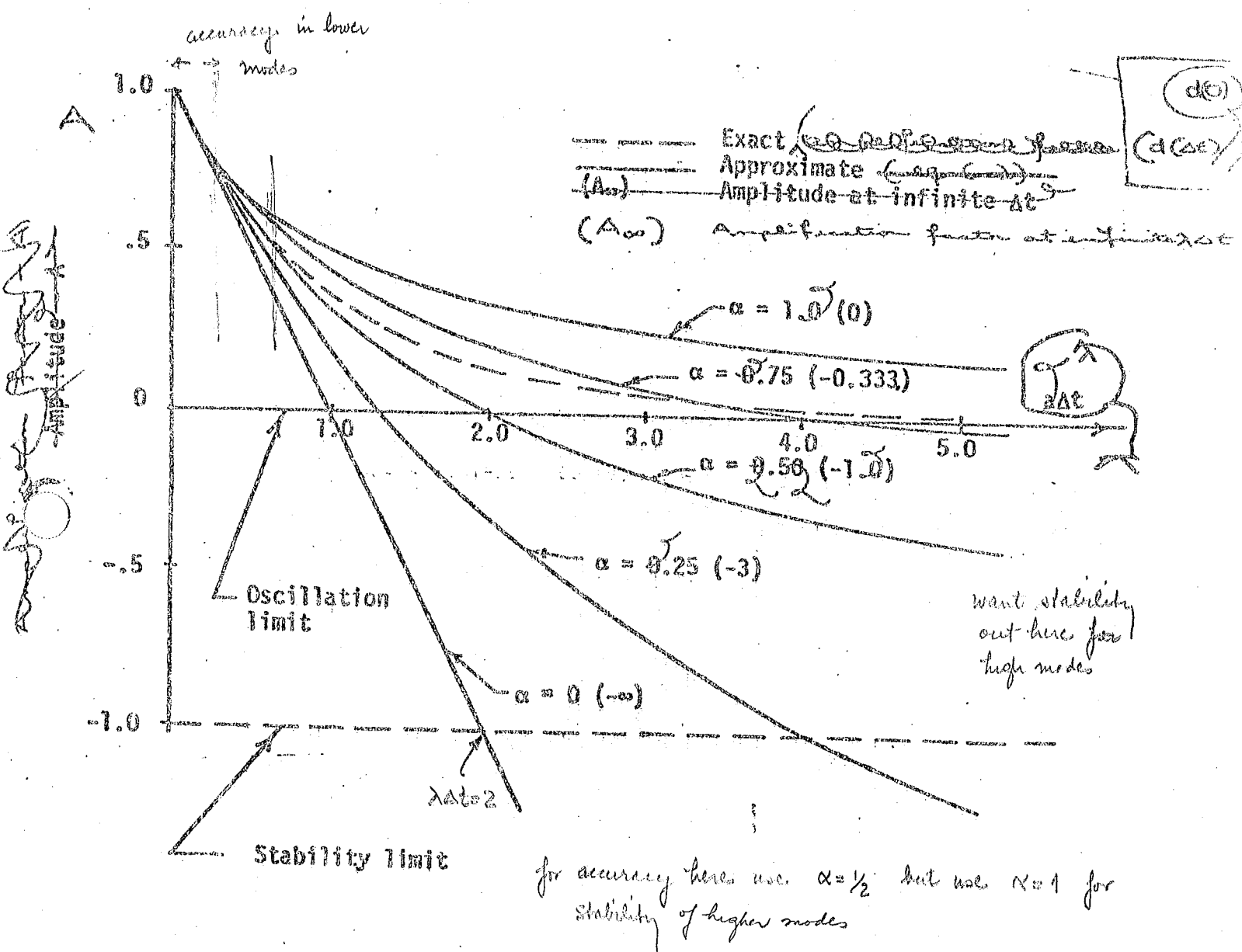


Fig. 1 Integration Formula Characteristics

$$e(t_{n+1}) = A^{n+1} e(0) - \sum_{i=0}^n A^i \tau(t_{n-i})$$

$$\text{but } e(0) = d_0 - d(0) = 0$$

$$\therefore e(t_n) = - \sum_{i=0}^{n-1} A^i \tau(t_{n-i}) \Rightarrow |e(t_n)| = \left| \sum \right| \leq \sum_{i=0}^{n-1} |A|^i \max_{t \in [0, t]} |\tau(t)|$$

using stability $|A| < 1 \Rightarrow$

$$|e(t_n)| \leq \sum_{i=0}^{n-1} \max_{t \in [0, T]} |\tau(t)| \leq n c \Delta t^{k+1} \stackrel{\text{use consistency}}{\leq} t_n \cdot c \Delta t^k \stackrel{\text{fixed}}{=} O(\Delta t^k)$$

$$\text{as } \Delta t \rightarrow 0 \quad |e(t_n)| \rightarrow 0$$

Since we've shown this for a single mode, & we said that the discrete system is a superpos. of different modes, then $\Rightarrow$ the discrete system converges to the actual problem if $\Delta t \rightarrow 0$.

Let us look at the hypothetical case in elastodynamics. This is an IBVP.

$$\begin{aligned} f_i &: \Omega \times]0, T[\rightarrow \mathbb{R} \quad \text{friction} \quad p: \Omega \rightarrow \mathbb{R} \quad (\text{mass density}) \\ (s) \quad \begin{cases} g_i: \Gamma_{g_i} \times]0, T[\rightarrow \mathbb{R} \\ h_i: \Gamma_{h_i} \times (0, T) \rightarrow \mathbb{R} \end{cases} & \left. \vphantom{\begin{matrix} g_i \\ h_i \end{matrix}} \right\} \text{BC's} \end{aligned}$$

$$\text{Given IC: } \begin{aligned} u_{0,i} &: \Omega \rightarrow \mathbb{R} \quad (\text{displ}) \\ \dot{u}_{0,i} &: \Omega \rightarrow \mathbb{R} \quad (\text{velocity}) \end{aligned}$$

we want to find $\Rightarrow u_i: \bar{\Omega} \times [0, T] \rightarrow \mathbb{R} \quad \exists.$

$$\rho u_{i,tt} = \sigma_{ij,j} + f_i \quad \text{on } \Omega \times (0, T)$$

$$\text{and } \begin{aligned} u_i &= g_i \quad \text{on } \Gamma_{g_i} \times (0, T) \\ \sigma_{ij} n_j &= h_i \quad \text{on } \Gamma_{h_i} \times (0, T) \end{aligned}$$

domain is fixed w/ time
but data may vary

$$\left. \begin{aligned} u_i(\underline{x}, 0) &= u_{0,i}(\underline{x}) \\ u_{i,t}(\underline{x}, 0) &= \dot{u}_{0,i}(\underline{x}) \end{aligned} \right\} \underline{x} \in \Omega$$

$$\sigma_{ij} = C_{ijkl} u_{(k,l)} \quad \text{where } C_{ijkl} \neq C_{ijkl}(t)$$

(w) Given $u_i: [0, T] \rightarrow \mathcal{A}_i \ni \forall w_i \in V_i$ (no time dependence here)

Then

$$(\underline{w}, \rho \ddot{\underline{u}}) + a(\underline{w}, \underline{u}) = (\underline{w}, \underline{f}) + (\underline{w}, \underline{h})_{\Gamma_1} \quad \left] \text{ gives PDE + BC} \right.$$

$\searrow \int_{\Omega} w_i \rho \ddot{u}_i d\Omega$

also $\left. \begin{aligned} (\underline{w}, \underline{u}(0) - \underline{u}_0) &= 0 \\ (\underline{w}, \dot{\underline{u}}(0) - \dot{\underline{u}}_0) &= 0 \end{aligned} \right\} \text{ IC in weak form}$

We can establish under appropriate hypothesis $(S) \Leftrightarrow (w)$.

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Review of last week

We were looking at the elastodynamics problem in its $(S) \Leftrightarrow (w)$ forms. we will write the (G1) form next

Given: $f_i, g_i, h_i, (u_{0i}, \dot{u}_{0i} \in \mathbb{R})$

Find:

$$u_i^h: [0, T] \rightarrow \mathcal{A}_i^h \ni \forall w_i^h \in U_i^h$$

$\mathcal{C}$ with g -type bc included

$U_i^h, \mathcal{A}_i^h$ same as static problem

$$(\underline{w}^h, \rho \ddot{\underline{u}}^h) + a(\underline{w}^h, \underline{u}^h) = (\underline{w}^h, \underline{f}) + (\underline{w}^h, \underline{h})_{\Gamma_1}$$

with IC $\left. \begin{aligned} (\underline{w}^h, \underline{u}^h(0) - \underline{u}_0^h) &= 0 \\ (\underline{w}^h, \dot{\underline{u}}^h(0) - \dot{\underline{u}}_0^h) &= 0 \end{aligned} \right\}$

Now we separate into unknown = known form

$$\text{let } v_i^h(\underline{x}, t) = \sum_{\eta \in \eta_i} N_a(\underline{x}) d_{iA}(t) \in U_i^h$$

$$(G2) \quad g_i^h(\underline{x}, t) = \sum_{\eta \in \eta_i} N_a(\underline{x}) g_{iA}(t) \in \mathcal{A}_i^h$$

Then we get if $\underline{u}^h = \underline{v}^h + \underline{g}^h$

$$(\underline{w}^h, \rho \ddot{\underline{u}}^h) + a(\underline{w}^h, \underline{u}^h) = (\underline{w}^h, \underline{f}) + (\underline{w}^h, \underline{h})_{\Gamma_1} - (\underline{w}^h, \rho \ddot{\underline{g}}^h) - a(\underline{w}^h, \underline{g}^h)$$

and

$$\begin{aligned} (\underline{w}^h, \underline{u}^h(0)) &= (\underline{w}^h, \underline{u}_0^h - \underline{g}(0)) \\ (\underline{w}^h, \dot{\underline{u}}^h(0)) &= (\underline{w}^h, \dot{\underline{u}}_0^h - \dot{\underline{g}}(0)) \end{aligned}$$

The matrix formulation then becomes:

Find: $\underline{d}: [0, T] \rightarrow \mathbb{R}^{n_{eq}}$ s.t.

$$\underline{M} \ddot{\underline{d}} + \underline{K} \underline{d} = \underline{F}(t)$$

where $\underline{K}$ is same as before for elastostatic

where $\underline{M} = \underset{e=1}{\overset{n_{el}}{A}} (\underline{m}^e)$

$$\underline{m}^e = [m_{pq}^e]$$

$$p = n_{ed}(a-1) + i$$

$$q = n_{ed}(b-1) + j$$

$$n_{ed} = n_{sd}$$

$$[m_{pq}^e] = \int_{\Omega^e} \rho N_a N_b \delta_{ij} d\Omega$$

now $\underline{f}^e = \{f_p^e\}$; $f_p^e = \int_{\Omega^e} N_a f_i(t) d\Omega + \int_{\Gamma^e} N_a h_i d\Gamma$

$$= \sum_{q=1}^{n_{ee}} (k_{pq}^e g_q^e + m_{pq}^e \ddot{g}_q^e)$$

Now the IC's are that

$$\underline{d}(0) = \underline{d}_0; \quad \dot{\underline{d}}(0) = \dot{\underline{d}}_0$$

the discretizing is for the IC's

$$\underline{d}_0 = \underline{\bar{M}}^{-1} \left\{ \underset{e=1}{\overset{n_{el}}{A}} (\bar{\underline{d}}^e) \right\}$$

$$\underline{\bar{M}} = A(\bar{\underline{m}}^e)$$

$$\bar{m}_{pq}^e = \int_{\Omega^e} N_a N_b \delta_{ij} d\Omega$$

$$\bar{\underline{d}}^e = \{\bar{d}_p^e\}$$

$$\bar{d}_p^e = \int_{\Omega^e} N_a u_{0,i} d\Omega - \sum_q \bar{m}_{pq}^e g_q^e(0)$$

$$\dot{\underline{d}}_0 = \underline{\bar{M}}^{-1} \left\{ \underset{e=1}{\overset{n_{el}}{A}} (\dot{\bar{\underline{d}}}^e) \right\}$$

where $\dot{\bar{d}}_p^e = \int_{\Omega^e} N_a \dot{u}_{0,i} d\Omega - \sum \bar{m}_{pq}^e \dot{g}_q^e(0)$

before we had

$$\underline{M} \ddot{\underline{d}} + \underline{K} \underline{d} = \underline{F}(t)$$

and IC

now let's add a viscous damping matrix

$$\textcircled{*} \quad \underline{M} \ddot{\underline{d}} + \underline{C} \dot{\underline{d}} + \underline{K} \underline{d} = \underline{F}(t)$$

where $\underline{C}$ is a viscous damping matrix

if $\underline{K} = 0$ then we really have a 1st order system on $\dot{\underline{d}}$ but this also leads to nonlinear analysis also. We will now look at integrators of the "NEWMARK Family" i.e. "one-step" algorithms.

We can write $\textcircled{*}$ as 3 eqns

$$\underline{M} \underline{a}_{n+1} + \underline{C} \underline{v}_{n+1} + \underline{K} \underline{d}_{n+1} = \underline{F}_{n+1} = \underline{F}(t_{n+1})$$

where $\underline{a}_n \approx \underline{\ddot{d}}(t_n)$ $\underline{v}_n \approx \underline{\dot{d}}(t_n)$ $\underline{d}_n \approx \underline{d}(t_n)$
 $\underline{a}_n \approx \underline{\ddot{d}}(t_n)$ $\underline{v}_n \approx \underline{\dot{d}}(t_n)$ $\underline{d}_n \approx \underline{d}(t_n)$

These are Taylor series approx.
$$\begin{aligned} \underline{d}_{n+1} &= \underline{d}_n + \Delta t \underline{v}_n + \frac{\Delta t^2}{2} \{ (1-2\beta) \underline{a}_n + 2\beta \underline{a}_{n+1} \} \\ \underline{v}_{n+1} &= \underline{v}_n + \Delta t \{ (1-\gamma) \underline{a}_n + \gamma \underline{a}_{n+1} \} \end{aligned}$$
 where β is a parameter
" γ " " " } they control accuracy & stability of algorithm.

For the algorithmic eqns suppose we are given everything at d_n, x_n, a_n and we want to find $a_{n+1}, d_{n+1}, x_{n+1}$

for the initial step.

$$M_{a_0} = F_{a_0} - C_{v_{a_0}} - K_{a_0} \quad \text{now this is known.}$$

Now look at members of newmark family

Now look at members of newmark family

Ex. 1. suppose we look at the (constant) ^{implicit} average accel. method $\rho = 1/4$ $\delta = 1/2$

this is equiv to applying trapez rule to first order form of $M\ddot{d} + C\dot{d} + Kd = \ddot{u}$

what is first order form.

$$\text{let } \underline{y} = (\underline{d}, \underline{\dot{d}})^T \quad \text{let } \underline{f} = \underline{f}(\underline{y}, t) = \left\{ \begin{array}{l} \underline{\dot{d}} \\ \underline{M}^{-1} (\underline{F}(t) - \underline{C}\underline{\dot{d}} - \underline{K}\underline{d}) \end{array} \right\}$$

Thus $y = f(y, t)$ this first order form. Now apply trapez form.

now $\tilde{z}_n \times \dot{y}(t_n) \Rightarrow \tilde{y}_{n+1} = \tilde{y}_n + \frac{\Delta t}{2} (\tilde{z}_n + \tilde{z}_{n+1})$ trapez. rule $\tilde{y}_n \approx y(t_n)$

$$z_{n+1} = f(z_n, t_{n+1})$$

trap is uncond. stable & 2nd order accurate.

exercice show that the above integrator is the newmark member $w/\beta = 1/4$, $\gamma = 1/2$.

$$A_{\alpha} = \frac{1}{2} \left(\frac{1}{\alpha} \int_0^1 \frac{1}{\sqrt{1-t}} dt \right) \left(\frac{1}{\alpha} \int_0^1 \frac{1}{\sqrt{1-t}} dt \right) = \frac{1}{2} \left(\frac{1}{\alpha} \right) \left(\frac{1}{\alpha} \right) = \frac{1}{2\alpha^2}.$$

Ex. 2. if we use linear accel method $\beta = \frac{1}{6}$, $\gamma = \frac{1}{6}$ conditionally stable, 2nd order
this is an implicit method.

what is implicit or explicit

Implicit is a matrix problem needs to be solved to advance sol. from steps to steps.

explicit - no matrix problem needs to be solved.

The conditionally stable system is dependent on the highest e. value.

Ex. 3. Central difference method $\beta = 0$ $\gamma = \frac{1}{2}$

i.e. $\frac{1}{n} \frac{(d_{n+1} - d_{n-1})}{2\Delta t}$

$$a_n = (d_{n+1} - 2d_n + d_{n-1})/\Delta t^2$$

2nd order accurate

explicit; conditionally stable
attractive due to 2nd order & explicit

if $\underline{c} = 0$ and $\underline{M}$ is diag: $\underline{M} \underline{a}_{n+1} + \underline{K} \underline{d}_{n+1} = \underline{F}_{n+1}$
by some method,

if $\beta = 0$ $\underline{d}_{n+1} = \underline{d}_n + \Delta t \underline{v}_n + \frac{\Delta t^2}{2} \underline{a}_n$; if $\underline{d}_n, \underline{v}_n, \underline{a}_n$ are known; $\underline{d}_{n+1}$ is found

$(\underline{F}_{n+1} - \underline{K} \underline{d}_{n+1}) / [\underline{M}]$ explicit in nature
can get $\underline{v}_{n+1}$

3

$$\left. \begin{aligned} \underline{d}_{n+1} &= \underline{d}_n + \Delta t \underline{v}_n + (1-\beta) \frac{\Delta t^2}{2} \underline{a}_n \\ \underline{v}_{n+1} &= \underline{v}_n + (1-\gamma) \Delta t \underline{a}_n + \gamma \Delta t \underline{a}_{n+1} \end{aligned} \right\} \text{Newmark formulas}$$

are predictors
eliminate we can get

$\underline{d}_{n+1}^v$ known.
diag $\underline{K}$ hence we have a
 $(\underline{M} + \beta \Delta t^2 \underline{K})$ is implicit

substitution for $\underline{v}_{n+1}$

$$\underline{C} \underline{v}_{n+1}^v = \underline{K} \underline{d}_{n+1}^v$$

if $\underline{M}$ is diag
we are implicit system.

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$$\left. \begin{aligned} \underline{d}_{n+1} &= \underline{F}_{n+1} \\ \underline{v}_n + \frac{\Delta t^2}{2} \{ (1-2\beta) \underline{a}_n + 2\beta \underline{a}_{n+1} \} \\ \gamma \underline{a}_n + \gamma \underline{a}_{n+1} \\ \underline{d}_0 \text{ given} \end{aligned} \right\}$$

after this incop

We want to reduce this to an SDOF system. Recall EV problem.

Undamped EV problem $(\underline{K} - \lambda \underline{M}) \underline{\psi} = 0 \quad \Rightarrow \quad \{ \lambda_l, \underline{\psi}_l \} \quad 1 \leq l \leq n_{eq}$

also $\underline{\psi}_l^T \underline{M} \underline{\psi}_m = \delta_{lm}$ and $\underline{\psi}_l^T \underline{K} \underline{\psi}_m = \lambda_l \delta_{lm}$ (no sum) where $\underline{d} = \sum \underline{d}_i \underline{\psi}_i$

Look at the Rayleigh matrix damping is $\underline{C} = a \underline{M} + b \underline{K}$

then $\underline{\psi}_l^T \underline{C} \underline{\psi}_m = a \delta_{lm} + b \lambda_l \delta_{lm}$ λ is an EV of $(\underline{K} - \lambda \underline{M}) \underline{\psi} = 0$

we will define $\omega_l = (\lambda_l)^{1/2}$ if we remember $\underline{M}$ & $\underline{K}$ are pos. semidef & def respect.
 $= l^{th}$ undamped frequency ≥ 0

define $\xi_l = (a/\omega_l + b\omega_l)/2 = l^{th}$ modal damping factor.

now $\underline{C}$ is diag as well as $\underline{M}$ & $\underline{K}$

If we do modal decomp (as in heat eqn.) on the original problem

(3) $\begin{cases} \ddot{d} + 2\xi\omega\dot{d} + \omega^2 d = F & \text{these are the } \begin{smallmatrix} \text{typical} \\ \text{form} \end{smallmatrix} \text{ coeff} \\ d(0) = d_0 \\ \dot{d}(0) = \dot{d}_0 \end{cases}$

If we apply modal decomp on the temporally discretized version we get

(4) $\underline{a}_{n+1} + 2\xi\omega\underline{v}_{n+1} + \omega^2 \underline{d}_{n+1} = \underline{T}_{n+1}$
 same eqn for $\underline{d}_{n+1}, \underline{v}_{n+1}$ w/ $\underline{d}_0$ given & $\underline{v}_0 = \dot{\underline{d}}_0$ given

By taking (3) and temporally discretize we will get (4)

if we look at $\underline{d}_n$ and $\underline{d}_n \rightarrow \underline{d}(t_n)$ (t_n fixed) $\Rightarrow \underline{d}_n \rightarrow \underline{d}(t_n)$
 $\underline{v}_n \rightarrow \dot{\underline{d}}(t_n) \Rightarrow \underline{v}_n \rightarrow \dot{\underline{d}}(t_n)$

Now let's generalize the arguments from heat conduction by analyzing the algorithm. Now look at (4) in first order form.

$$\underline{y}_{n+1} = \underline{A} \underline{y}_n + \underline{L}_n$$

$\underline{y}_n = \begin{Bmatrix} \underline{d}_n \\ \underline{v}_n \end{Bmatrix}$ $\underline{A}$ amplif matrix $\underline{L}_n$ load vector

eqns for $\underline{d}_{n+1}$ & $\underline{v}_{n+1}$ and subst for $\underline{a}_{n+1}$ from $\underline{a}_{n+1} + \dots$ & for $\underline{a}_n$ write eqn for previous time step. This leaves eqns in $\underline{d}_n, \underline{d}_{n+1}, \underline{v}_n, \underline{v}_{n+1}$. This

leads to $A = \tilde{A}_1^{-1} \tilde{A}_2$

$$\tilde{A}_1 = \begin{bmatrix} 1 + \Delta t^2 \beta \omega^2 & 2\Delta t^2 \beta \xi \omega \\ \Delta t \gamma \omega^2 & 1 + 2\Delta t \gamma \xi \omega \end{bmatrix}$$

$$\tilde{A}_2 = \begin{bmatrix} 1 - \frac{\Delta t^2}{2} (1-2\beta) \omega^2 & \Delta t (1 - \Delta t (1-2\beta) \xi \omega) \\ -\Delta t (1-\gamma) \omega^2 & 1 - 2\Delta t (1-\gamma) \xi \omega \end{bmatrix}$$

$$\underline{L}_n = \tilde{A}_1^{-1} \left\{ \begin{array}{l} \frac{\Delta t^2}{2} [(1-2\beta)F_n + 2\beta F_{n+1}] \\ \Delta t [(1-\gamma)F_n + \gamma F_{n+1}] \end{array} \right\}$$

We now do as before: define $\underline{\tau}$ [truncation error vector - local] by replacing $\underline{y}_{n+1}$ by $\underline{y}(t_{n+1})$

$$\therefore \underline{y}(t_{n+1}) = A \underline{y}(t_n) + \underline{L}_{n+1} + \underline{\tau}(t_n) \quad \text{where } \underline{y}(t_n) = \begin{Bmatrix} d(t_n) \\ \dot{d}(t_n) \end{Bmatrix}$$

for newmark we can show that

$$\|\underline{\tau}\| = O(\Delta t^{k+1}) \quad \begin{array}{ll} k=2 & \text{when } \gamma = \frac{1}{2} \\ k=1 & \text{otherwise} \end{array} \quad \begin{array}{l} \text{ie. doesn't depend} \\ \text{on } \beta \text{ through} \\ \text{2nd order.} \end{array}$$

aside $\gamma = \frac{1}{2}$ $\beta = \frac{1}{12}$ $C=0 \Rightarrow k=3$ the fox-goodwin method
implicit and conditionally stable

We pick β to get unconditional stability or to get explicit type algorithm
($\beta = \frac{1}{4}$ implicit) ($\beta = 0$ explicit stable)

Now subtract $\underline{d}(t_n)$ from $\underline{d}_n$ then

$$\underline{y}_n - \underline{y}(t_n) = \underline{e}(t_n) = \tilde{A}^n \underline{e}(0) - \sum_{i=0}^{n-1} \tilde{A}^i \underline{\tau}(t_{n-1-i})$$

we want the matrix norm ≤ 1 for convergence.

must now define consistent matrix norm.

Consider $\underline{y} = \begin{Bmatrix} y_1 \\ \vdots \\ y_m \end{Bmatrix} \in \mathbb{R}^m$ vector examples of $\|\cdot\|$ norm (Euclidean) $(\underline{y}^T \underline{y})^{1/2} = \|\underline{y}\|_2 \in \mathbb{R}$

matrix norm for $A \in \mathbb{R}^{m \times m}$

$$\|\underline{A}\| \in \mathbb{R}^{m \times m} \rightarrow \mathbb{R} \quad \text{e.g. (i) } \|\underline{A}\| \geq 0 \quad \|\underline{A}\| = 0 \text{ iff } \underline{A} = \underline{0}$$

$$(ii) \quad \|\underline{A} \alpha\| = |\alpha| \|\underline{A}\|$$

$$(iii) \quad \|\underline{A} + \underline{B}\| \leq \|\underline{A}\| + \|\underline{B}\|$$

$$(iv) \quad \|\underline{A} \underline{B}\| \leq \|\underline{A}\| \|\underline{B}\|$$

so far matrix norm & vector norm are defined separately. Want now to define vector & matrix norm that are compatible, ie $\|\underline{A} \underline{y}\| \leq \|\underline{A}\| \|\underline{y}\| \Leftarrow$ compatibility

We now ask if $\exists$ compatible matrix & vector norm,
yes $\exists$ one that can be naturally defined.

example $\|\underline{y}\|_2 = (\underline{y}^T \underline{y})^{1/2}$ $\|\underline{A}\|_2 = (\max \text{ eigenvalue of } \underline{A}^T \underline{A})^{1/2}$

convergence in one norm $\Rightarrow$ convergence in any other norm. We will henceforth use the $\|\cdot\|_2$ norms.

now look at $\underline{A}$ & look at its EV. $\{\lambda^i\}_{i=1}^m$ $\underline{A}$ need not be sym. in general.

$\therefore$ its EV are complex thus $\bar{\lambda}^i = \text{Re}(\lambda^i) - i \text{Im}(\lambda^i)$

now $|\lambda^i| = (\lambda^i \bar{\lambda}^i)^{1/2}$ & pick $\max_{i=1, \dots, m} |\lambda^i| = \rho(\underline{A})$ this is the spectral rad. of $\underline{A}$

stability is defined by that no. $\rho(\underline{A})$

consider any natural matrix norm: $\rho(\underline{A}) \leq \|\underline{A}\|$

In linear systems $\rho(\underline{A})$ is enough to define stability but for non lin systems need also norm.

example EX: $\begin{bmatrix} 0 & k \\ 0 & 0 \end{bmatrix}$ $k \gg 0$ $\rho(\underline{A}) = 0$ $\|\underline{A}\| = k$
 $\underline{A}^T \underline{A} = \begin{bmatrix} k^2 & 0 \\ 0 & 0 \end{bmatrix}$ or solve $\lambda^2 - k^2 = 0 \therefore \|\underline{A}\|_2 = k$

suppose $\underline{A}$ has linearly indep eigenvectors.

$$\underline{A} = \underline{P} \underline{\Lambda} \underline{P}^{-1}$$

$\underline{P}$ matrix of EV $\underline{\Lambda} = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_m \end{bmatrix}$ matrix of EV. $\lambda_m = \max_i \lambda_i$

$$\underline{A}^n = \underline{P} \underline{\Lambda}^n \underline{P}^{-1} \text{ if } \lambda_m < 1 \Rightarrow \lambda_m^n \rightarrow 0 \text{ & } \underline{A}^n \rightarrow 0$$

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from last time

$$\|\underline{A}\|_{(2)} = (\max \text{ EV of } \underline{A}^T \underline{A})^{1/2}$$

$$\|\underline{y}\|_{(2)} = (\underline{y} \cdot \underline{y})^{1/2} \text{ corresponding vector norm.}$$

$$\{\lambda^i\}_{i=1}^n = \text{EV of } \underline{A} ; \rho(\underline{A}) = \max_i |\lambda_i| \leq \|\underline{A}\|$$

Define: the Jordan Block $\underline{J}$ is a square matrix

$$\underline{J} = \begin{bmatrix} \lambda & 1 & 0 \\ 0 & \lambda & \vdots \\ 0 & 0 & \lambda \end{bmatrix} \quad \lambda \text{ complex}$$

L. $\underline{A}$ is given. $\exists$ a matrix $\underline{Q}$ s.t. $\underline{Q}^{-1} \underline{A} \underline{Q} = \begin{bmatrix} \underline{J}_1 & & \\ & \ddots & \\ & & \underline{J}_k \end{bmatrix}$

where $\underline{J}_i$ are jordan blocks

$$= \begin{bmatrix} \lambda_i & 1 & 0 \\ 0 & \lambda_i & 1 \\ & & \ddots \\ 0 & & \lambda_i \end{bmatrix}$$

This is a jordan form.

Thm: If $\underline{A}$ has L.I. $\vec{e}_i$ then for $\underline{P} = [\vec{e}_1, \vec{e}_2, \dots] \Rightarrow \underline{P}^{-1} \underline{A} \underline{P} = \underline{\Lambda}$

$$\text{where } \underline{\Lambda} = \begin{bmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_m \end{bmatrix}$$

We're doing this work to provide the groundwork to show that $\underline{M}d + \underline{K}d = \underline{F}$ is convergent. We do this to find out how $\underline{A}^n$ grows.

we will now look at the (spectral) stability -

we want $\|\underline{A}^n\| \leq \text{const} \forall n$ since error contains powers of $\underline{A}$

To achieve this we will require spectral stability of $\underline{A}$

ie (i) $\rho(\underline{A}) \leq 1$

(ii) λ_i of multiplicity > 1 satisfy that $|\lambda_i| < 1$

look at $\underline{A}$ for 2x2 pertaining to the Newmark Case.

$$\text{For newmark case } \underline{A}^n = (\underline{P} \underline{\Lambda} \underline{P}^{-1})^n = \underline{P} \underline{\Lambda}^n \underline{P}^{-1}$$

$$\text{then } \|\underline{A}^n\| \leq \|\underline{\Lambda}^n\| \|\underline{P}\| \|\underline{P}^{-1}\| = \text{const} \|\underline{\Lambda}^n\|$$

$$\text{for 2x2 } \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix}^n = \begin{bmatrix} \lambda_1^n & 0 \\ 0 & \lambda_2^n \end{bmatrix}; \text{ if } |\lambda_i| \leq 1 \Rightarrow |\lambda_1^n|, |\lambda_2^n| \leq 1$$

$$\text{now } \rho(\underline{A}^n) = |\lambda_2^n| \leq 1$$

but $\lambda_1 \neq \lambda_2 \therefore 2^{\text{nd}}$ condition is inoperative

for the most general canonical form.

$$\underline{A}^n = (\underline{Q} \underline{J} \underline{Q}^{-1})^n = (\underline{Q} \underline{J}^n \underline{Q}^{-1})$$

$$\|\underline{A}^n\| \leq \|\underline{Q}\| \|\underline{J}^n\| \|\underline{Q}^{-1}\| = \text{const} \cdot \|\underline{J}^n\|$$

$$\text{suppose } \underline{J} = \begin{bmatrix} \lambda & 1 \\ 0 & \lambda \end{bmatrix}$$

$$\underline{J}^n = \begin{bmatrix} \lambda^n & n\lambda^{n-1} \\ 0 & \lambda^n \end{bmatrix}$$

$\therefore$ if $|\lambda| = 1$ then

$$\rho(\underline{J}^n) = \lambda^n \text{ and } \|\underline{A}\| \Rightarrow |(n-1)\lambda^{n-1}| < 1$$

then if $|\lambda| = 1$ $\|\underline{A}\|$ blows up

$\therefore$ need (ii) condition

the 2 conditions $\rho(A) \leq 1$ & multiplicity are sufficient but not necessary conditions for stability. ex $[I] = A$ $I^n = I$ & is still stable $|\lambda| = 1$

The Now suppose we fix t_n ; assume A is spectrally stable

Assume

$$\|\underline{e}(t)\|_{(2)} \leq c \Delta t^{k+1} \quad \forall t \in [0, T]$$

then $\underline{e}(t_n) = \underline{a}_n - \underline{e}(t_n) \rightarrow 0$ as $\Delta t \rightarrow 0$

Pf: we found $\underline{e}(t_n) = A^n \underline{e}(0) - \sum_{i=0}^{n-1} A^i \underline{e}(t_{n-1-i})$
if $\underline{e}(0)$ is fixed then

$$\begin{aligned} \|\underline{e}(t_n)\| &\leq \sum_{i=0}^{n-1} \|A^i\| \|\underline{e}(t_{n-1-i})\| \leq n \cdot \rho(A) \|\underline{e}\| \leq n \cdot \overbrace{c \rho(A)}^{\text{const.}} \Delta t^{k+1} \\ &\leq \text{const } n \Delta t^{k+1} = \text{const } t_n \Delta t^k \therefore \|\underline{e}(t_n)\| \text{ is } O(k) \end{aligned}$$

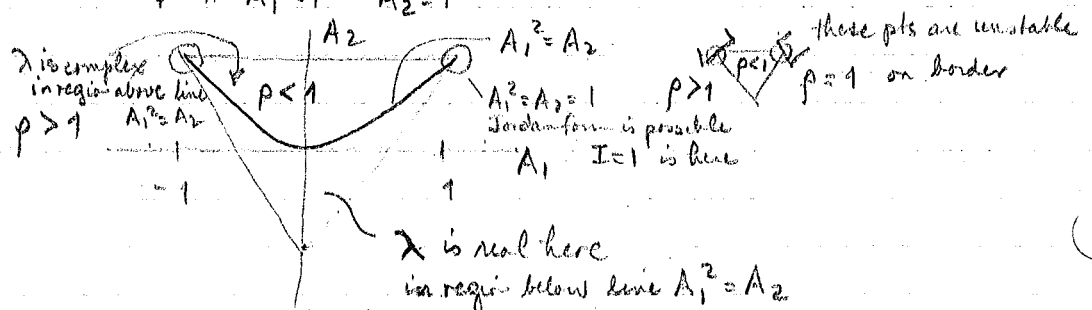
Spectral stability of newmark ^{family} A being 2×2
look at EV. $0 = \det(A - \lambda I) = \lambda^2 - 2A_1\lambda + A_2$

$$\begin{aligned} A_1 &= \frac{1}{2} \text{tr } A = \frac{1}{2}(a_{11} + a_{22}) \\ A_2 &= \det A \\ \therefore \lambda_{\pm} &= A_1 \pm \sqrt{A_1^2 - A_2} \end{aligned}$$

if $A_1^2 < A_2$ λ is complex
if $A_1^2 = A_2$ multiple root
if $A_1^2 > A_2$ λ is distinct real roots

for spectral stability are equivalent to $-(A_2+1)/2 \leq A_1 \leq (A_2+1)/2$ $-1 \leq A_2 \leq 1$
& $-1 < A_1 < 1$ $A_2 = 1$

exam question



if we define $\Omega = \omega \Delta t =$ "sampling frequency"

$$\begin{aligned} A_1 &= 1 - [\xi \Omega + \Omega^2 (\gamma + \frac{1}{2}) / 2] / D \\ A_2 &= 1 - [2\xi \Omega + \Omega^2 (\gamma - \frac{1}{2}) / 2] / D \\ D &= 1 + 2\xi \Omega + \beta \Omega^2 \end{aligned}$$

Summary of Spectral stability for newmark
unconditionally stable

$$\frac{1}{2} \leq \gamma \leq 2\beta$$

w/ K has positive est

$$Q = aM + bK \quad \text{assume } a, b > 0$$

$\xi^2 \geq 0$ system is physically stable

conditional stability

$$\gamma \geq \frac{1}{2} \neq (2\beta \leq \gamma)$$

$\gamma < \frac{1}{2}$ numerical amplifier

$$\Rightarrow \Omega = \omega \Delta t \leq \Omega_{crit} = \frac{\xi(\gamma - \frac{1}{2}) + [\gamma/2 - \beta + \xi^2(\gamma - \frac{1}{2})^2]^{1/2}}{\gamma/2 - \beta}$$

we assume that $\beta \geq 0$

your max frequency ω determines Δt . If we define in terms of period of vibration ie $\omega T = 2\pi$ T if fixed time then

$$\Delta t \cdot \frac{\omega}{2\pi} = \frac{1}{T} \cdot \Delta t \Rightarrow \frac{\Delta t}{T} \leq \frac{\Omega_{crit}}{2\pi}$$

Consider $\xi = 0$ physically undamped case. then $\Omega_{crit} = (\gamma/2 - \beta)^{-1/2}$
for $\xi > 0$ $\gamma > \frac{1}{2}$ we increase stable time step since $\xi(\gamma - \frac{1}{2}) > 0$

Look at some known methods $\xi = 0$

accuracy determines Δt trap rule (1) (average accel) $\beta = 1/4$ $\gamma = 1/2$ stability cond. unconditionally $2\beta \geq \gamma \geq 1/2$ $1/2 \geq 1/2 \geq 1/2$ yes

Stability determines Δt	(2) lin accel	$\beta = 1/6$	$\gamma = 1/2$	$\Omega_{crit} = 2\sqrt{3} \approx 3.464$	IMPLICIT
	(3) Fox Goodwin	$\beta = 1/2$	$\gamma = 1/2$	$\Omega_{crit} = \sqrt{6} \approx 2.449$	
	(4) Central Diff	$\beta = 0$	$\gamma = 1/2$	$\Omega_{crit} = 2$	
					EXPLICIT

between implicit methods use (1). Between condit stable systems use (2) since explicit if you don't get much better Δt by using (2), (3).

for linear systems use (1). For non linear & wave prop type problems use (4).

5/26/81

Office Hrs. Today only 3-4

Stability conditions: Newmark

requires implicit Unconditional: $2\beta \geq \gamma \geq 1/2$

Conditional: $\gamma \geq 1/2$, $0 \leq \beta < \gamma/2$;
can be related to
dissipative

$$\omega \Delta t \leq \frac{2\pi}{\Omega} \Delta t \quad \Omega \leq \Omega_{crit} = \frac{\gamma(\gamma - 1/2) + \sqrt{[\gamma/2 - \beta] + \xi^2(\gamma - 1/2)^2}}{[\gamma/2 - \beta]}$$

conservative bound $\Omega \leq (\gamma/2 - \beta)^{-1/2}$ undamped; $\xi > 0$ damped case ($\xi = 0$)

"Oscillatory Response" $A_1^2 < A_2$ involves complex conjugate.

the above is true & so: we are governed by highest λ $\therefore (\Delta t)_{\lambda_{neg}} \leq (\Delta t)_{\lambda_{neg-1}}$

$\therefore (\Delta t)_{\lambda_{neg}}$ is conservative wrt whole system.

Stronger (slightly) notion of stability $\gamma \geq 1/2$ $\xi < 1$ undamped.

unconditional requires $\beta \geq (\gamma + 1/2)^2/4$

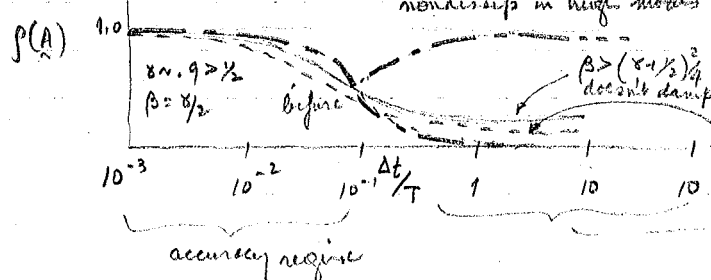
conditional $\beta < (\gamma + 1/2)^2/4$ $\Omega < \Omega_{bifurcation} = \frac{\xi(\gamma - 1/2)/2 + \sqrt{[(\gamma + 1/2)^2/4 - \beta] + \xi^2(\gamma - 1/2)^2}}{[\xi(\gamma - 1/2)/2]}$

bifurcation is pt at which complex conj roots become real $[(\gamma + 1/2)^2 - \beta]$

for $\xi = 0$ $\Omega_{bif} = [(\gamma + 1/2)^2 - \beta]^{-1/2}$

$[p(A)]^n$ ~ amplification "factor" of the mode in question. If $p(A) < 1$ then we get dissipation. For example

this is for unconditional stability with $\beta = 1/2$ being the lower limit non-damping in high modes



in accuracy we will take 10 or more steps per cycle.
high modes - artifacts of spatial discretization occur here; we are interested in damping system.
so...

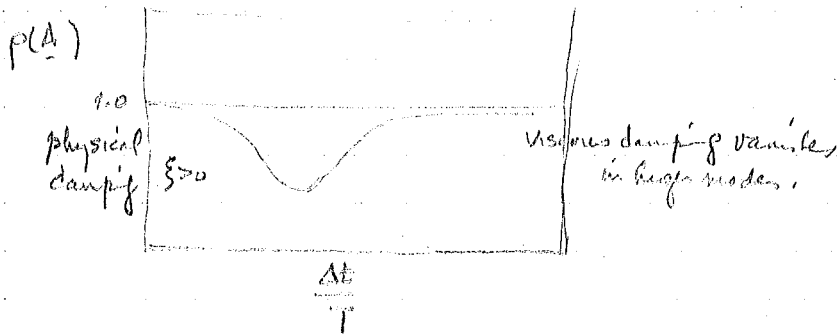
if we take $\beta = (\gamma + 1/2)^2/4$ no bifurcation and get dissipation

if $\beta > 1/2$ & $(\gamma + 1/2)^2 > \beta$ then we get bifurcation but don't return back to 1.

Asymptote is ≥ 1 . We will get damped oscillatory response.

trapez has a $\rho(\underline{A}) = 1 \quad \forall \omega, \Delta t$ $\gamma = 1/2 \quad \beta = 1/4$
 newmark tends to damp lower modes (which isn't good). Other methods tend to do better in lower modes.

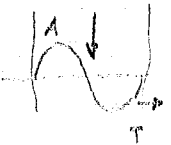
Suppose we fix $\xi \neq 0$ (VISCOUS DAMPING)



what we want is accuracy in low modes, dissipate in high modes

Accuracy Measures

2 types of phenomena for sinusoidal responses $T \uparrow$ or amplitude $\downarrow$
 dispersive (dissipative)



Look at $\ddot{d} + 2\xi\omega\dot{d} + \omega^2 d = 0$ S.D.O.F. model problem (*)
 + I.C.

suppose $0 \leq \xi < 1$ underdamped oscillatory

$$d(t) = e^{-\xi\omega t} (d_0 \cos \omega_d t + \frac{\dot{d}_0 + \xi\omega d_0}{\omega_d} \sin \omega_d t)$$

$$\omega_d = (1 - \xi^2)^{1/2} \omega$$

case (ii) critically damped $\xi = 1$

(iii) overdamped $\xi > 1$

look at

Continuous System
 underd

discrete analog
 complex conj λ 's

critic

Real & equal

Overdamped

Real & distinct

discrete algorithm (ie newmark) mimics (*)

$$d_{n+1} = A_{11} d_n + A_{12} v_n \quad \text{for no external forces}$$

$$v_{n+1} = A_{21} d_n + A_{22} v_n$$

at step n

$$d_n = A_{11} d_{n-1} + A_{12} v_{n-1}$$

$$v_n = A_{21} d_{n-1} + A_{22} v_{n-1}$$

desire
want to find a difference equation form

$$d_{n+1} - 2A_1 d_n + A_2 d_{n-1} = 0$$

$\frac{1}{2} \text{trace } A$ $\det A$

$\lambda_{1,2}$ eigenvalues of A $\lambda_1 \neq \lambda_2$
soln $d_n = C_1 \lambda_1^n + C_2 \lambda_2^n$
(const)

want to write this in terms of (*) assuming complex conjugate λ

desired write exact soln. $y(t_{n+1}) = \tilde{E} y(t_n)$ $\tilde{E}$ is the evolution operator
if is the exact amplification matrix
looking at EV of $\tilde{E}$ then $\lambda(\tilde{E}) = e^{-\xi \Omega \pm i \bar{\Omega}}$ $\Omega, \bar{\Omega} = \omega \Delta t, \bar{\omega} \Delta t$

thus $d(t_n) = C_1 \lambda_1(\tilde{E})^n + C_2 \lambda_2(\tilde{E})^n$ complete parallel of discrete sol.

since we have A then
we can now assume $\lambda = e^{-\xi \bar{\Omega} \pm i \bar{\Omega}}$ & solve for $\xi, \bar{\omega}$
using the following

for $\xi < 1$

$$\bar{\Omega}_d = \arctan \left(\frac{A_2}{A_1} - 1 \right)^{1/2} \quad \xi = \frac{-1}{2\bar{\Omega}_d} \ln A_2 \quad \bar{\Omega}_d = (1 - \xi^2)^{1/2} \bar{\Omega}$$

actual value of $\bar{\Omega}$

thus $d_n = e^{-\xi \bar{\omega} t_n} (d_0 \cos \bar{\omega}_d t_n + \bar{c} \sin \bar{\omega}_d t_n)$ from which a plot can get a picture of
dissip dispersion what mode is actual doing at each time step.

if set $d_1, d_0 \Rightarrow \bar{c} = \frac{d_1 - A_1 d_0}{(A_2 - A_1^2)^{1/2}} = \frac{1}{2} \frac{(A_{11} - A_{22}) d_0 + A_{12} v_0}{(A_2 - A_1^2)^{1/2}}$ if $\xi = 0$ $\bar{\omega} = \bar{\omega}_d$

thus $\xi, \bar{\omega}$ are the parameters to look at

people have looked at the following to characterize an algorithm

is ξ algorithmic damping ratio

$\bar{\omega}$ " frequency, however people use $\frac{\bar{T} - T}{T}$ the relative period error
where $T\bar{\omega} = 2\pi$ $\bar{T}\bar{\omega} = 2\pi$

we mostly look at the cases where $\xi = 0$ even if algorithm introduces damping
for newmark damping is weakest feature of algorithm.

for newmark $\bar{\xi} = \xi + (8 - 1/2) \Omega_{/2} = \omega \Delta t = (\Delta t/T)^{1/2} \pi$ thus we get a 1st order error in damping.

500

$$-\frac{h}{2} \begin{bmatrix} \lambda & 0 \\ 0 & \lambda \end{bmatrix} + \frac{c^2}{h} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} = \underline{K} - \lambda \underline{M} = \begin{bmatrix} \frac{c^2}{h} - \frac{\lambda h}{2} & -\frac{c^2}{h} \\ -\frac{c^2}{h} & \frac{c^2}{h} - \frac{\lambda h}{2} \end{bmatrix}$$

$$\therefore \left(\frac{c^2}{h} - \frac{\lambda h}{2} \right)^2 - \left(\frac{c^2}{h} \right)^2 = 0$$

$$\frac{c^4}{h^2} - \frac{2c^2\lambda h^2}{2h^2} + \frac{\lambda^2 h^4}{4h^2} - \frac{c^4}{h^2} \Rightarrow \lambda = 0 \text{ or } \frac{\lambda h^2}{4} - c^2 = 0$$

$$\lambda = \frac{4c^2}{h^2}$$

$$\omega = \sqrt{\lambda} = \frac{2c}{h}$$

Lumped

$$\frac{c^2}{h} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} - \frac{h}{6} \begin{bmatrix} 2\lambda & \lambda \\ \lambda & 2\lambda \end{bmatrix} = \begin{bmatrix} \left(\frac{c^2}{h} - \frac{\lambda h}{3} \right) & -\frac{c^2}{h} - \frac{h\lambda}{6} \\ -\frac{c^2}{h} - \frac{h\lambda}{6} & \frac{c^2}{h} - \frac{\lambda h}{3} \end{bmatrix}$$

$$\left(\frac{c^2}{h} - \frac{\lambda h}{3} \right)^2 - \left(\frac{c^2}{h} + \frac{h\lambda}{6} \right)^2 = 0$$

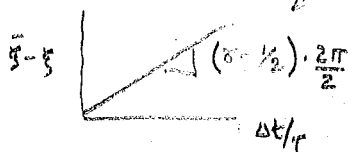
$$\therefore \left[\frac{c^2}{h} - \frac{\lambda h}{3} - \frac{c^2}{h} - \frac{h\lambda}{6} \right] \left[\frac{c^2}{h} - \frac{\lambda h}{3} + \frac{c^2}{h} + \frac{h\lambda}{6} \right] = 0$$

$$\therefore \lambda = 0 \text{ for here or } \frac{2c^2}{h} - \frac{\lambda h}{6} = 0 \text{ or } \lambda = \frac{12c^2}{h^2}$$

$$\omega = \lambda^{\frac{1}{2}} = 2\sqrt{3} \frac{c}{h}$$

(Higher order test)

If we plotted the graph we want to match modal functions. However we don't get it.



but plotting $\frac{\bar{T} - T}{T}$ turns out to be $O(\Omega^2)$

Other things people have looked at has been amplitude decay. This is essentially equivalent to $\bar{\xi}$

Recall for conditionally stable algorithms.

* is for heat conduction for forward differences

$$\alpha = 0 \quad \Delta t \leq 2/\lambda_{\max} \quad \text{in eqn on pg 16}$$

* is for elastodynamics for central differences

$$\Delta t \leq 2/\omega_{\max} \quad \xi = 0 \quad \beta = 0 \quad \delta = 1/2 \quad \text{in eqn on pg 22}$$

$$\lambda_{\max} \leq \max_e \lambda_{\max}^e (\text{elemental } \lambda)_{\max} \quad \omega_{\max} \leq \max_e (\omega_{\max}^e)$$

we want to look at this in the lumped mass context; explicit schemes cond. stable.

Example (1) LIN elem

elastic situation for rod $u_{tt} = c^2 u_{xx}$ $c = \sqrt{E/\rho}$ bar wave vel

then $k^e = \frac{c^2}{h} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}$
from 1st quarter

$m^e = \frac{\text{exact}}{\text{constant}} = \left[\int_{\Omega^e} N_A N_B d\Omega \right] = \frac{h}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix}$
See Wingt handout pg 8

sum of $\delta_{ij} = \text{mass} = \frac{2+1+1+2}{6} = 1$

the Lumped mass matrix $\tilde{m}^e = \frac{h}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$ $\Delta \quad \rho \omega_{ij}^e = \lambda_1 = \lambda_2 = 1$

$N_1 = \frac{1-x}{2}$ $N_2 = \frac{x+1}{2}$ then we get $(N_1^2|_{x=-1}, N_1 + N_2, N_2^2|_{x=1})^{1/2} = \frac{h}{2}$ also same for N_2^2 ; $N_1 N_2|_{x=-1,1} = 0$

Lumped $\tilde{m}^e$ can also be obtained using trapez rule i.e. $\xi_1 = 1, \xi_2 = -1$ $W_1 = W_2 = 1$ on $\int N_A N_B d\Omega$
leads to $\tilde{m}^e = \delta_{ab} h/2$ $\frac{1}{2} [N_A N_B|_{\xi=-1} W_1 + N_A N_B|_{\xi=1} W_2] \cdot h$

we will get $\omega_2 = 0$ (freq of rigid modes) $\& \quad \omega_1 \neq 0 = \omega_{\max}$ obtained from $\det(\tilde{K} - \lambda \tilde{M}) = 0$
see wingt notes pg

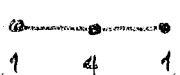
LUMPED $\lambda^e = \omega_{\max}^e = 2c/h$ $\checkmark$ now $\Delta t \leq \frac{2}{\omega_{\max}^e} = \frac{h}{c}$ and the smallest value governs

Δt is the transit time for a bar wave to travel 1 element

for the consistent case $\omega_{\max}^e = 2\sqrt{3} c/h^e$ $\& \quad$ thus $\Delta t = \frac{h^e}{\sqrt{3} c} < \frac{h^e}{c}$

General rule Consistent case tend to produce highest frequencies i.e. shorter ^{critical} time steps than ^{lumped}

thus may require 2x no. of time steps

Example 2  quadrature

$$\xi_1 = -1, \xi_2 = 0, \xi_3 = 1$$

ratio of masses produced by Simpson's rule if using lumped masses

$$m = \frac{h}{6} \cdot \begin{bmatrix} 1 & 0 & 0 \\ 0 & 4 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Exercise do the consistent case, $\int N^T N d\xi$

$$N_1 = \frac{\xi(\xi-1)}{2}, N_2 = (\xi^2-1), N_3 = \frac{(\xi+1)\xi}{2}$$

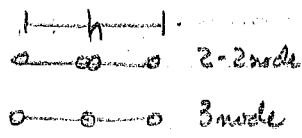
$$d\Omega = j d\xi, \quad j = X, \xi = \frac{\partial X}{\partial \xi} \sum N_i x_i = \left[\frac{2\xi-1}{2} x_1 - 2\xi x_2 + \frac{2\xi+1}{2} x_3 \right]$$

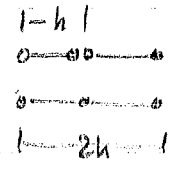
if $x_1 = -h/2, x_2 = 0, x_3 = h/2, \xi = 1/2$

now $\omega_{max} = 2\sqrt{6} c/h$

$$\Rightarrow \Delta t \leq \frac{2}{\omega} = \frac{h}{\sqrt{6} c} = .4082 \times \Delta t \text{ (of 2 node lumped)}$$

this result is compared to 2-node elements $\therefore$



we can compare $\frac{1-h}{2}$ also as follows 

$$\Delta t_{3node} = .4082 \Delta t_{2node}$$

this leads to $\Delta t = .8164$ of Δt for 2-node

in both, crit Δt is more stringent for 3 node element.

for the heat cond. solutions $\lambda_{max} = \omega_{max}^2$ replace c^2 in elasticity by k

thus $\Delta t \leq 2/\lambda_{max} = \frac{h^2}{2k}$ for the 2 node lumps.

note that $\Delta t \sim O(h^2)$ this is a more stringent condition than the elasticity problem.

For 4 node quadrilateral 1 pt quadrature on stiffness (ie accepts linear fns exactly) hence $\Rightarrow$ trapezoidal rule on m (thus lumped mass)).

$m \sim$ lumped $\Rightarrow 1/4$ at each node.

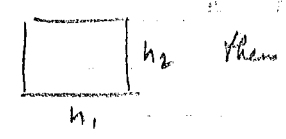
$$w_{max} \leq c_d g^{1/2}, \quad c_d = \left(\frac{\lambda + 2\mu}{\rho} \right)^{1/2}$$

g is a geometric parameter $= \frac{4}{A^2} \sum_{i=1}^2 \sum_{a=1}^4 b_{ia}^2 \rho_{ia}$

area/cell

$$[b_{ia}] = \frac{1}{2} \begin{bmatrix} y_1 - y_4 & y_3 - y_1 & y_4 - y_2 & y_1 - y_3 \\ x_4 - x_2 & x_1 - x_3 & x_2 - x_4 & x_3 - x_1 \end{bmatrix}$$

Now $\Delta t \leq 2/(c_d g^{1/2})$ if we assume



$$\leq \frac{1}{\sqrt{2}} \left[c_d (h_1^2 + h_2^2) \right]^{1/2}$$

if $h_1 = h_2 = h$ then $\Delta t \leq \sqrt{2} h_c$

however this is worse; yet this result is a sufficient condition

Linear Multistep Methods (LMS) methods

we will look at first order system.

$$\dot{\underline{y}} = \underline{f}(\underline{y}, t) = \underline{G} \underline{y} + \underline{H}(t)$$

A k -step LMS method can be written as.

$$\sum_{i=0}^k \{ \alpha_i \underline{y}_{n+1-i} + \Delta t \beta_i \underline{f}(\underline{y}_{n+1-i}, t_{n+1-i}) \} = \underline{0}$$

LMS terminology

LMS method explicit case: $\beta_0 = 0$

implicit: if $\beta_0 \neq 0$.

backward diff: $\beta_i = 0 \quad i \geq 1$

α & β are param. that define algorithm.

[example generalized trapez family can be expressed in this notation
 $k=1 \quad \alpha_0 = -\alpha, \alpha_1 = 1 \quad \beta_0 = -\alpha \quad \beta_1 = \alpha - 1$ where α is param. of generalized trapezoidal family.

for the heat eqn. if $\underline{G} = -\underline{M}^{-1} \underline{K}$ & $\underline{H}(t) = \underline{M}^{-1} \underline{F}(t)$
 and $\underline{y} = \underline{d}$ we get the generalized trap. family.

6/2/81

Linear Multistep Methods for 1st Order Systems - Continued

$$\dot{\underline{y}} = \underline{f}(\underline{y}, t) = \underline{G} \underline{y} + \underline{H}(t) \quad (1)$$

for the semi discrete heat eqn.

$$\underline{M} \underline{\dot{d}} + \underline{K} \underline{d} = \underline{F}$$

if we let

$$\underline{G} = -\underline{M}^{-1} \underline{K} \quad (4)$$

$$\underline{y} = \underline{d} \quad (3)$$

and $\underline{H} = \underline{M}^{-1} \underline{F} \quad (5)$ then we get the above

$\underline{M}$ is positive def

$\underline{K}$ " " semi def

SEE ABOVE FOR REVIEW :

- if $\alpha = 0$ we get explicit method.
- $\alpha = 1$ we get a backward difference
- if $k \rightarrow \infty$ we get computational difficulties since we must keep more vectors (which costs money).

let us look at the spectrum of $\underline{G}$

$$\text{ie } (\underline{G} - \lambda \underline{I}) \underline{\psi} = 0$$

subst. for $\underline{G}$ to get

$$(\underline{K} + \lambda \underline{M}) \underline{\psi} = 0$$

$$\text{thus the } \lambda_j = -\lambda_j^* \leq 0 \quad \lambda \left. \begin{array}{l} \text{was soln for } (\underline{K} - \lambda \underline{M}) \underline{\psi} = 0 \\ \text{and } \underline{\psi} = \underline{\psi} \end{array} \right\}$$

the λ_j fall on negative real axis.

Now we look at the semi-discrete eqns of motion.

General form was

$$\underline{M} \ddot{\underline{d}} + \underline{C} \dot{\underline{d}} + \underline{K} \underline{d} = \underline{F}$$

$$\underline{M} > 0$$

$$\underline{C}, \underline{K} \geq 0$$

$$\text{Now let } \underline{y} = \begin{pmatrix} \underline{d} \\ \dot{\underline{d}} \end{pmatrix}$$

$$\text{where } \underline{G} = \begin{pmatrix} 0 & \underline{I} \\ -\underline{M}^{-1}\underline{K} & -\underline{M}^{-1}\underline{C} \end{pmatrix}$$

$$\underline{H} = \begin{pmatrix} 0 \\ \underline{M}^{-1}\underline{F} \end{pmatrix}$$

$$\text{then } \dot{\underline{y}} = \underline{G} \underline{y} + \underline{H}$$

$$\text{Now let } \underline{\psi} = \begin{pmatrix} \underline{\psi}_1 \\ \underline{\psi}_2 \end{pmatrix}$$

the generalized EV

$$(\underline{G} - \lambda \underline{I}) \underline{\psi} = \left(\begin{bmatrix} 0 & \underline{I} \\ -\underline{M}^{-1}\underline{K} & -\underline{M}^{-1}\underline{C} \end{bmatrix} - \lambda \begin{bmatrix} \underline{I} & 0 \\ 0 & \underline{I} \end{bmatrix} \right) \begin{pmatrix} \underline{\psi}_1 \\ \underline{\psi}_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$\text{using: } \underline{\psi}_2 - \lambda \underline{\psi}_1 = 0 \quad (1)$$

$$-\underline{M}^{-1}\underline{K} \underline{\psi}_1 - (\underline{M}^{-1}\underline{C} + \lambda \underline{I}) \underline{\psi}_2 = 0$$

subst for $\underline{\psi}_2$ for (1)

$$\therefore -\underline{K} \underline{\psi}_1 - (\underline{C} + \lambda \underline{M}) \lambda \underline{\psi}_1 = 0$$

$$\text{or } (\lambda^2 \underline{M} + \lambda \underline{C} + \underline{K}) \underline{\psi}_1 = 0$$

$$\text{define now } \hat{\underline{\psi}}_l^T \underline{M} \hat{\underline{\psi}}_m = \delta_{lm}$$

$$\text{and } \hat{\underline{\psi}}_l^T \underline{K} \hat{\underline{\psi}}_m = \omega_l^2 \delta_{lm}$$

if $\underline{C}$ diagonalized then we can find a simple sol to the above
ie use rayleigh damping $\underline{C} = a \underline{M} + b \underline{K}$

then $\hat{\psi}_L^T C \hat{\psi}_m = a \delta_{lm}^2 + b \omega_L^2 \delta_{lm} = 2 \xi_L \omega_L \delta_{lm}$

thus $\lambda_g^2 + 2 \xi_L \omega_L \lambda_g + \omega_L^2 = 0$ $\lambda = \lambda_g$ KL

for a SDOF $\lambda_g^2 + 2 \xi \omega \lambda_g + \omega^2 = 0$

oscillatory

if $0 \leq \xi < 1$ system is underdamped then $\lambda_{1,2} = -\xi \omega \pm i \omega \sqrt{1-\xi^2}$ conjugates

$\xi = 1$ critically damped then $\lambda_{1,2} = -\xi \omega = -\omega$ (double root)

$\xi > 1$ overdamped $\lambda_{1,2} = -\xi \omega \pm \omega \sqrt{\xi^2 - 1}$ real

1. Remark if $\xi = 0$ (no damping) $\lambda = \pm i \omega$

2. Since $\xi \neq \omega \geq 0$ then $\lambda_{1,2}$ are elements of the negative half plane of the complex plane.

3. Considering LMS methods stability region contains spectrum of G_n

Stability of LMS methods

• Assume G possesses LI $\vec{\psi}$ (assumes diagonalizability)

a) we will reduce to SDOF

let $\lambda = \lambda_g$ $\vec{P} = \vec{P}_g = (\psi_1, \psi_2, \dots, \psi_n)$
 $\vec{\psi} = \vec{\psi}_g$ $\Lambda = \begin{bmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_n \end{bmatrix}$

because of assumption $\vec{P}^{-1} G \vec{P} = \Lambda$

let $\vec{z} = \vec{P}^{-1} \vec{y}$ we will consider $H_n(t) = 0$

then $\dot{\vec{z}} = \Lambda \vec{z} \Rightarrow \dot{z} = \lambda z$ because of diag $\left\{ \begin{array}{l} \vec{P}^{-1} \dot{\vec{y}} = \vec{P}^{-1} G \vec{P} \vec{P}^{-1} \vec{y} \\ \dot{\vec{z}} = \Lambda \vec{z} \end{array} \right.$

Putting into LMS k-step

$$\sum_{i=0}^k (\alpha_i z_{n+1-i} + \Delta t \beta_i \Lambda z_{n+1-i}) = 0$$

for SDOF since diag & $\vec{\psi}$ are LI

$$\sum_{i=0}^k (\alpha_i z_{n+1-i} + \Delta t \beta_i \lambda z_{n+1-i}) = 0$$

these are linear diff eqn.

the solution: $z_{n+1-i} = C \xi^{n+1-i}$. Put into k-step to get

$$\sum_{i=0}^k (\alpha_i + \Delta t \beta_i \lambda) C \xi^{n+1-i} = 0 \quad \text{now look at roots of this polynomial}$$

for stability analysis

Definition 1: An LMS method of the form (2) (ie $\sum_{i=0}^k () = 0$)

is said to be absolutely stable for a fixed $\Delta t \lambda$

if $\forall z$ satisfying $|z| \leq 1$, ^{eqn then} is lies of unit circle.

(similar to spectral stability). There is an equiv between this & spectral stability of A_0 .

Definition 2: Region of absolute stability

is the set of all $\lambda \Delta t \in \mathbb{C}$ at which the LMS is absolutely stable.

Now we can take the k-step method & rewrite it as

$$\begin{Bmatrix} z_{n+1} \\ z_n \\ \vdots \\ z_{n+2-k} \end{Bmatrix} = \underset{k \times k}{\tilde{A}} \begin{Bmatrix} z_n \\ \vdots \\ z_{n+1-k} \end{Bmatrix} \quad \text{where}$$

$$\tilde{A} = \begin{bmatrix} \frac{-(\alpha_1 + \Delta t \lambda \beta_1)}{(\alpha_0 + \Delta t \lambda \beta_0)} & \frac{-(\alpha_2 + \Delta t \lambda \beta_2)}{(\alpha_0 + \Delta t \lambda \beta_0)} & \dots & \frac{-(\alpha_k + \Delta t \lambda \beta_k)}{(\alpha_0 + \Delta t \lambda \beta_0)} \\ 1 & 0 & & 0 \\ 0 & 1 & & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & 1 & 0 \end{bmatrix}$$

looking at $p(\tilde{A})$ its polynomial will be the same as the

δ polynomial we derived before.

Definition 3: A-stable algorithm

An LMS method is A-stable if the solutions of (*) $\rightarrow 0$ as $n \rightarrow \infty$ when $\text{Re } \lambda < 0$

$$(*) \quad \sum_{i=0}^k (\alpha_i + \Delta t \lambda \beta_i) z_{n+1-i} = 0$$

Appendix II - Proof of (7.25)

A preliminary lemma.

Let $\tilde{A}$ and $\tilde{B}$ be symmetric matrices of dimension n , let $\tilde{B}$ be positive definite and consider the eigenvalue problem

$$\tilde{A}\tilde{u} - \lambda\tilde{B}\tilde{u} = \tilde{0}. \quad (\text{II1})$$

There are n real eigenvalues

$$\lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_n, \quad (\text{II2})$$

and corresponding orthonormal (with respect to $\tilde{B}$) eigenvectors $\tilde{u}_1, \dots, \tilde{u}_n$. The Rayleigh quotient is defined by

$$R(\tilde{u}) = \frac{\tilde{u}^T \tilde{A} \tilde{u}}{\tilde{u}^T \tilde{B} \tilde{u}}, \quad \tilde{u} \neq \tilde{0}. \quad (\text{II3})$$

Lemma

$$\lambda_n = \max_{\substack{\tilde{x} \in \mathbb{R}^n \\ \tilde{v} \neq \tilde{0}}} R(\tilde{v})$$

and the maximum occurs for $\tilde{v} = c\tilde{u}_n$, where c is an arbitrary non-zero constant.

Proof. Since the eigenvectors $\tilde{u}_1, \dots, \tilde{u}_n$ form a basis for $\mathbb{R}^n$, we can expand any vector $\tilde{v}$ as

$$\tilde{v} = \sum_{i=1}^n c_i \tilde{u}_i,$$

and, therefore,

$$\tilde{v}^T \tilde{B} \tilde{v} = \sum_{i=1}^n |c_i|^2$$

$$\tilde{v}^T \tilde{A} \tilde{v} = \sum_{i=1}^n \lambda_i |c_i|^2.$$

For $\tilde{v} \neq 0$,

$$R(\tilde{v}) = \frac{\sum_{i=1}^n \lambda_i |c_i|^2}{\sum_{i=1}^n |c_i|^2} \leq \lambda_n, \quad \text{now since } \lambda_n \geq \lambda_i \quad \forall i=1,2,\dots$$

$$\Leftarrow \sum_{i=1}^n \lambda_n |c_i|^2 \geq \sum_{i=1}^n \lambda_i |c_i|^2$$

by (II2). Since $R(c\tilde{u}_n) = \lambda_n$, the proof of the lemma is complete. ■

Let $\tilde{K}$ and $\tilde{M}$ be global matrices formed by assembling the collections of matrices $\{\tilde{k}^e\}_1^{nel}$ and $\{\tilde{m}^e\}_1^{nel}$ according to some connectivity array.

Consider the symmetric, block-diagonal matrices

$$\tilde{\tilde{K}} = \begin{bmatrix} \tilde{k}^1 & 0 & \dots & 0 \\ 0 & \tilde{k}^2 & \ddots & \vdots \\ \vdots & \ddots & \ddots & 0 \\ 0 & \dots & 0 & \tilde{k}^{n_{el}} \end{bmatrix} \quad (II4)$$

$$\tilde{\tilde{M}} = \begin{bmatrix} \tilde{m}^1 & 0 & \dots & 0 \\ 0 & \tilde{m}^2 & \ddots & \vdots \\ \vdots & \ddots & \ddots & 0 \\ 0 & \dots & 0 & \tilde{m}^{n_{el}} \end{bmatrix} \quad (II5)$$

The dimension of $\tilde{\tilde{K}}$ and $\tilde{\tilde{M}}$ is $\tilde{n}_{eq} = \sum_{e=1}^{n_{el}} n_{eq}^e$, where n_{eq}^e denotes the number of equations for element e . If n_{eq} is the dimension of $\tilde{K}$ and $\tilde{M}$ then, clearly, $\tilde{n}_{eq} \geq n_{eq}$.

The eigenvalue problems for the $\tilde{K}$, $\tilde{M}$ and $\tilde{\tilde{K}}$, $\tilde{\tilde{M}}$ systems are (resp.)

$$(\tilde{K} - \lambda \tilde{M}) \tilde{u} = 0 \quad (II6)$$

and

$$(\tilde{\tilde{K}} - \lambda \tilde{\tilde{M}}) \tilde{\tilde{u}} = 0. \quad (II7)$$

All the matrices involved are assumed symmetric and $\underline{\underline{M}}$ and $\underline{\underline{\tilde{M}}}$ are assumed positive-definite. Thus both (II6) and (II7) fit into the framework of (III1) and, therefore, the lemma holds in each case, viz.

$$\lambda_{n_{eq}} = \max_{\substack{\underline{\underline{y}} \in \mathbb{R}^{n_{eq}} \\ \underline{\underline{y}} \neq \underline{\underline{0}}}} \frac{\underline{\underline{v}}^T \underline{\underline{K}} \underline{\underline{v}}}{\underline{\underline{v}}^T \underline{\underline{M}} \underline{\underline{v}}} \quad (II8)$$

$$\tilde{\lambda}_{\tilde{n}_{eq}} = \max_{\substack{\underline{\underline{y}} \in \mathbb{R}^{\tilde{n}_{eq}} \\ \underline{\underline{y}} \neq \underline{\underline{0}}}} \frac{\underline{\underline{v}}^T \underline{\underline{\tilde{K}}} \underline{\underline{v}}}{\underline{\underline{v}}^T \underline{\underline{\tilde{M}}} \underline{\underline{v}}} \quad (II9)$$

Let $\mathbb{R}_{constrained}^{n_{eq}} = \{ \underline{\underline{y}} \in \mathbb{R}^{n_{eq}} \mid \text{components of } \underline{\underline{y}} \text{ which correspond to the same degree-of-freedom in } \mathbb{R}^{n_{eq}} \text{ under the identification of the connectivity array used in assembling } \underline{\underline{K}} \text{ and } \underline{\underline{M}} \text{ are to be equal} \} .$

With this definition it is clear that

$$\lambda_{n_{eq}} = \max_{\substack{\underline{\underline{y}} \in \mathbb{R}_{const.}^{n_{eq}} \\ \underline{\underline{y}} \neq \underline{\underline{0}}}} \frac{\underline{\underline{v}}^T \underline{\underline{\tilde{K}}} \underline{\underline{v}}}{\underline{\underline{v}}^T \underline{\underline{\tilde{M}}} \underline{\underline{v}}} \leq \tilde{\lambda}_{\tilde{n}_{eq}} . \quad (II10)$$

Remarks

1. The equality in (III0) may be intuitively obvious to some readers.

To make it convincing to those readers who are dubious, consider the following simple example:

Let

$$\tilde{k}^e = \begin{bmatrix} k_{11}^e & k_{12}^e \\ k_{21}^e & k_{22}^e \end{bmatrix}$$

$$\tilde{m}^e = \begin{bmatrix} m_{11}^e & m_{12}^e \\ m_{21}^e & m_{22}^e \end{bmatrix}$$

$$e = 1, 2$$

$$\tilde{K} = \begin{bmatrix} k_{11}^1 & k_{12}^1 & 0 \\ k_{21}^1 & (k_{22}^1 + k_{11}^2) & k_{12}^2 \\ 0 & k_{21}^2 & k_{22}^2 \end{bmatrix}$$

$$\tilde{M} = \begin{bmatrix} m_{11}^1 & m_{12}^1 & 0 \\ m_{21}^1 & (m_{22}^1 + m_{12}^2) & m_{12}^2 \\ 0 & m_{21}^2 & m_{22}^2 \end{bmatrix}$$

$$\underbrace{\tilde{K}}_{4 \times 4} = \begin{bmatrix} \tilde{k}^1 & 0 \\ 0 & \tilde{k}^2 \end{bmatrix}$$

$$\underbrace{\tilde{M}}_{4 \times 4} = \begin{bmatrix} \tilde{m}^1 & 0 \\ 0 & \tilde{m}^2 \end{bmatrix}$$

$$\tilde{\varphi} = \begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \varphi_3 \end{pmatrix}$$

$$\tilde{\varphi} = \begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \varphi_2 \\ \varphi_3 \end{pmatrix} \in \mathbb{R}_{\text{const.}}^4$$

The equality in (II10) is established for the example under consideration by the following calculations:

$$\begin{aligned} \tilde{\varphi}^T \tilde{K} \tilde{\varphi} &= \begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix}^T \tilde{k}^1 \begin{pmatrix} \varphi_1 \\ \varphi_2 \end{pmatrix} + \begin{pmatrix} \varphi_2 \\ \varphi_3 \end{pmatrix}^T \tilde{k}^2 \begin{pmatrix} \varphi_2 \\ \varphi_3 \end{pmatrix} \\ &= \tilde{\varphi}^T \tilde{K} \tilde{\varphi} \end{aligned}$$

$$\begin{aligned}
& \tilde{\varphi}^T \tilde{M} \tilde{\varphi} = \begin{Bmatrix} \varphi_1 \\ \varphi_2 \end{Bmatrix}^T \tilde{m}_1 \begin{Bmatrix} \varphi_1 \\ \varphi_2 \end{Bmatrix} + \begin{Bmatrix} \varphi_2 \\ \varphi_3 \end{Bmatrix}^T \tilde{m}_2 \begin{Bmatrix} \varphi_2 \\ \varphi_3 \end{Bmatrix} \\
& = \tilde{\varphi}^T \tilde{M} \tilde{\varphi}
\end{aligned}$$

The essence of a formal proof may be deduced from the above calculations. The interested reader may wish to provide the details.

2. The inequality in (II10) follows directly from (II9).

In terms of the element matrices, the eigenvalue problem (II7) can be written as

$$\begin{bmatrix} \tilde{k}^1 - \tilde{\lambda}_m^1 & 0 & \dots & 0 \\ 0 & \tilde{k}^2 - \tilde{\lambda}_m^2 & & \vdots \\ \vdots & & \ddots & 0 \\ 0 & \dots & 0 & \tilde{k}^{n_{el}} - \tilde{\lambda}_m^{n_{el}} \end{bmatrix} \begin{Bmatrix} \tilde{u}^1 \\ \tilde{u}^2 \\ \vdots \\ \tilde{u}^{n_{el}} \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ \vdots \\ 0 \end{Bmatrix}$$

(II11)

From this expression it can be seen that the eigenvalues of (II7) are the same as the totality of eigenvalues of the uncoupled element problems

$$(\tilde{k}^e - \lambda_{\tilde{m}}^e) \tilde{u}^e = \tilde{0}, \quad e = 1, \dots, n_{el} \quad (\text{II12})$$

Therefore

$$\tilde{\lambda}_{n_{eq}}^e = \max_e (\lambda_{\max}^e), \quad (\text{II13})$$

where

$$\lambda_{\max}^e = \max_{\substack{\tilde{v} \in \mathbb{R}^{n_{eq}} \\ \tilde{v} \neq 0}} \frac{\tilde{v}^T \tilde{k}^e \tilde{v}}{\tilde{v}^T \tilde{m}^e \tilde{v}}.$$

Combining (II10) and (II13), we have that

$$\lambda_{n_{eq}} \leq \max_e (\lambda_{\max}^e),$$

which completes the proof of the theorem. $\blacksquare$

FINITE ELEMENT NOTES

1) LINEAR MULTISTEP METHODS FOR FIRST - ORDER EQUATION SYSTEMS

CONSIDER A SYSTEM OF LINEAR, ORDINARY DIFFERENTIAL EQUATIONS

$$\dot{\underline{y}} = \underline{f}(\underline{y}, t) = \underline{G} \underline{y} + \underline{H}(t) \quad (1)$$

WHERE

$\underline{G} \dots$ IS A GIVEN CONSTANT MATRIX

$\underline{H} \dots$ IS A GIVEN VECTOR-VALUED FUNCTION OF t

DEFINITION: A k -STEP LINEAR MULTISTEP METHOD FOR (1) IS DEFINED BY

$$\sum_{i=0}^k \{ \alpha_i \underline{y}_{n+1-i} + \Delta t \beta_i \underline{f}(\underline{y}_{n+1-i}, t_{n+1-i}) \} = 0 \quad (2)$$

WHERE THE α_i 's AND β_i 's ARE GIVEN PARAMETERS WHICH DEFINE THE METHOD

NOTE: THE TERM "LINEAR" IN LINEAR MULTISTEP METHOD REFERS TO THE FORM OF (2) NOT TO THE LINEARITY OF $\underline{f}$ AS GIVEN ABOVE. IN FACT (2) HOLDS FOR $\underline{f}$ NONLINEAR.

REMARKS:

i) AN LMS METHOD IS CALLED EXPLICIT IF $\beta_0 = 0$ OTHERWISE IT IS CALLED IMPLICIT.

ii) AN LMS METHOD IS CALLED A BACKWARD-DIFFERENCE METHOD IF $\beta_i = 0$ FOR ALL $i \geq 1$

$$\alpha_0 \underline{y}_{n+1} + \Delta t \beta_0 \underline{f}(\underline{y}_{n+1}, t_{n+1}) + \alpha_1 \underline{y}_n + \alpha_2 \underline{y}_{n-1} + \dots + \alpha_k \underline{y}_{n+1-k} = 0$$

then $\underline{y}_{n+1}$ is dependent only on all $\underline{y}_i$ $i \leq n+1$

(iii) AN EXAMPLE OF A 1-STEP LMS METHOD IS THE GENERALIZED TRAPEZOIDAL FAMILY, FOR WHICH

$$h=1 \quad \alpha_0 = 1 \quad \beta_0 = -\alpha \quad \Rightarrow y_{n+1} = \alpha y_n + \Delta t f(y_n, t_n)$$

$$\alpha_1 = -1 \quad \beta_1 = \alpha - 1 \quad \Rightarrow y_{n+1} = \alpha y_n + \Delta t f(y_n, t_n) = y_n$$

$$\alpha_0 = 1, \alpha_1 = -1, \beta_0 = 0, \beta_1 = -1$$

RECALL $\alpha = 0 \Rightarrow$ EXPLICIT $\Rightarrow y_{n+1} = y_n + \Delta t f(y_n, t_n)$

$\alpha = 1 \Rightarrow$ BACKWARD DIFFERENCE $y_{n+1} - \Delta t f(y_{n+1}, t_{n+1}) = y_n$

$$\alpha_0 = 1, \alpha_1 = -1, \beta_0 = -1, \beta_1 = 0$$

(v) THE GREATER THE NUMBER OF STEPS INVOLVED ("h")

THE GREATER THE "HISTORICAL DATA POOL" WHICH MUST BE STORED $\rightarrow$ A PRACTICAL DISADVANTAGE!

2) SEMI-DISCRETE HEAT EQUATION AND EQUATION OF MOTION IN FIRST-ORDER FORM

IN ORDER TO STUDY LMS METHODS AS APPLIED TO THE HEAT EQUATION AND THE EQUATION OF MOTION, WE NEED TO FIRST PUT THESE EQUATIONS IN THE FORM OF (1), AND THEN STUDY THE SPECTRAL PROPERTIES OF THE RESULTING G MATRICES.

Q) SEMI-DISCRETE HEAT EQUATION

GIVEN :

$$M \dot{d} + K d = F$$

$$d(0) = D_0$$

$$\dot{y} = G y$$

$$\text{let } y = \psi e^{\lambda t} \quad \text{for homogeneous case}$$

$$\dot{y} = \lambda \psi e^{\lambda t}$$

$$\Rightarrow \lambda \psi e^{\lambda t} = G \psi e^{\lambda t}$$

$$\Rightarrow (\lambda I - G) \psi e^{\lambda t} = 0$$

$$\Rightarrow \text{ev problem is find } \lambda \text{ s.t. } [\lambda I - G] = 0$$

WHERE :

M ... CAPACITY MATRIX, SYMM. POS. DEF.

K ... CONDUCTIVITY MATRIX, SYMM. POS. SEMI-DEF.

$F = F(t), t \in [0, T] \dots$ HEAT SUPPLY VECTOR (PRESC.)

$d \dots$ TEMPERATURE VECTOR

$\dot{d}$... TIME (t) DERIVATIVE OF d

D_0 ... GIVEN INITIAL TEMPERATURE.

LETTING $\underline{u} = \underline{d}$, WE FIND (3)

$$\underline{Q} = -\underline{M}^{-1} \underline{K} \quad (4)$$

$$\underline{H}(t) = \underline{M}^{-1} \underline{F}(t) \quad (5)$$

THE SPECTRUM OF $\underline{Q}$ IS DETERMINED FROM THE FOLLOWING EIGENPROBLEM

$$(\underline{Q} - \lambda(\underline{Q}) \underline{I}) \underline{\Psi}(\underline{Q}) = \underline{Q} \quad (6)$$

SUBSTITUTING (4) AND MULTIPLYING BY $\underline{M}$ YIELDS

$$(\underline{K} + \lambda(\underline{Q}) \underline{M}) \underline{\Psi}(\underline{Q}) = \underline{Q} \quad (7)$$

COMPARISON WITH PREVIOUS RESULTS (I6) YIELDS

$$\lambda(\underline{Q}) = -\lambda \leq 0 \quad (8)$$

$$\underline{\Psi}(\underline{Q}) = \underline{\Psi} \quad (9)$$

THUS IN THIS CASE THE SPECTRUM OF $\underline{Q}$ ($\lambda(\underline{Q})$) RESIDES UPON THE NEGATIVE REAL AXIS IN THE COMPLEX PLANE, $\mathbb{C}$.

b) SEMI-DISCRETE EQUATIONS OF MOTION

GIVEN :

$$\underline{M} \ddot{\underline{d}} + \underline{C} \dot{\underline{d}} + \underline{K} \underline{d} = \underline{F}$$

WHERE :

let $\dot{d} = v$ $d = u$

$$\underline{M} \dot{\underline{v}} + \underline{C} \underline{u} + \underline{K} \underline{u} = \underline{F}$$

$$\dot{\underline{v}} = \underline{M}^{-1} \underline{F} - \underline{M}^{-1} \underline{C} \underline{u} - \underline{M}^{-1} \underline{K} \underline{u}$$

$$\begin{pmatrix} \dot{\underline{u}} \\ \dot{\underline{v}} \end{pmatrix} = \begin{bmatrix} \underline{0} & \underline{I} \\ -\underline{M}^{-1} \underline{K} & -\underline{M}^{-1} \underline{C} \end{bmatrix} \begin{pmatrix} \underline{u} \\ \underline{v} \end{pmatrix} + \begin{pmatrix} \underline{0} \\ \underline{M}^{-1} \underline{F} \end{pmatrix}$$

$\underline{M}$... MASS MATRIX, SYMM. POS. DEF.

$\underline{C}$... DAMPING MATRIX, SYMM, POS. SEMI-DEF.

$\underline{K}$... STIFFNESS MATRIX, SYMM. POS. SEMI-DEF.

$\underline{F} = \underline{F}(t)$... EXTERNAL FORCE VECTOR (PRESC.)

$\underline{d}$... DISPLACEMENT VECTOR.

100-202-00
100-202-00
100-202-00
(100-202-00)

$\dot{\underline{q}} \dots$ VELOCITY VECTOR (DERIV. OF $\underline{q}$ WRT t)

$\ddot{\underline{q}} \dots$ ACCELERATION VECTOR (2^{nd} DERIV. OF $\underline{q}$ WRT t)

IN THIS CASE WE LET

$$\underline{y} = \begin{Bmatrix} \underline{q} \\ \dot{\underline{q}} \end{Bmatrix} \quad (10)$$

THUS

$$\underline{G} = \begin{bmatrix} \underline{0} & \underline{I} \\ -\underline{M}^{-1}\underline{K} & -\underline{M}^{-1}\underline{C} \end{bmatrix}, \quad \underline{H}(t) = \begin{Bmatrix} \underline{0} \\ \underline{M}^{-1}\underline{F}(t) \end{Bmatrix} \quad (11), (12)$$

IF WE LET

$$\underline{\psi}(\underline{\omega}) = \begin{Bmatrix} \psi_1(\underline{\omega}) \\ \psi_2(\underline{\omega}) \end{Bmatrix} \quad \text{THEN THE EIGENPROBLEM CAN BE EXPRESSED AS} \quad (13)$$

$$\left(\begin{bmatrix} \underline{0} & \underline{I} \\ -\underline{M}^{-1}\underline{K} & -\underline{M}^{-1}\underline{C} \end{bmatrix} - \begin{bmatrix} \lambda(\underline{\omega})\underline{I} & \underline{0} \\ \underline{0} & \lambda(\underline{\omega})\underline{I} \end{bmatrix} \right) \begin{Bmatrix} \psi_1(\underline{\omega}) \\ \psi_2(\underline{\omega}) \end{Bmatrix} = \begin{Bmatrix} \underline{0} \\ \underline{0} \end{Bmatrix} \quad (14)$$

OR EQUIVALENTLY,

$$\underline{\psi}_2(\underline{\omega}) - \lambda(\underline{\omega}) \underline{\psi}_1(\underline{\omega}) = \underline{0} \quad (15)$$

$$-\underline{M}^{-1}\underline{K} \underline{\psi}_1(\underline{\omega}) - (\underline{M}^{-1}\underline{C} + \lambda(\underline{\omega})\underline{I}) \underline{\psi}_2(\underline{\omega}) = \underline{0} \quad (16)$$

COMBINING (15) AND (16) AND MULTIPLYING BY $\underline{M}$ YIELDS A QUADRATIC EIGENPROBLEM:

$$(\underline{K} + \lambda(\underline{\omega}) \underline{C} + \lambda(\underline{\omega})^2 \underline{M}) \underline{\psi}_1(\underline{\omega}) = \underline{0} \quad (17)$$

FROM PREVIOUS WORK WE KNOW THAT THERE EXIST A SET OF VECTORS $\hat{\underline{\psi}}_l$ SUCH THAT

$$\hat{\underline{\psi}}_l^T \underline{M} \hat{\underline{\psi}}_m = \delta_{lm}$$

$$\hat{\underline{\psi}}_l^T \underline{K} \hat{\underline{\psi}}_m = \delta_{lm} \omega_l^2 \quad (\text{NO SUM})$$

NOW IF WE ASSUME THAT ζ IS DIAGONAL IN THE $\hat{\psi}$ BASIS (WHICH IS THE CASE FOR RAYLEIGH DAMPING WHERE $\zeta = aM + bK$ a, b CONSTANTS) THEN

$$\hat{\psi}_l^T \zeta \hat{\psi}_m = 2 \xi_l \omega_l \delta_{lm} \quad (\text{NO SUM})$$

THUS PREMULTIPLYING (17) BY $\hat{\psi}_l^T$ AND EXPANDING $\hat{\psi}(t)$ IN TERMS OF THE $\hat{\psi}_l$ WE FIND

$$\lambda(\zeta)^2 + 2 \xi_l \omega_l \lambda(\zeta) + \omega_l^2 = 0 \quad (\text{NO SUM})$$

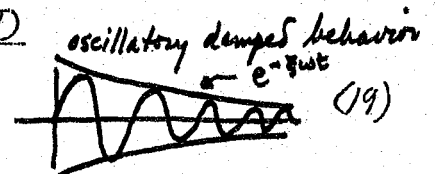
DROPPING THE l SUBSCRIPT WE FIND THE UNDERLYING FORM OF THE $\lambda(\zeta)$ EIGENPROBLEM TO BE

$$\lambda(\zeta)^2 + 2 \xi \omega \lambda(\zeta) + \omega^2 = 0 \quad (18)$$

THE SOLUTION OF (18) TAKES ON THE FOLLOWING CHARACTERISTIC FORMS:

(i) $0 \leq \xi < 1 \Rightarrow$ UNDERDAMPED

$$\lambda_{1,2}(\zeta) = -\xi \omega \pm i \omega \sqrt{1 - \xi^2}$$



(19)

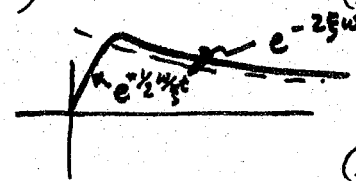
(ii) $\xi = 1 \Rightarrow$ CRITICALLY DAMPED

$$\lambda_{1,2}(\zeta) = -\omega \quad (\text{DOUBLE ROOT})$$

(20)

(iii) $\xi > 1 \Rightarrow$ OVERDAMPED

$$\lambda_{1,2}(\zeta) = -\xi \omega \pm \omega \sqrt{\xi^2 - 1}$$



(21)

REMARKS:

i) NO DAMPING $\Rightarrow \xi = 0 \Rightarrow \lambda_{1,2}(\zeta) = \pm i \omega$ (CONJUGATE IMAGINARY ROOTS).

ii) IN ALL CASES, SINCE ξ AND $\omega \geq 0$ THE EIGENVALUES OF ζ ARE CONFINED TO THE NEGATIVE HALF-PLANE OF $\mathbb{C}$ INCLUDING THE IMAGINARY AXIS.

iii) IN CONSIDERING LMS METHODS IT IS IMPORTANT TO INSURE THAT THE STABILITY REGION IN THE COMPLEX PLANE ENCLOSES THE SPECTRUM OF THE MATRIX ζ FOR THE CASE IN QUESTION.

3) STABILITY OF LMS METHODS

a) REDUCTION TO AN SDOF PROBLEM

AS WITH PREVIOUS ANALYSES THE REDUCTION TO A MODEL PROBLEM PLAYS AN ESSENTIAL ROLE. WE WILL ASSUME THAT $\underline{G}$ POSSESSES LINEARLY INDEPENDENT EIGENVECTORS.

NOTATION:

$\lambda = \lambda(\underline{G}) \dots$ EIGENVALUE

$\psi = \psi(\underline{G}) \dots$ EIGENVECTOR

$\underline{P} = \underline{P}(\underline{G}) \dots$ MATRIX OF EIGENVECTORS

$\underline{\Lambda} = \underline{\Lambda}(\underline{G}) \dots$ DIAGONAL MATRIX OF EIGENVALUES

THUS

$$\underline{P} = [\psi_1, \psi_2, \dots, \psi_N]$$

$$\underline{\Lambda} = \text{diag}[\lambda_1, \lambda_2, \dots, \lambda_N] = \begin{bmatrix} \lambda_1 & & \\ & \lambda_2 & \\ & & \ddots \\ & & & \lambda_N \end{bmatrix}$$

WHERE N IS THE DIMENSION OF $\underline{G}$

THUS $\underline{P}$ AND $\underline{\Lambda}$ SATISFY

$$\underline{P}^{-1} \underline{G} \underline{P} = \underline{\Lambda} \quad (22)$$

DEFINING

$$\underline{z} = \underline{P}^{-1} \underline{y} \quad (23)$$

ENABLES US TO OBTAIN FROM THE HOMOGENOUS FORM OF (1), THE FOLLOWING UNCOUPLED SYSTEM

$$\dot{\underline{z}} = \underline{\Lambda} \underline{z} \quad (24)$$

WHICH IS CHARACTERIZED BY THE GENERIC SCALAR EQUATION

$$P^{-1} \cdot e_{q,n}(z) = \sum_{i=0}^k (\underbrace{\alpha_i P^{-1}}_{\alpha_i} \underbrace{y_{n+1-i}}_{y_{n+1-i}} + \Delta t \beta_i \underbrace{P^{-1} G}_{\underbrace{P^{-1} \lambda}_{\lambda}} \underbrace{y_{n+1-i}}_{y_{n+1-i}}) \quad (25)$$

$$\dot{z} = \lambda z$$

LIKEWISE THE HOMOGENEOUS FORM OF THE LMS ALGORITHM (2), CAN BE UNCOUPLED

$$\sum_{i=0}^k (\alpha_i z_{n+1-i} + \Delta t \beta_i \lambda z_{n+1-i}) = 0 \quad (26)$$

WHICH IS CHARACTERIZED BY

$$\sum_{i=0}^k (\alpha_i + \Delta t \lambda \beta_i) z_{n+1-i} = 0 \quad (27)$$

TO DETERMINE STABILITY ONE ASSUMES THE SOLUTION OF THE LINEAR DIFFERENCE EQUATION (27) TO BE OF THE FORM

$$z_{n+1-i} = C \zeta^{n+1-i}, \quad C \dots \text{A CONSTANT}$$

SUBSTITUTION INTO (27) YIELDS A k^{TH} ORDER POLYNOMIAL IN ζ , THE ROOTS OF WHICH DETERMINE THE STABILITY OF THE LMS METHOD.

if $k=1$

$$\sum_{i=0}^k (\alpha_i + \Delta t \lambda \beta_i) \zeta^{n+1-i} = 0 \quad (28)$$

$$(\alpha_0 + \Delta t \lambda \beta_0) \zeta^{n+1} + (\alpha_1 + \Delta t \lambda \beta_1) \zeta^n = 0 \Rightarrow \zeta^n [(\alpha_1 + \Delta t \lambda \beta_1) + \zeta (\alpha_0 + \Delta t \lambda \beta_0)] = 0$$

b) CONCEPTS OF STABILITY or $\zeta=0$ (n roots) or $\zeta = -\frac{\alpha_1 + \Delta t \lambda \beta_1}{\alpha_0 + \Delta t \lambda \beta_0}$
 explicit $\zeta = \frac{1 + \Delta t \lambda}{1 - \Delta t \lambda}$; implicit $\zeta = \frac{1}{1 - \lambda \Delta t}$

DEFINITION 1: A LMS METHOD OF THE FORM (2) IS SAID TO BE ABSOLUTELY STABLE, FOR A FIXED $\Delta t \lambda$ IF ALL ζ SATISFYING (28) ARE SUCH THAT $|\zeta| \leq 1$
 for explicit $\Rightarrow$ for $\lambda < 0$ any Δt will do; for BD $+\lambda \Delta t \leq 0$

DEFINITION 2: THE REGION OF ABSOLUTE STABILITY OF A LMS METHOD IS THE SET OF $\Delta t \lambda \in \mathbb{C}$ AT WHICH IT IS ABSOLUTELY STABLE.

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REMARK: THE DEFINITION OF ABSOLUTE STABILITY IS VERY CLOSE TO THAT OF SPECTRAL STABILITY GIVEN PREVIOUSLY, WITH THE EXCEPTION THAT NO SPECIAL ACCOUNT IS GIVEN TO MULTIPLE ROOTS OF UNIT MODULUS, WHICH POTENTIALLY GIVE RISE TO WEAK INSTABILITIES. ONE CAN SHOW THAT THE ROOTS OF (28) AND THE EIGENVALUES OF AN ASSOCIATED AMPLIFICATION MATRIX ARE THE SAME BY CONSTRUCTING A ONE-STEP MULTIVALUE METHOD EQUIVALENT TO THE GIVEN LMS METHOD, VIZ.

$$\underbrace{\begin{Bmatrix} z_{n+1} \\ z_n \\ \vdots \\ z_{n+2-k} \end{Bmatrix}}_{k \times 1} = \underbrace{\tilde{A}}_{k \times k} \underbrace{\begin{Bmatrix} z_n \\ z_{n-1} \\ \vdots \\ z_{n+1-k} \end{Bmatrix}}_{k \times 1} \quad (29)$$

WHERE

$$\tilde{A} = \begin{bmatrix} \frac{-(\alpha_1 + \Delta t \lambda \beta_1)}{(\alpha_0 + \Delta t \lambda \beta_0)} & \frac{-(\alpha_2 + \Delta t \lambda \beta_2)}{(\alpha_0 + \Delta t \lambda \beta_0)} & \dots & \dots & \frac{-(\alpha_k + \Delta t \lambda \beta_k)}{(\alpha_0 + \Delta t \lambda \beta_0)} \\ 1 & 0 & \dots & 0 & 0 \\ 0 & 1 & & 0 & \vdots \\ \vdots & & \ddots & \vdots & \vdots \\ 0 & 0 & \dots & 1 & 0 \end{bmatrix} \quad (30)$$

DEFINITION 3: A LMS METHOD IS A-STABLE IF SOLUTIONS OF (27) $\rightarrow 0$ AS $n \rightarrow \infty$ WHEN $\text{Re } \lambda < 0$.

REMARKS:

i) PHYSICALLY SPEAKING, AN A-STABLE ALGORITHM IS ONE WHICH PRODUCES SOLUTIONS WHICH DECAY TO ZERO WHENEVER THE CORRESPONDING EXACT SOLUTIONS OF (25) DECAY TO ZERO. CLEARLY IF AN ALGORITHM IS A-STABLE THEN THE REGION OF ABSOLUTE STABILITY CONTAINS THE LEFT HALF-PLANE OF $\mathbb{C}$.

ii) THE CONDITION OF A-STABILITY PLACES NO LIMITATION ON THE SIZE OF Δt , CONSEQUENTLY IT IS CLOSELY

ie look at trapezoidal algorithm no condition on Δt exists except > 0

RELATED TO WHAT WE PREVIOUSLY TERMED UNCONDITIONAL SPECTRAL STABILITY. THE ONLY DIFFERENCE BEING THAT SPECTRAL STABILITY ALLOWS EIGENVALUES OF UNIT MODULUS WHICH POTENTIALLY COULD CONSERVE A SOLUTION RATHER THAN CONTRACT IT TO ZERO.

IT SHOULD BE APPARENT THAT A-STABLE ALGORITHMS ARE IMPORTANT FOR MANY PHYSICAL PROBLEM CLASSES. HOWEVER, WITHIN THE CLASS OF LMS METHODS, THE SUBCLASS OF A-STABLE METHODS IS SEVERELY DELIMITED. THIS IS THE CONTENT OF A CELEBRATED THEOREM DUE TO DAHLQUIST.

DAHLQUIST'S THEOREM ON A-STABLE LMS METHODS

- 1) AN EXPLICIT A-STABLE LMS METHOD DOES NOT EXIST.
- 2) A THIRD-ORDER ACCURATE A-STABLE LMS METHOD DOES NOT EXIST.
- 3) THE SECOND-ORDER ACCURATE A-STABLE LMS METHOD WITH THE SMALLEST ERROR CONSTANT IS THE TRAPEZOIDAL RULE. (RECALL THAT THE ERROR CONSTANT IS THE COEFFICIENT OF Δt^{k+1} IN THE LOCAL TRUNCATION ERROR.)

REMARKS:

- i) THE UPSHOT OF THIS THEOREM IS THAT IF WE SEEK AN A-STABLE LMS METHOD WHICH POSSESSES SOME SPECIAL FEATURE (SUCH AS HIGH-FREQUENCY NUMERICAL DISSIPATION) WE NECESSARILY ENTAIL SOME LOSS OF ACCURACY WITH RESPECT TO THE TRAPEZOIDAL RULE.
- ii) SINCE THE SPECTRUM OF THE G-MATRIX FOR THE EQUATION OF MOTION FALLS WITHIN THE LEFT-HALF OF THE COMPLEX PLANE (AND FOR THE UNDAMPED CASE FALLS ON THE IMAGINARY AXIS) IT WOULD APPEAR THAT THE TRAPEZOIDAL RULE IS THE CANONICAL A-STABLE LMS METHOD FOR STRUCTURAL DYNAMICS. THE WIDESPREAD USE OF THE TRAPEZOIDAL RULE IN THIS CONTEXT APPEARS TO CONFIRM THIS OBSERVATION.
- iii) THERE ARE A-STABLE LMS METHODS OTHER THAN TRAPEZOIDAL RULE WHICH ARE USEFUL FOR STRUCTURAL DYNAMICS (FOR EXAMPLE PARKS METHOD)

(v) MANY STRUCTURAL DYNAMICS ALGORITHMS DO NOT FALL WITHIN THE CLASS OF FIRST-ORDER LMS METHODS CONSIDERED THUS FAR. (FOR EXAMPLE THE NEWMARK FAMILY OF ALGORITHMS)

REFERENCES :

- (D1) G. DAHLQUIST, "A SPECIAL STABILITY PROBLEM FOR LINEAR MULTISTEP METHODS", BIT 3, 27-43 [1963].
- (G1) C. W. GEAR, NUMERICAL INITIAL VALUE PROBLEMS IN ORDINARY DIFFERENTIAL EQUATIONS, PRENTICE-HALL, ENGLEWOOD CLIFFS, N. J., [1971]

Remarks

- If an algorithm is A-stable region of absolute stability contains left half of complex $\mathbb{C}$ plane
- Condition of A-stability places no restrictions on Δt (similar to unconditional spectral stability - however it allows $|\lambda| = 1$).

A-stable places severe restrictions on physical systems.

Dahlquist's Theorem

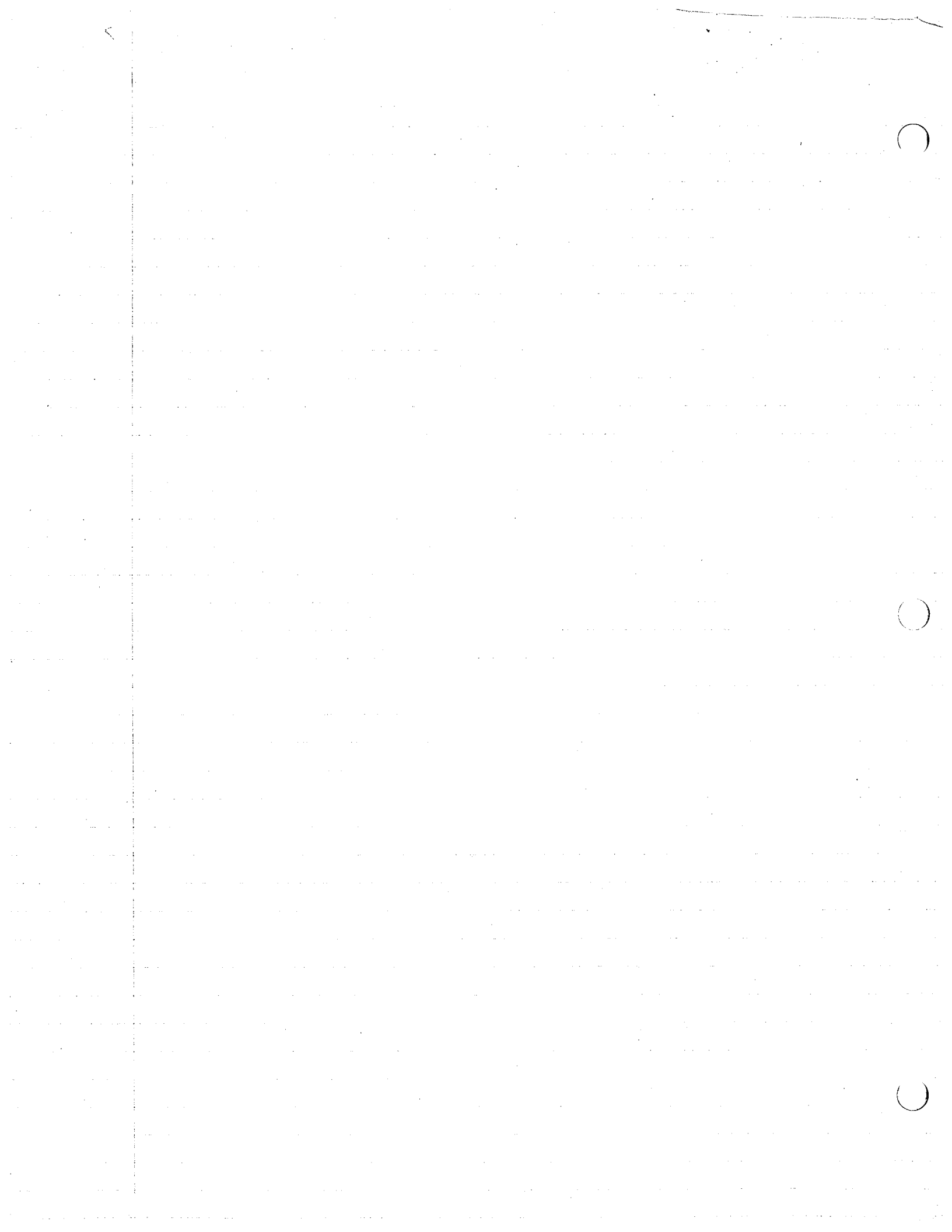
- An explicit A-stable LMS method doesn't exist.
 - That a 3rd order accurate LMS method does not exist (A-stable)
 - The 2nd order accurate LMS (A-stable) method w/ best error constant trapezoidal rule
- remember $\|e\| \leq C \Delta t^{k+1}$

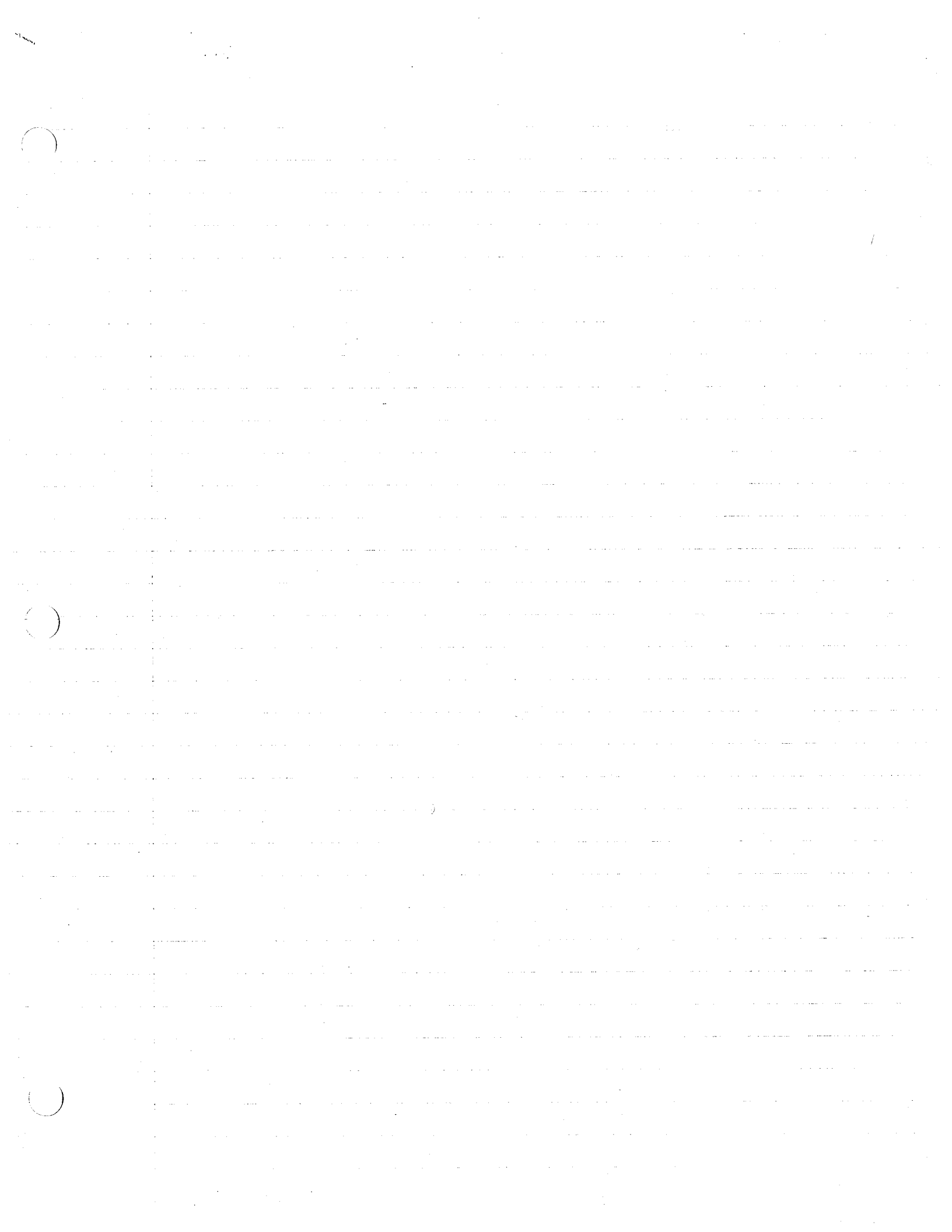
Remarks

1. A Special LMS method that has A-stability & desirable effect (i.e. high frequency numerical dissipation) will be less accurate than trapezoidal
2. Since spec of G is in left hand plane, trapez rule is canonical rule.
i.e. For eqn of motion trapezoidal rule is canonical LMS A-stable method.
3. Other LMS methods, that are used extensively for eqns of motion (Park's method - keeps more historical data - more computational).

→ * Newmark cannot be used into 1st order system

4. There are other integ. algorithm for structural dynamics (i.e. Newmark's)







HOMEWORK

CESAR LEVY ✓

Consider the equations of two-dimensional, classical, linear isotropic, elasticity theory on the domain illustrated in Figure 1

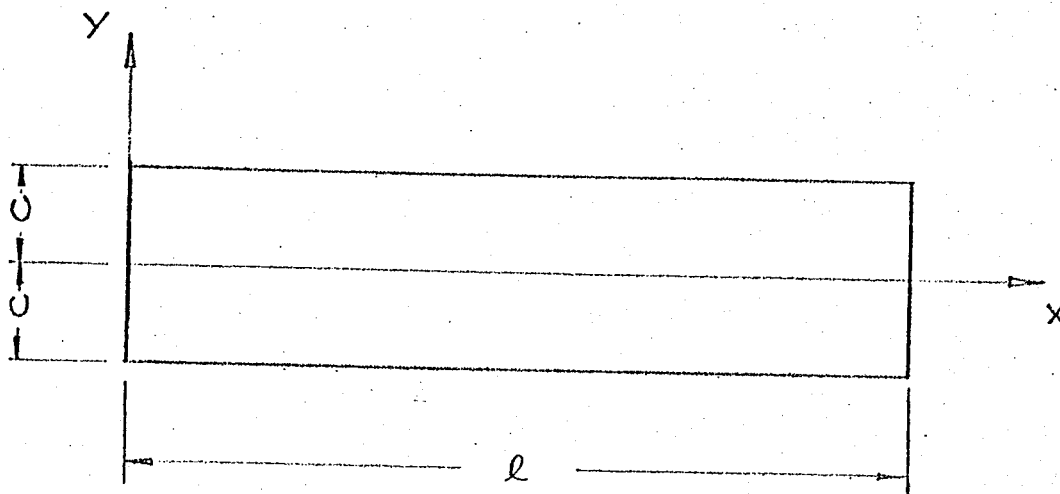


Figure 1

Assume there is no body force, and the boundary conditions are given as follows:

(displacement)

$$u_1(0,0) = u_2(0,0) = 0$$

$$u_1(0,\pm c) = 0$$

(traction)

$$t_1(x,\pm c) = t_2(x,\pm c) = 0, \quad x \in]0, l[$$

The cross diag and bisect patterns were found not to give good results for $\nu=3$ $1 \times 1 \lambda, \mu$ w/hourglass / 2×2 w/compatible modes give same results

$u_2(l,0)$: actual/num is $-244.14 / -242.98$

$u_1(0,y)$	actual/num
$y=0$	0 / 0
.5	-.1295 / -.1372
1.0	-.2072 / -.2207
1.5	-.1813 / -.1953
2.0	0 / 0

$\sigma_{11}(7,y)$ actual/num

y	actual/num
0.25	.4299 / .4218
.75	1.266 / 1.266
1.25	2.109 / 2.110
1.75	2.953 / 2.953

$\sigma_{12}(7,y)$ actual/num

y	actual/num
0.25	-.369 / -.367
.75	-.3223 / -.3204
1.25	-.2285 / -.2286
1.75	-.0879 / -.08577

2×2 quad gave large errors in $u_1(0,y)$ $8-10\%$ in σ_{11} $3 \sim 14\%$ in τ_{12}
 $1 \times 1 \lambda$ $2 \times 2 \mu$ quad " " in $u_1(0,y)$ $8-10\%$ in σ_{11} $3 \sim 14\%$ in τ_{12}

$\nu=.499$ $1 \times 1 \lambda, \mu$ w/hourglass & 2×2 w/comp

$u_2(l,0)$: actual/num is $-205.47 / -204.11$

$u_1(0,y)$	act/num
$y=0$	0 / 0
.5	-.1317 / -.1466
1.0	-.2807 / -.2369
1.5	-.7843 / -.2117
2.0	0 / 0

$\sigma_{11}(7,y)$ actual/num

y	actual/num
0.25	.4219 / .4227
0.75	1.266 / 1.264
1.25	2.109 / 2.112
1.75	2.953 / 2.952

$\tau_{12}(7,y)$

-.369 / -.367
-.3223 / -.3205
-.2285 / -.2268
-.0879 / -.0889

$1 \times 1 \lambda$ $2 \times 2 \mu$ gave large errors in $u_1(0, y)$ 5% in σ_{11} 1-6% in σ_{12}
 2×2 quad 67% in $u_2(1, 0)$ large errors in $u_2(0, y)$ large errors in σ_{11}

would recommend quad overall w/ 2×2 w/incomp
or 1×1 w/hourglass

	QUAD	BISECT	CROSS DIAG	
even worse 1 x 2 μ WORSE THAN (++)	$V = .3$ 2001. $V = .499$ 2003.	not good very large very bad σ_{12} has wrong sign $V = .499$ 2012. $V = .3$ 2011.	$V = .499$ 2017. $V = .3$ 2016.	(like 1 $\lambda \mu$ + hour) not good 2 x 2
Same as Thompson V.G. Same as #	$V = .3$ 2005. $V = .499$ 2008.	like the above. $V = .499$ 2020.	$V = .499$ 2019.	(like 1 $\lambda \mu$ + hour) 2 x 2 w/ uncomp
not good comp diag NOT GOOD (++)	$V = .3$ 2004. $V = .499$ 2007.	$V = .499$ 2014. like the above like the above	$V = .499$ 2022 not good errors in $u_1 \sim 150\%$ in σ_{11} & σ_{12} high (++)	1 x 1 λ 2 x 2 μ
very good on σ_{11} V.G.	$V = .3$ 2006. $V = .499$ 2009.	$V = .499$ 2018. like the above not good	$V = .499$ 2018. same just as bad as $\lambda 1$ & $\mu 2$ w/ σ_{12} changing sign also	1 x 1 w/ hourglass
	2000 data check 2002 2028 2010 purged.	2010 data check	2015 data check	

DATE: May 6, 1981

TO : Students in Applied Mechanics

FROM : Tom Hughes and Charles Steele

SUBJECT: Applied Mechanics Industrial Affiliates Dinner

You are invited for cocktails and dinner on Wednesday evening at 6:00 p.m. at the Stanford Faculty Club you are strongly urged also to attend the presentations on Wednesday and Thursday, May 13, 14, 1981. Please sign up with Theo, Room 252, Durand Bldg. by Friday MAY 8, 1981

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$$\left. \begin{aligned} t_1(\ell, y) &= 0 \\ t_2(\ell, y) &= \frac{P}{2I} (c^2 - y^2) \\ t_1(0, y) &= \int_0^+ \frac{P\ell y}{I} \\ t_2(0, y) &= -\frac{P}{2I} (c^2 - y^2) \end{aligned} \right\} \begin{aligned} &y \in]-c, c[\\ &y \in]-c, 0[\cup]0, c[\end{aligned}$$

where P is a given constant and $I = 2c^3/3$.

The traction boundary conditions are those encountered in simple bending theory for a cantilever beam with root section at $x = 0$, that is, parabolically varying end shear and linearly varying bending stress at the root. The displacement boundary conditions allow the root section to warp.

The exact solution of this problem is easily derived and is given as follows:

$$t_{11} = -\frac{P\tilde{x}y}{I}$$

$$t_{22} = 0$$

$$t_{12} = \frac{P}{2I} (c^2 - y^2)$$

$$\frac{6E^1 I}{P} u_1(x, y) = -y\{3(\ell^2 - \tilde{x}^2) + (2 + \nu^1)(y^2 - c^2)\}$$

$$\frac{6E^1 I}{P} u_2(x, y) = (\tilde{x}^3 - \ell^3) - \{(4 + 5\nu^1)c^2 + 3\ell^2\}(\tilde{x} - \ell) + 3\nu^1 \tilde{x}y^2$$

where $\tilde{x} = \ell - x$ and

	E^1	ν^1
plane stress	E	ν
plane strain	$E/(1-\nu^2)$	$\nu/(1-\nu)$

Obtain approximate solutions of the above elasticity problem by using the linear static analysis program "LEARN."

- Employ the following data in your calculations:

$$P = -1, \quad l = 16, \quad c = 2, \quad E = 1, \quad \nu = .3 \text{ and } .499$$

(The latter value of Poissons' ratio corresponds to the nearly-incompressible case.)

- The mesh to be used is depicted in Figure 2. (Only half the domain need be modelled since the x-axis is a line of anti-symmetry.) You will need to compute nodal forces corresponding to the traction boundary conditions along $x = 0$ and $x = l$.
- Assume plane strain conditions.
- Use 4-node quadrilaterals and employ each of the following quadrature rules:
 - (i) 2×2
 - (ii) 2×2 μ -term, 1×1 λ -term
 - (iii) 2×2 with incompatible modes

Also employ (iv) the hourglass stiffness option.

- Use 3-node triangles in the
 - (v) bisection pattern
 - (vi) cross-diagonal pattern

(There are twelve cases in all.)

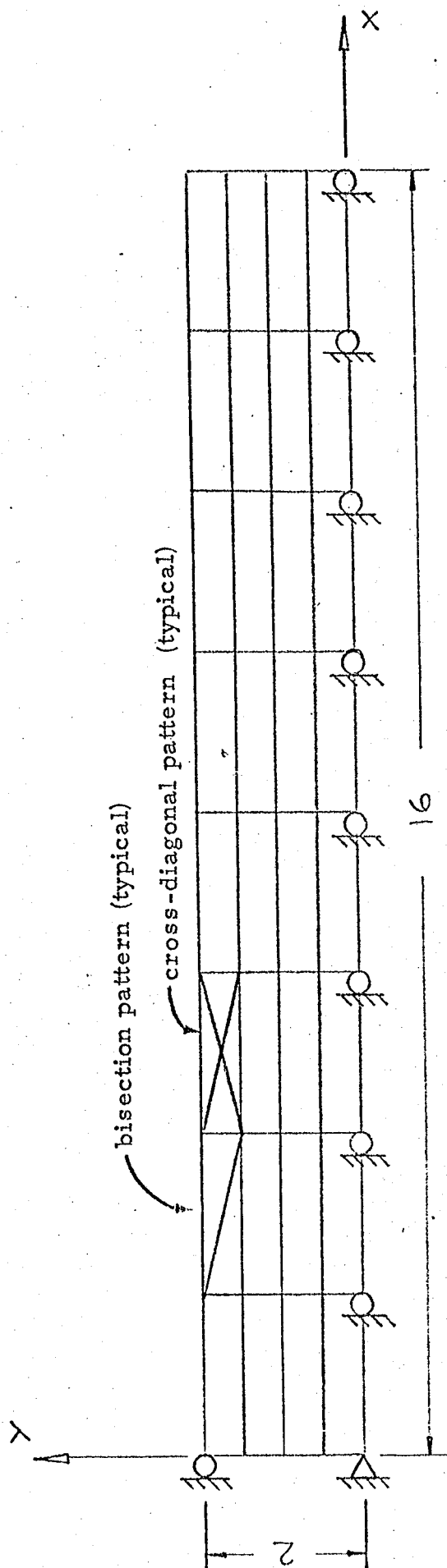


Figure 2

- In each case, compare the computed vertical tip displacement [i.e. $u_2(16,0)$], the horizontal displacement of the root section (i.e. $u_1(0,y)$, $y \in [-c,c]$) and the bending and shear stress (i.e. σ_x and τ_{xy} , respectively) for $x = 7$, $y \in [-c,c]$, with the exact solution.
- Write a brief report summarizing your findings and include an evaluation of the performance of each element.

Instructions for Using "LEARN"

Consult the LEARN User's Manual regarding deck set up.

To save keypunch labor, use the generation options on coordinate and element data.

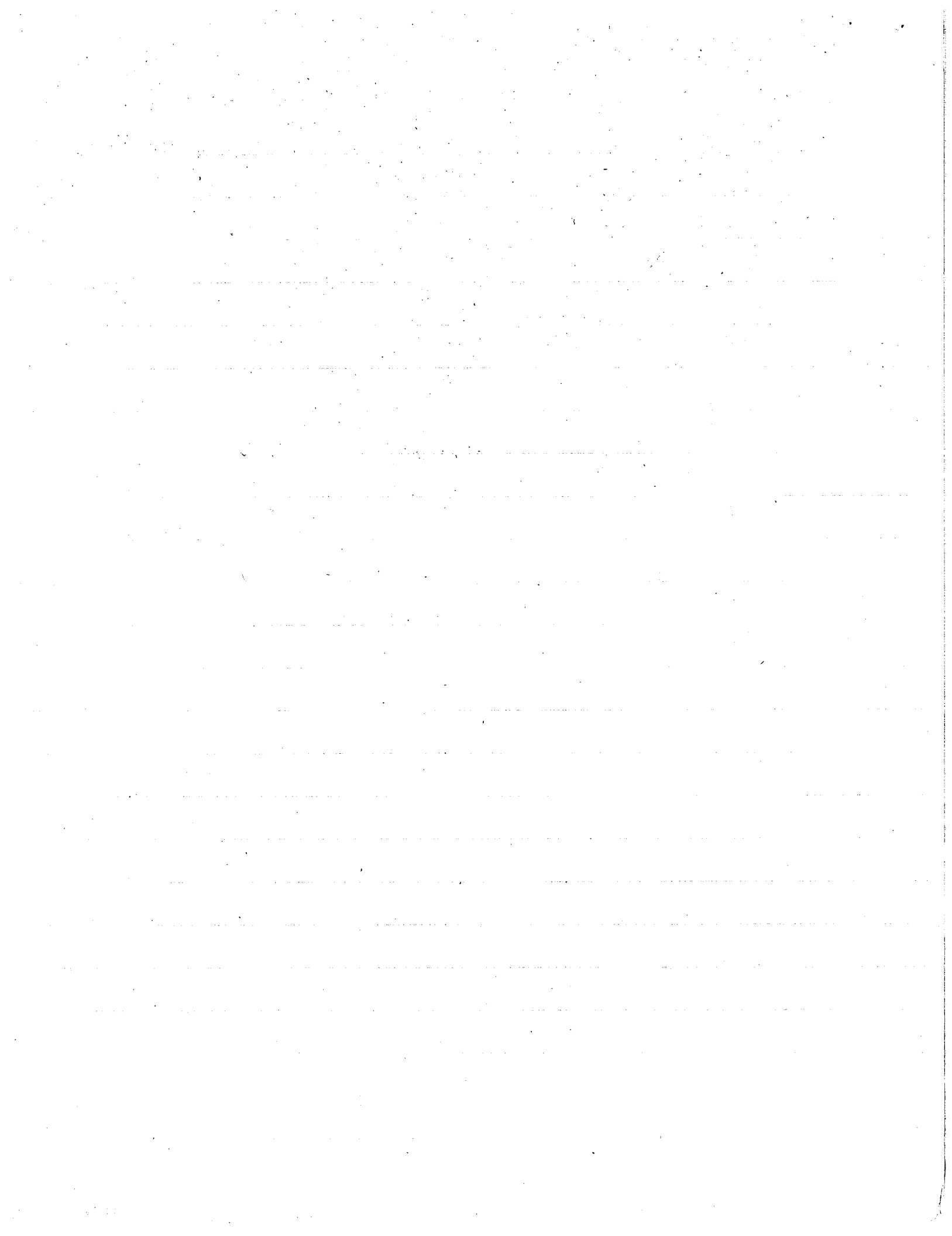
Your first run should be a data check run (i.e., MODE.EQ.0 on the master control card). Run the 2×2 quadrature case for $\nu = .3$. Carefully check the output to insure that you have keypunched your data correctly.

Keep using the data check mode until you are satisfied that your input data is correct. Then set MODE.EQ.1 which will cause the program to execute. The vertical displacement $u_2(16,0) = -220.4$. If you obtain this value, begin running the other cases. If you don't obtain the correct value, or if you are having any trouble whatsoever, see me or the T.A. immediately and bring your output and deck.

4 node quad: based on the results for all the compared items (ie $u_2(16,0)$, $u_1(0,y)$, $\sigma_x(7,y)$, $\tau_{xy}(7,y)$) it was found that the 2x2 w/incompressible and the 1x1 w/hourglass elements were the best and provided fairly accurate results for the displacements and stresses. The results degraded somewhat for the incompressible case. The next best case was the reduced selective integration case ($1 \times 1 \lambda$ $2 \times 2 \mu$) which was fairly accurate in the stresses but not in the displacements. The results were somewhat upgraded for the incompressible case. The 2x2 case shows degradation of results as we approach the incompressible case. Severe lock-up of the elements occurs; the extensional stresses degrade but the shear stresses remain fairly accurate. This was considered to be the worst case. In terms of overall time, for each value of ν , they are comparable.

3 node bisection: The first thing we note is that the results are the same for all the cases run (ie $\nu = .499$: 2×2 , 2×2 w/incomp, 1×1 w/hourglass & reduced selective int.) since $\nabla \cdot u = \text{constant}$.
 As the value of $\nu \rightarrow .5$ The elements suffer severe lock-up, but overall the displacements agree better; σ_x becomes worse and τ_{xy} becomes better. The time comparisons are somewhat worse than the 4-node quad, with respect to the 4 node quad as $\nu \rightarrow .5$ the displacement $u_1(0,y)$ are better and the stresses σ_x , τ_{xy} degrade.

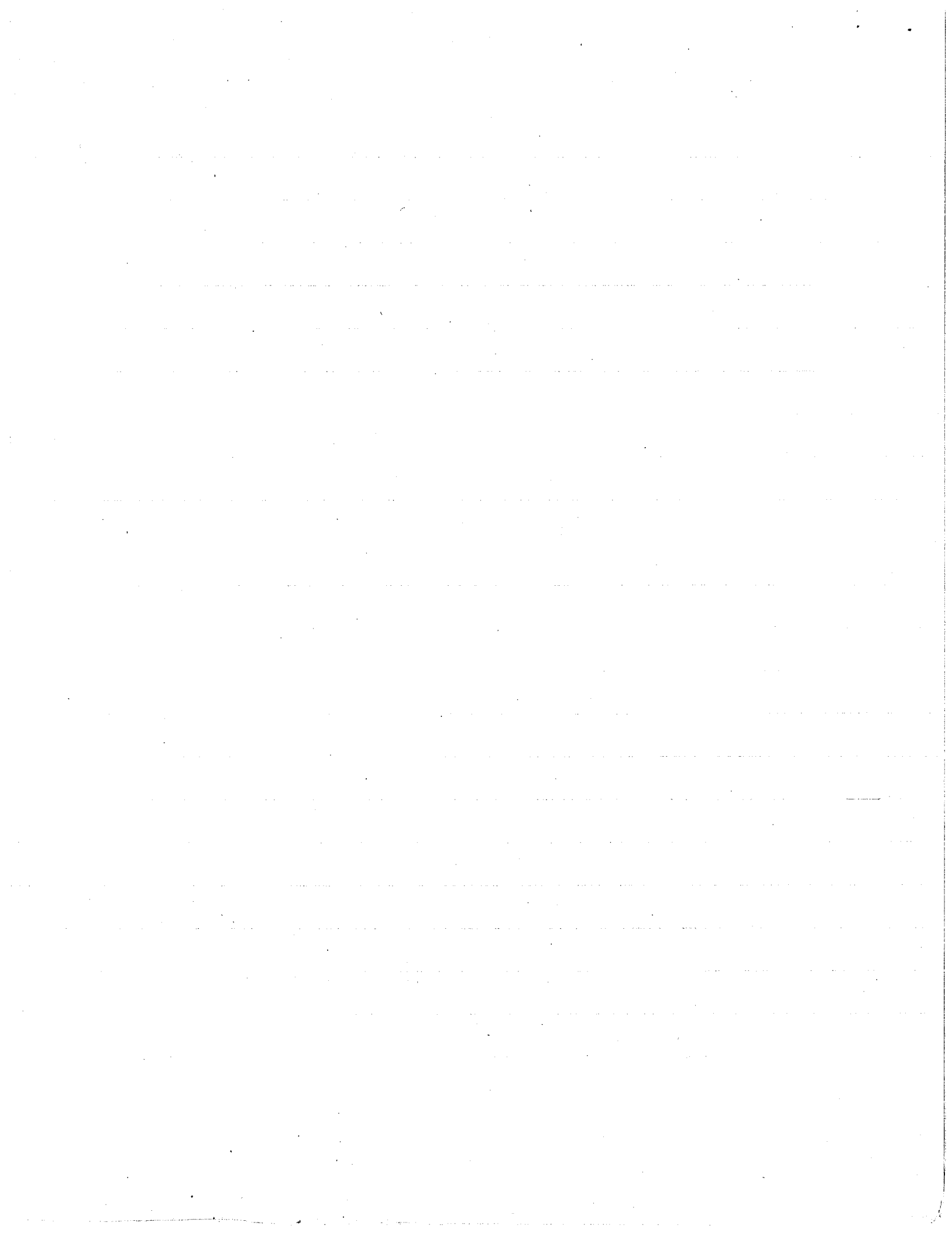
3 node cross-drag: as in the bisection case we note that the results are the same for all cases for the same value of ν .^{Since $\nabla \cdot u = \text{constant}$, that} We note that the errors in $u_1(0,y)$ decrease, that the errors in σ_x decrease, but the errors in τ_{xy} increase as $\nu \rightarrow .5$. We also note a



reversal in sign in τ_{xy} . By looking at a quadrilateral, we find that τ_{xy} exhibits a checkerboard pattern which should not happen. ^{For the 2x2 case,} with respect to the 4-node quad as $\nu \rightarrow .5$, $u_1(0,y)$ is better; $\sigma_x(7,y)$ is better but τ_{xy} is worse. With respect the 3 node bisection, as $\nu \rightarrow .5$, $u_1(0,y)$ is worse, σ_x is better and τ_{xy} becomes worse.

Overall comparisons:

The 4-node quad (2x2 w/incomp and 1x1 $\lambda \mu$ w/hourglass) by far provides the most superior results. In general for $\nu = .3$ the 4-node quad does much better than the other two elements. In general for $\nu = .499$ the 4 node quad 2x2 case gives better results in τ_{xy} ; the 2x2 cross doing gives better results in σ_x & $u_2(1)$ and the 2x2 bisection gives better results in $u_1(0,y)$.



4 node quad $u_2(16,0)$

y	EXACT	2×2	$2 \times 2 \mu$ $1 \times 1 \lambda$	2×2 W/INCOMP	$1 \times 1 \lambda$ W/HOURLASS
.3	-244.14	-220.78	-222.54	-242.98	-242.98
ERROR		(10)	(9)	(0.5)	(0.5)
.499	-205.47	-68.68	-192.65	-204.12	-204.12
ERROR		(66)	(6)	(0.7)	(0.7)

$u_1(0, y)$ $v = .3$ (ERROR IN PARENTH)

y	EXACT				
0	0	0.0	0.0	0.0 (0)	0.0 (0)
.5	-.1295	-.3455 (167)	-.3375 (161)	-.1372 (6)	-.1372 (6)
1.0	-.2072	-.5809 (180)	-.5677 (174)	-.2207 (7)	-.2207 (7)
1.5	-.1813	-.5606 (209)	-.5490 (203)	-.1953 (8)	-.1953 (8)
2.0	0	0.0	0.0	0.0 (0)	0.0 (0)

STIFF & LOAD FORM

.06

.06

.06

.06

$v = .499$

STIFF FACTOR/12

.00

.00

.00

.00

y	EXACT				
0	0	0	0	0 (0)	0 (0)
.5	-.1317	-.5089 (286)	-.2727 (107)	-.1466 (11)	-.1466 (11)
1.0	-.2107	-.8248 (291)	-.4595 (118)	-.2369 (12)	-.2369 (12)
1.5	-.1843	-.7391 (301)	-.4571 (148)	-.2117 (10)	-.2117 (10)
2.0	0	0	0	0 (0)	0 (0)

STIFF & LOAD FORM

.04

.06

.04

.04

STIFF FACTOR/12

.04

.00

.04

.04

4 node quad $\sigma_x(7, y)$ $v = .3$ (error in parent.)

<u>y</u>	<u>EXACT</u>	<u>2x2</u>	<u>2x2 μ</u>	<u>1x1 λ</u>	<u>2x2 w/inecomp</u>	<u>1x1 λ</u>	<u>λ μ w/HOUR</u>
.25	.4219	.3871 (8)	.3903 (7)	.4218 (20)	.4218 (20)	.4218 (20)	.4218 (20)
.75	1.2656	1.161 (8)	1.171 (7)	1.266 (20)	1.266 (20)	1.266 (20)	1.266 (20)
1.25	2.1094	1.934 (8)	1.949 (7)	2.110 (20)	2.110 (20)	2.110 (20)	2.110 (20)
1.75	2.9531	2.706 (8)	2.726 (8)	2.953 (20)	2.953 (20)	2.953 (20)	2.953 (20)

$v = .499$ (error in parent)

<u>y</u>	<u>EXACT</u>						
.25	.4219	.1252 (70)	.4005 (5)	.4227 (20)	.4227 (20)	.4227 (20)	.4227 (20)
.75	1.2656	.3934 (69)	1.206 (5)	1.264 (20)	1.264 (20)	1.264 (20)	1.264 (20)
1.25	2.1094	.7029 (66)	1.946 (5)	2.112 (20)	2.112 (20)	2.112 (20)	2.112 (20)
1.75	2.9531	1.081 (63)	2.800 (5)	2.952 (20)	2.952 (20)	2.952 (20)	2.952 (20)

$\tau_{xy}(7, y)$ $v = .3$ (error in parent)

<u>y</u>	<u>EXACT</u>						
.25	-.3691	-.3580 (3)	-.3596 (3)	-.3673 (~.5)	-.3673 (~.5)	-.3673 (~.5)	-.3673 (~.5)
.75	-.3223	-.3150 (2)	-.3160 (2)	-.3204 (~.6)	-.3204 (~.6)	-.3204 (~.6)	-.3204 (~.6)
1.25	-.2285	-.2280 (~0)	-.2285 (~0)	-.2266 (~1)	-.2266 (~1)	-.2266 (~1)	-.2266 (~1)
1.75	-.0879	-.0991 (13)	-.09695 (10)	-.0858 (~2)	-.0858 (~2)	-.0858 (~2)	-.0858 (~2)

$v = .499$ (error in parent)

.25	-.3691	-.3523 (4)	-.3629 (2)	-.3671 (~.5)	-.3671 (~.5)	-.3671 (~.5)	-.3671 (~.5)
.75	-.3223	-.3103 (4)	-.3184 (1)	-.3205 (~.6)	-.3205 (~.6)	-.3205 (~.6)	-.3205 (~.6)
1.25	-.2285	-.2284 (~0)	-.2266 (~1)	-.2268 (~1)	-.2268 (~1)	-.2268 (~1)	-.2268 (~1)
1.75	-.0879	-.1089 (24)	-.0921 (5)	-.0856 (~3)	-.0856 (~3)	-.0856 (~3)	-.0856 (~3)

3-node triangle $u_2(16,0)$

BISECTION

v	EXACT	2x2	2x2 w/INCOMP	1x1 λ/μ w/hour	1x1 λ 2x2 μ
.3	-244.14	-183.15 (25)			
.499	-205.47	-39.03 (81)	-39.03	-39.03	-39.03

$u_1(0,y)$ $v=.499$ (error in parenthesis)

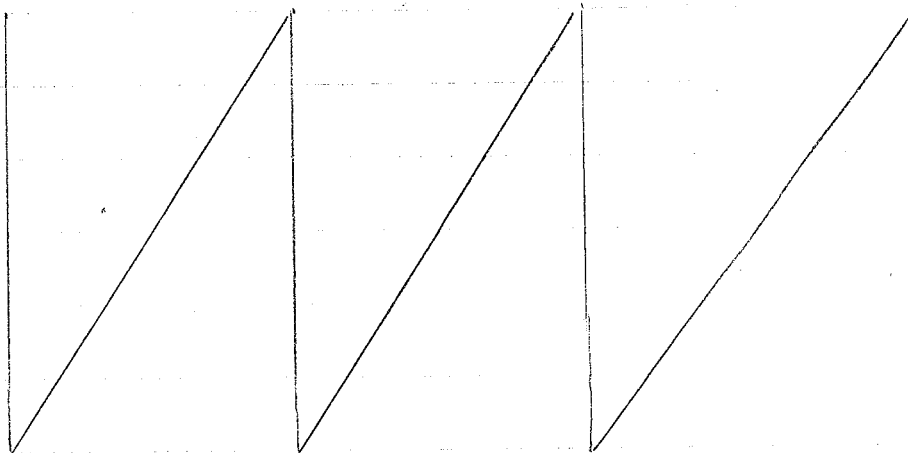
y	EXACT				
0	0	0	0	0	0
.5	-.1317	-.1630 (23)	-.1630	-.1630	-.1630
1.0	-.2107	-.4392 (108)	-.4392	-.4392	-.4392
1.5	-.1843	-.5128 (178)	-.5128	-.5128	-.5128
2.0	0	0	0	0	0
STIFF & LOAD FORM		.04	.04	.04	.04
STIFF FACTOR 12		.04	.04	.04	.04

$v=.3$ (error in parenthesis)

y	EXACT	
0	0	0
.5	-.1295	-.5485 (324)
1.0	-.2072	-.9056 (337)
1.5	-.1813	-.8333 (360)
2.0	0	0.0

STIFF & LOAD FORM .04

STIFF FACTOR 12 .04



3-node triangle

BISECTION

$\sigma_x(7, y)$ $\nu = .3$ (error in parenthesis)

<u>y</u>	<u>EXACT</u>	<u>2x2</u>	<u>2x2</u> <u>w/ INCOMP</u>	<u>1x1 λ/μ</u> <u>w/ HOURGLASS</u>	<u>1x1 λ/μ</u> <u>2x2 μ</u>
.375	.6328	.7142 (13)			
.875	1.4765	1.378 (7)			
1.375	2.3203	2.041 (12)			
1.875	3.164	2.703 (15)			

$\nu = .499$

<u>y</u>	<u>EXACT</u>				
.375	.6328	11.99 (1795)	11.99	11.99	11.99
.875	1.4765	9.966 (575)	9.966	9.966	9.966
1.375	2.3203	8.191 (253)	8.191	8.191	8.191
1.875	3.164	5.411 (71)	5.411	5.411	5.411

$\tau_{xy}(7, y)$ $\nu = .3$ (error in parenthesis)

<u>y</u>	<u>EXACT</u>				
.375	-.3618	.1143 (132)			
.875	-.3032	.1529 (150)			
1.375	-.1978	.2275 (215)			
1.875	-.0458	.3381 (838)			

$\nu = .499$ (error in parenthesis)

<u>y</u>	<u>EXACT</u>				
.375	-.3618	-.2741 (24)	-.2741	-.2741	-.2741
.875	-.3032	-.2195 (28)	-.2195	-.2195	-.2195
1.375	-.1978	-.1693 (14)	-.1693	-.1693	-.1693
1.875	-.0458	-.1140 (149)	-.1140	-.1140	-.1140

3-node triangle

$u_2(16,0)$

CROSS-DIAG.

y	EXACT	2x2	2x2 W/INCOMP	1x1 λ 2x2 μ	1x1 λ μ W/HOUR
.3	-224.14	-212.93 (13)			-212.93
.499	-205.47	-183.17 (11)	-183.17	-183.17	-183.17

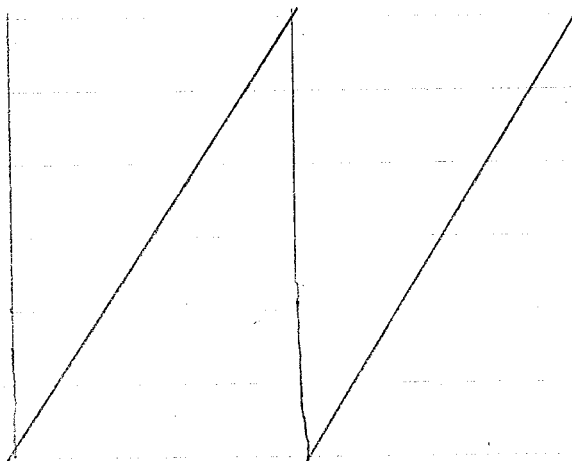
$u_1(0,y)$ $v=.499$

(error in parenthesis)

y	EXACT				
0	0	0	0	0	0
.5	-.1317	-.3112 (136)	-.3112	-.3112	-.3112
1.0	-.2107	-.5195 (147)	-.5195	-.5195	-.5195
1.5	-.1843	-.5139 (178)	-.5139	-.5139	-.5139
2.0	0	0.	0	0	0
STIFF & LOAD FORM.		.04	.04	.04	.04
STIFF FACTOR 12		.04	.04	.04	.04

$v=.3$ (error in parenthesis)

y	EXACT			
0	0	0.		0.
0.5	-.1295	-.3960 (206)		-.3960
1.0	-.2072	-.6625 (220)		-.6625
1.5	-.1813	-.6304 (248)		-.6304
2.0	0	0.		0
STIFF & LOAD FORM		0.0		0.06
STIFF FACTOR 12		.06		.06



3-node triangle

CROSS-DIAG

$\sigma_x(7,y)$ $\nu = .3$ (error in parenthesis)

y	EXACT	2x2	2x2 w/ INCOMP	1x1 λ	2x2 μ	1x1 λ/μ w/ HOUR
.125	.2109	.0537 (75)				.0537
.625	1.0547	.8006 (24)				.8006
1.125	1.8984	1.547 (19)				1.547
1.625	2.7422	2.294 (16)				2.294

$\nu = .499$ (error in parenthesis)

y	EXACT				
.125	.2109	.0798 (62)	.0798	.0798	.0798
.625	1.0547	.8737 (17)	.8737	.8737	.8737
1.125	1.8984	1.630 (14)	1.630	1.630	1.630
1.625	2.7422	2.415 (12)	2.415	2.415	2.415

$\tau_{xy}(7,y)$ $\nu = .3$

y	EXACT			
.125	-.3735	-.3521 (6)		-.3521
.625	-.3384	-.3077 (9)		-.3077
1.125	-.2564	-.2211 (14)		-.2211
1.625	-.1274	-.0931 (27)		-.0931

$\nu = .499$ (error in parenthesis)

y	EXACT			
.125	-.3735	-.3162 (15)	-.3162	-.3162
.625	-.3384	-.2021 (40)	-.2021	-.2021
1.125	-.2564	-.0354 (86)	-.0354	-.0354
1.625	-.1274	.1722 (235)	.1722	.1722

ME 235B
Finite Element Methods
Spring Quarter 1981

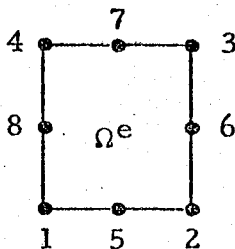
Stanford University

Instructor: T. Hughes

Homework Assignment
(due 1 week after handout)

Instructions: To cut down on computer costs, attempt to run all the data sets in problems 2 and 3 in one run.

1. Consider the 8-node isoparametric element shown below



Assume the nodal coordinates are as follows:

a	x_a	y_a	
1	$-\frac{h}{2}$	$-\frac{h}{2}$	$\frac{h}{2} (-1, -1)$
2	$\frac{h}{2}$	$-\frac{h}{2}$	$\frac{h}{2} (1, -1)$
3	$\frac{h}{2}$	$\frac{h}{2}$	$\frac{h}{2} (1, 1)$
4	$-\frac{h}{2}$	$\frac{h}{2}$	$\frac{h}{2} (-1, 1)$
5	0	$-\frac{h}{2}$	$\frac{h}{2} (0, -1)$
6	$\frac{h}{2}$	0	$\frac{h}{2} (1, 0)$
7	0	$\frac{h}{2}$	$\frac{h}{2} (0, 1)$
8	$-\frac{h}{2}$	0	$\frac{h}{2} (-1, 0)$

Consider the following body force

$$\tilde{\mathbf{f}} = \begin{Bmatrix} 0 \\ -g \end{Bmatrix} ; \quad g \text{ const.}$$

10000

10000

10000

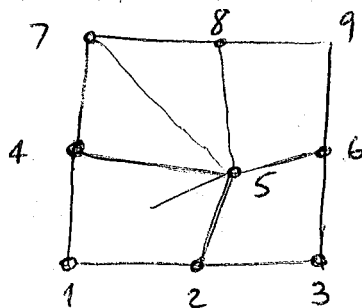
10000

10000

10000

10000

10000



1.	PATCH	TEST #1	D1A=1	D2A=0					
2	9	2	2	1	0	1	← execution; 0 for data check.		
3	1	1		0.0		0.0			
4	3	0		2.0		0.0			
5	4	0		0.0		1.0			
6	5	8		0.1		0.8			
7	6	0		2.0		1.0			
	7	1		0.0		2.0			
	9	0		2.0		2.0			
0	blank								
1	1	1	1	1				change these	
2	4	0	1	1					
3	6	1	1	1					
4	9	0	1	1					
5	blank								
6	1	1		1.0		0.0		↑ Change these	
7	4	0		1.0		0.0			
8	6	1		1.0		0.0			
9	9	0		1.0		0.0		↓	
10	blank								
11	3	4	1	0	1	0	0	0	0
12	1		1	1.0	0.3	0.3	0.3	0.0	1.
13		0.0		0.0					
14	1	1	1	2	5	4	0		
15	2	1	2	3	6	5	0		
16	3	1	4	5	8	7	0		
17	4	1	5	6	9	8	0		
18	blank								
19	PATCH TEST #2	D1A=0	D2A=1						
20-43	Copy Cards 2-15 here								
4	1	1		0.0	56	30	1.0	598	or copy cards 2-28 FROM
5	4	0		0.0	116	30	1.0	626	FILE1 to end
6	6	1		0.0	146	30	1.0	656	& change in active cards
7	9	0		0.0	176	30	1.0	686	44-47
8-56	copy cards 20-28				190	14	1.0	716	
						14		746	
						14		760	
						14		774	
						14		788	

0.6709941E+01	0.9813007E+00	0.344501E-03	0.1157379E-01	0.2901493E+04	1258
0.6714940E+01	0.9813573E+00	0.3430704E-03	0.1152086E-01	0.2914853E+04	1259
0.6719939E+01	0.9814137E+00	0.3414967E-03	0.1146936E-01	0.2928285E+04	1260
0.6724938E+01	0.9814693E+00	0.3399483E-03	0.1141739E-01	0.2941622E+04	1261
0.6729938E+01	0.9815241E+00	0.3384261E-03	0.1136684E-01	0.2954854E+04	1262
0.6734937E+01	0.9815784E+00	0.3369192E-03	0.1131630E-01	0.2968070E+04	1263
0.6739936E+01	0.9816331E+00	0.3354033E-03	0.1126623E-01	0.2981465E+04	1264
0.6744935E+01	0.9816880E+00	0.3338847E-03	0.1121569E-01	0.2995045E+04	1265
0.6749934E+01	0.9817426E+00	0.3323767E-03	0.1116514E-01	0.3008634E+04	1266
0.6754933E+01	0.9817963E+00	0.3308975E-03	0.1111555E-01	0.3022083E+04	1267
0.6759933E+01	0.9818500E+00	0.3294174E-03	0.1106644E-01	0.3035661E+04	1268
0.6764932E+01	0.9819034E+00	0.3279494E-03	0.1101732E-01	0.3049249E+04	1269
0.6769931E+01	0.9819567E+00	0.3264886E-03	0.1096869E-01	0.3062893E+04	1270
0.6774930E+01	0.9820101E+00	0.3250251E-03	0.1092005E-01	0.3076685E+04	1271
0.6779929E+01	0.9820637E+00	0.3235594E-03	0.1087141E-01	0.3090622E+04	1272
0.6784928E+01	0.9821170E+00	0.3221033E-03	0.1082277E-01	0.3104603E+04	1273
0.6789927E+01	0.9821693E+00	0.3206749E-03	0.1077509E-01	0.3118422E+04	1274
0.6794927E+01	0.9822209E+00	0.3192709E-03	0.1072836E-01	0.3132135E+04	1275
0.6799926E+01	0.9822721E+00	0.3178790E-03	0.1068163E-01	0.3145850E+04	1276
0.6804925E+01	0.9823236E+00	0.3164795E-03	0.1063538E-01	0.3159760E+04	1277
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0.6814923E+01	0.9824267E+00	0.3136885E-03	0.1054192E-01	0.3187874E+04	1279
0.6819922E+01	0.9824782E+00	0.3122971E-03	0.1049519E-01	0.3202078E+04	1280
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0.6829921E+01	0.9825793E+00	0.3095719E-03	0.1040411E-01	0.3230266E+04	1282
0.6834920E+01	0.9826295E+00	0.3082244E-03	0.1035929E-01	0.3244410E+04	1283
0.6839919E+01	0.9826795E+00	0.3068780E-03	0.1031494E-01	0.3258623E+04	1284
0.6844918E+01	0.9827298E+00	0.3055308E-03	0.1027012E-01	0.3272991E+04	1285
0.6849917E+01	0.9827802E+00	0.3041625E-03	0.1022530E-01	0.3287499E+04	1286
0.6854917E+01	0.9828303E+00	0.3028410E-03	0.1018047E-01	0.3302062E+04	1287
0.6859916E+01	0.9828799E+00	0.3015287E-03	0.1013708E-01	0.3316432E+04	1288
0.6864915E+01	0.9829297E+00	0.3002407E-03	0.1009369E-01	0.3330660E+04	1289
0.6869914E+01	0.9829795E+00	0.2989622E-03	0.1005125E-01	0.3344902E+04	1290
0.6874913E+01	0.9830242E+00	0.2976782E-03	0.1000834E-01	0.3359332E+04	1291
0.6879912E+01	0.9830728E+00	0.2963908E-03	0.9965420E-02	0.3373922E+04	1292
0.6884912E+01	0.9831193E+00	0.2951121E-03	0.9922028E-02	0.3388542E+04	1293
0.6889911E+01	0.9831693E+00	0.2938369E-03	0.9879589E-02	0.3403248E+04	1294
0.6894910E+01	0.9832166E+00	0.2925855E-03	0.9837627E-02	0.3417805E+04	1295
0.6899909E+01	0.9832640E+00	0.2913359E-03	0.9796143E-02	0.3432463E+04	1296
0.6904908E+01	0.9833110E+00	0.2900993E-03	0.9754658E-02	0.3447094E+04	1297
0.6909907E+01	0.9833579E+00	0.2888662E-03	0.9713173E-02	0.3461809E+04	1298
0.6914907E+01	0.9834050E+00	0.2876289E-03	0.9672642E-02	0.3476688E+04	1299
0.6919906E+01	0.9834521E+00	0.2863940E-03	0.9631157E-02	0.3491691E+04	1300
0.6924905E+01	0.9834992E+00	0.2851633E-03	0.9589672E-02	0.3506762E+04	1301
0.6929904E+01	0.9835462E+00	0.2839344E-03	0.9548664E-02	0.3521937E+04	1302
0.6934903E+01	0.9835932E+00	0.2827381E-03	0.9508610E-02	0.3536916E+04	1303
0.6939902E+01	0.9836375E+00	0.2815549E-03	0.9469032E-02	0.3551703E+04	1304
0.6944901E+01	0.9836822E+00	0.2803993E-03	0.9430408E-02	0.3566457E+04	1305
0.6949901E+01	0.9837274E+00	0.2792166E-03	0.9390631E-02	0.3581448E+04	1306
0.6954900E+01	0.9837728E+00	0.2780380E-03	0.9351254E-02	0.3596628E+04	1307
0.6959899E+01	0.9838179E+00	0.2768710E-03	0.9312153E-02	0.3611788E+04	1308
0.6964898E+01	0.9838629E+00	0.2757069E-03	0.9273529E-02	0.3627038E+04	1309
0.6969897E+01	0.9839081E+00	0.2745402E-03	0.9234428E-02	0.3642454E+04	1310
0.6974896E+01	0.9839532E+00	0.2733995E-03	0.9196281E-02	0.3657651E+04	1311
0.6979895E+01	0.9839985E+00	0.2722619E-03	0.9158134E-02	0.3672932E+04	1312
0.6984894E+01	0.9840433E+00	0.2711355E-03	0.9120464E-02	0.3688192E+04	1313
0.6989893E+01	0.9840884E+00	0.2700135E-03	0.9083271E-02	0.3703516E+04	1314
0.6994892E+01	0.9841279E+00	0.2688873E-03	0.9045124E-02	0.3719030E+04	1315
0.6999892E+01	0.9841719E+00	0.2677594E-03	0.9007454E-02	0.3734694E+04	1316
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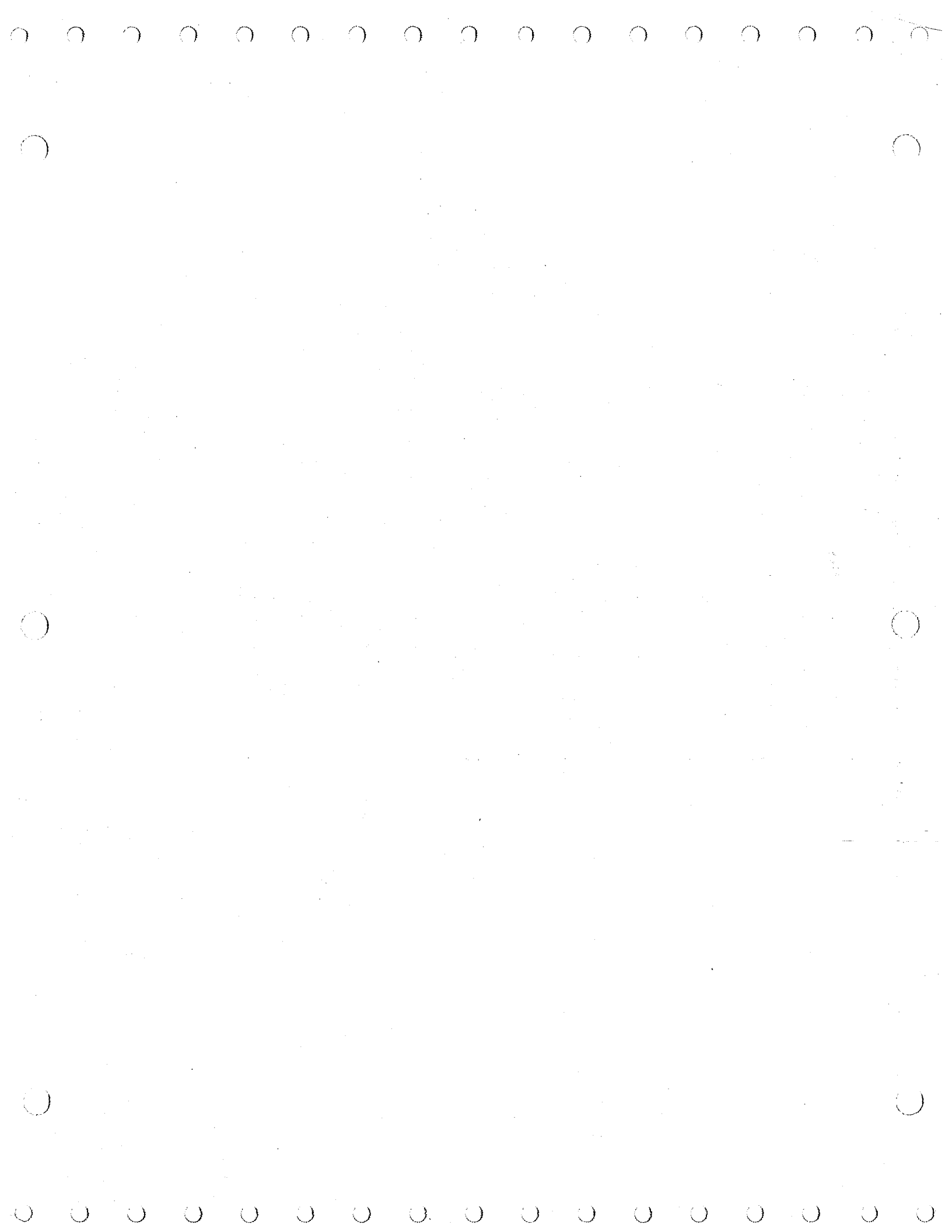
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0.6419990E+01	0.9775403E+00	0.4553176E-03	0.1524019E-01	0.2196268E+04	1200
0.6424989E+01	0.9776141E+00	0.4530386E-03	0.1516438E-01	0.2207317E+04	1201
0.6429988E+01	0.9776884E+00	0.4507517E-03	0.1508999E-01	0.2218515E+04	1202
0.6434987E+01	0.9777621E+00	0.4484854E-03	0.1501465E-01	0.2229727E+04	1203
0.6439986E+01	0.9778366E+00	0.4462593E-03	0.1494122E-01	0.2240850E+04	1204
0.6444985E+01	0.9779065E+00	0.4440576E-03	0.1486874E-01	0.2251959E+04	1205
0.6449984E+01	0.9779782E+00	0.4418669E-03	0.1479626E-01	0.2263125E+04	1206
0.6454983E+01	0.9780501E+00	0.4396727E-03	0.1472378E-01	0.2274419E+04	1207
0.6459982E+01	0.9781220E+00	0.4374811E-03	0.1465082E-01	0.2285813E+04	1208
0.6464981E+01	0.9781929E+00	0.4353262E-03	0.1458025E-01	0.2297127E+04	1209
0.6469980E+01	0.9782637E+00	0.4331396E-03	0.1450968E-01	0.2308536E+04	1210
0.6474979E+01	0.9783342E+00	0.4289087E-03	0.1436853E-01	0.2319972E+04	1211
0.6479978E+01	0.9784046E+00	0.4265032E-03	0.1429987E-01	0.2331498E+04	1212
0.6484977E+01	0.9784744E+00	0.4242704E-03	0.1422977E-01	0.2342999E+04	1213
0.6489976E+01	0.9785440E+00	0.4226228E-03	0.1416111E-01	0.2354576E+04	1214
0.6494975E+01	0.9786132E+00	0.4205445E-03	0.1409292E-01	0.2366178E+04	1215
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0.6504973E+01	0.9787517E+00	0.4164602E-03	0.1395607E-01	0.2389678E+04	1217
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0.6514971E+01	0.9788886E+00	0.4123796E-03	0.138208E-01	0.2413100E+04	1219
0.6519970E+01	0.9789554E+00	0.4103710E-03	0.1375532E-01	0.2424949E+04	1220
0.6524969E+01	0.9790229E+00	0.4083766E-03	0.1368952E-01	0.2436818E+04	1221
0.6529968E+01	0.9790900E+00	0.4064050E-03	0.1362419E-01	0.2448719E+04	1222
0.6534967E+01	0.9791565E+00	0.4044331E-03	0.1355886E-01	0.2460599E+04	1223
0.6539966E+01	0.9792231E+00	0.4024825E-03	0.1349449E-01	0.2472596E+04	1224
0.6544965E+01	0.9792891E+00	0.4005372E-03	0.1343012E-01	0.2484580E+04	1225
0.6549964E+01	0.9793550E+00	0.3985846E-03	0.1335527E-01	0.2496647E+04	1226
0.6554963E+01	0.9794213E+00	0.3966666E-03	0.1330185E-01	0.250877E+04	1227
0.6559962E+01	0.9794865E+00	0.3947557E-03	0.132843E-01	0.2521009E+04	1228
0.6564961E+01	0.9795516E+00	0.3928586E-03	0.1317549E-01	0.2533212E+04	1229
0.6569960E+01	0.9796163E+00	0.3909629E-03	0.1311207E-01	0.2545445E+04	1230
0.6574959E+01	0.9796811E+00	0.3890651E-03	0.1304913E-01	0.2557786E+04	1231
0.6579958E+01	0.9797460E+00	0.3872043E-03	0.1298761E-01	0.2570263E+04	1232
0.6584957E+01	0.9798099E+00	0.3853738E-03	0.1292658E-01	0.2582615E+04	1233
0.6589956E+01	0.9798728E+00	0.3835538E-03	0.128650E-01	0.2594882E+04	1234
0.6594955E+01	0.9799354E+00	0.3817240E-03	0.1280594E-01	0.2607196E+04	1235
0.6599954E+01	0.9800061E+00	0.3798965E-03	0.1274538E-01	0.2619693E+04	1236
0.6604953E+01	0.9800616E+00	0.3781074E-03	0.1268578E-01	0.2632296E+04	1237
0.6609952E+01	0.9801235E+00	0.3763479E-03	0.1262712E-01	0.2644750E+04	1238
0.6614951E+01	0.9801844E+00	0.3745994E-03	0.1256943E-01	0.2657115E+04	1239
0.6619950E+01	0.9802451E+00	0.3728401E-03	0.1251030E-01	0.2669518E+04	1240
0.6624949E+01	0.9803063E+00	0.3710790E-03	0.1245213E-01	0.2682114E+04	1241
0.6629948E+01	0.9803676E+00	0.3693614E-03	0.123938E-01	0.2694844E+04	1242
0.6634947E+01	0.9804275E+00	0.3676452E-03	0.1233864E-01	0.2707375E+04	1243
0.6639946E+01	0.9804875E+00	0.3659355E-03	0.1228189E-01	0.2720013E+04	1244
0.6644945E+01	0.9805473E+00	0.3642347E-03	0.1222563E-01	0.2732721E+04	1245
0.6649944E+01	0.9806069E+00	0.3625322E-03	0.1216888E-01	0.2745481E+04	1246
0.6654943E+01	0.9806667E+00	0.3608309E-03	0.1211214E-01	0.2758375E+04	1247
0.6659942E+01	0.9807265E+00	0.3591627E-03	0.1205683E-01	0.2771381E+04	1248
0.6664941E+01	0.9807853E+00	0.357538E-03	0.1200247E-01	0.2784253E+04	1249
0.6669940E+01	0.9808431E+00	0.3558365E-03	0.1194859E-01	0.2797015E+04	1250
0.6674939E+01	0.9809006E+00	0.3542581E-03	0.1189375E-01	0.2809804E+04	1251
0.6679938E+01	0.9809586E+00	0.3526211E-03	0.1183939E-01	0.2822799E+04	1252
0.6684937E+01	0.9810167E+00	0.350966E-03	0.1178503E-01	0.2835904E+04	1253
0.6689936E+01	0.9810744E+00	0.3493966E-03	0.1173162E-01	0.2849030E+04	1254
0.6694935E+01	0.9811312E+00	0.3478257E-03	0.1167917E-01	0.2862051E+04	1255
0.6699934E+01	0.9811873E+00	0.3462411E-03	0.1162720E-01	0.2875002E+04	1256
0.6704933E+01	0.9812438E+00			0.2888164E+04	1257

	5	10	15	20	25	30	35	40	45
186	5	1	1	0	1	0	0	0	0
187	1			1.0		0.3		0.0	
188		0.0		0.0					
189	1	1	1	2	3	4	0		
190	blank -								
191	Rank Check Run for 1x1 A and 1x1 M w/stiffness								
192-199	copy 178 to 187 here								
200	3	1	1	0	1	1	0	0	0 X
201-204	copy 189-192 here -								
205	Rank Check Run for 1x1 L and 1x1 M w/stiffness								
206-213	copy 178 to 187 here								
214	3	1	1	0	1	1	0	1	0 X
215-218	copy 189-192 here								
219	blank								

or copy 178-192 from file 1
 & change 216 of active file

0.6110042E+01	0.9720769E+00	0.6352291E-03	0.2111387E-01	0.1574235E+04	1138	1
0.6115041E+01	0.9721808E+00	0.6315390E-03	0.2099705E-01	0.1532835E+04	1139	1
0.6120040E+01	0.9722841E+00	0.6229999E-03	0.2088070E-01	0.1592357E+04	1140	1
0.6125039E+01	0.9723867E+00	0.6244357E-03	0.2076530E-01	0.1601446E+04	1141	1
0.6130038E+01	0.9724877E+00	0.6209309E-03	0.2065182E-01	0.1610485E+04	1142	1
0.6135037E+01	0.9725881E+00	0.6174541E-03	0.2053928E-01	0.1619553E+04	1143	1
0.6140037E+01	0.9726884E+00	0.6139888E-03	0.2042723E-01	0.1628694E+04	1144	1
0.6145036E+01	0.9727880E+00	0.6105558E-03	0.2031612E-01	0.1628694E+04	1145	1
0.6150035E+01	0.9728876E+00	0.6071297E-03	0.2020502E-01	0.1637852E+04	1146	1
0.6155034E+01	0.9729866E+00	0.6037317E-03	0.2009487E-01	0.1647094E+04	1147	1
0.6160033E+01	0.9730846E+00	0.5970427E-03	0.1998615E-01	0.1656623E+04	1148	1
0.6165032E+01	0.9731821E+00	0.5937303E-03	0.1987791E-01	0.1674921E+04	1149	1
0.6170032E+01	0.9732801E+00	0.5903345E-03	0.197667E-01	0.1684343E+04	1150	1
0.6175031E+01	0.9733773E+00	0.5871027E-03	0.1966238E-01	0.1703279E+04	1151	1
0.6180030E+01	0.9734742E+00	0.5838363E-03	0.195509E-01	0.1722346E+04	1152	1
0.6185029E+01	0.9735706E+00	0.5806035E-03	0.1944923E-01	0.1712809E+04	1153	1
0.6190028E+01	0.9736682E+00	0.5773988E-03	0.1934385E-01	0.1731905E+04	1154	1
0.6195027E+01	0.9737611E+00	0.5742230E-03	0.192390E-01	0.1741483E+04	1155	1
0.6200027E+01	0.9738554E+00	0.5710719E-03	0.1913643E-01	0.1751093E+04	1156	1
0.6205026E+01	0.9739492E+00	0.5679499E-03	0.1903391E-01	0.1760720E+04	1157	1
0.6210025E+01	0.9740423E+00	0.5648250E-03	0.1893234E-01	0.1770460E+04	1158	1
0.6215024E+01	0.9741356E+00	0.5617575E-03	0.1883030E-01	0.1780127E+04	1159	1
0.6220023E+01	0.9742275E+00	0.5587034E-03	0.1873064E-01	0.1789858E+04	1160	1
0.6225022E+01	0.9743191E+00	0.5556599E-03	0.1863098E-01	0.1799662E+04	1161	1
0.6230021E+01	0.9744106E+00	0.5526519E-03	0.1853180E-01	0.1809457E+04	1162	1
0.6235020E+01	0.9745012E+00	0.5496799E-03	0.1843367E-01	0.1819240E+04	1163	1
0.6240020E+01	0.9745909E+00	0.5467280E-03	0.1824045E-01	0.1829021E+04	1164	1
0.6245019E+01	0.9746802E+00	0.5437676E-03	0.1814365E-01	0.1839021E+04	1165	1
0.6250018E+01	0.9747699E+00	0.5408411E-03	0.1804829E-01	0.1848971E+04	1166	1
0.6255017E+01	0.9748588E+00	0.5379280E-03	0.1795292E-01	0.1858985E+04	1167	1
0.6260016E+01	0.9749474E+00	0.5350364E-03	0.1785851E-01	0.1869031E+04	1168	1
0.6265016E+01	0.9750355E+00	0.5321763E-03	0.1776505E-01	0.1879076E+04	1169	1
0.6270015E+01	0.9751229E+00	0.5292470E-03	0.1767254E-01	0.1889119E+04	1170	1
0.6275014E+01	0.9752095E+00	0.5263404E-03	0.1758051E-01	0.1899189E+04	1171	1
0.6280013E+01	0.9752956E+00	0.5237262E-03	0.1748848E-01	0.1909394E+04	1172	1
0.6285012E+01	0.9753821E+00	0.5209409E-03	0.1739740E-01	0.1919603E+04	1173	1
0.6290011E+01	0.9754679E+00	0.5181725E-03	0.1730680E-01	0.1929859E+04	1174	1
0.6295011E+01	0.9755533E+00	0.5154209E-03	0.1721668E-01	0.1940161E+04	1175	1
0.6300010E+01	0.9756383E+00	0.5127059E-03	0.1712751E-01	0.1950436E+04	1176	1
0.6305009E+01	0.9757224E+00	0.5100169E-03	0.1703930E-01	0.1960719E+04	1177	1
0.6310008E+01	0.9758059E+00	0.5073508E-03	0.1695204E-01	0.1971022E+04	1178	1
0.6315007E+01	0.9758888E+00	0.5046758E-03	0.1686430E-01	0.1981470E+04	1179	1
0.6320006E+01	0.9759721E+00	0.5020292E-03	0.1677752E-01	0.1991916E+04	1180	1
0.6325006E+01	0.9760547E+00	0.4993966E-03	0.1669121E-01	0.2002406E+04	1181	1
0.6330005E+01	0.9761370E+00	0.4967914E-03	0.1660538E-01	0.2012917E+04	1182	1
0.6335004E+01	0.9762187E+00	0.4942114E-03	0.1652098E-01	0.2023425E+04	1183	1
0.6340003E+01	0.9762997E+00	0.4916571E-03	0.1643705E-01	0.2033938E+04	1184	1
0.6345002E+01	0.9763801E+00	0.4890987E-03	0.1635313E-01	0.2044577E+04	1185	1
0.6350001E+01	0.9764607E+00	0.4865618E-03	0.1626968E-01	0.2055237E+04	1186	1
0.6355000E+01	0.9765408E+00	0.4840686E-03	0.1618767E-01	0.2065822E+04	1187	1
0.6360000E+01	0.9766197E+00	0.4816016E-03	0.1610661E-01	0.2076440E+04	1188	1
0.6365000E+01	0.9766979E+00	0.4791424E-03	0.1602554E-01	0.2087062E+04	1189	1
0.6369988E+01	0.9767759E+00	0.4768847E-03	0.1594496E-01	0.2097832E+04	1190	1
0.6374977E+01	0.9768541E+00	0.4742593E-03	0.1586485E-01	0.2108551E+04	1191	1
0.6379966E+01	0.9769315E+00	0.4718404E-03	0.1578522E-01	0.2119360E+04	1192	1
0.6384955E+01	0.9770087E+00	0.4694390E-03	0.1570654E-01	0.2130202E+04	1193	1
0.6389955E+01	0.9770855E+00	0.4670441E-03	0.1562738E-01	0.2141125E+04	1194	1
0.6394944E+01	0.9771623E+00	0.4646734E-03	0.1554871E-01	0.2152049E+04	1195	1
0.6399933E+01	0.9772384E+00	0.4623113E-03	0.1547050E-01	0.2163044E+04	1196	1
0.6404922E+01	0.9773144E+00				1197	1



0.5810092E+01	0.9646033E+00	0.9170666E-03	0.3082739E-01	0.1090433E+04	1078	1
0.5815091E+01	0.9647506E+00	0.9111117E-03	0.2964333E-01	0.1097506E+04	1079	1
0.5830090E+01	0.9648976E+00	0.9051857E-03	0.2965975E-01	0.1104746E+04	1080	1
0.5835089E+01	0.9650435E+00	0.8993191E-03	0.2947807E-01	0.1111952E+04	1081	1
0.5830089E+01	0.9651893E+00	0.8934708E-03	0.2929687E-01	0.1119230E+04	1082	1
0.5835088E+01	0.9653335E+00	0.8876689E-03	0.2917711E-01	0.1126520E+04	1083	1
0.5840087E+01	0.9654775E+00	0.8819643E-03	0.2893925E-01	0.1133833E+04	1084	1
0.5845086E+01	0.9656199E+00	0.8762968E-03	0.2876282E-01	0.1141166E+04	1085	1
0.5850085E+01	0.9657609E+00	0.8651423E-03	0.2858877E-01	0.1148949E+04	1086	1
0.5855084E+01	0.9660410E+00	0.8596387E-03	0.284354E-01	0.1163279E+04	1088	1
0.5860083E+01	0.9661805E+00	0.8541460E-03	0.2807188E-01	0.1170706E+04	1089	1
0.5870082E+01	0.9663188E+00	0.8487189E-03	0.2790260E-01	0.1178246E+04	1090	1
0.5875081E+01	0.9664563E+00	0.8433093E-03	0.2773333E-01	0.1185804E+04	1091	1
0.5880080E+01	0.9665939E+00	0.8379607E-03	0.2756596E-01	0.1193373E+04	1092	1
0.5885079E+01	0.9667299E+00	0.8326659E-03	0.2740002E-01	0.1200962E+04	1093	1
0.5890078E+01	0.9668646E+00	0.8274305E-03	0.2723551E-01	0.1208561E+04	1094	1
0.5895077E+01	0.9669989E+00	0.8222293E-03	0.2707243E-01	0.1216206E+04	1095	1
0.5900076E+01	0.9671320E+00	0.8170852E-03	0.2691078E-01	0.1223863E+04	1096	1
0.5905075E+01	0.9672652E+00	0.8119494E-03	0.2674913E-01	0.1231604E+04	1097	1
0.5910074E+01	0.9673971E+00	0.8068776E-03	0.2658939E-01	0.1239345E+04	1098	1
0.5915073E+01	0.9675288E+00	0.8018252E-03	0.2643061E-01	0.1247155E+04	1099	1
0.5920072E+01	0.9676595E+00	0.7968273E-03	0.2627277E-01	0.1254977E+04	1100	1
0.5925071E+01	0.9677891E+00	0.7918798E-03	0.2611685E-01	0.1262818E+04	1101	1
0.5930070E+01	0.9679175E+00	0.7869932E-03	0.2596235E-01	0.1270659E+04	1102	1
0.5935069E+01	0.9680454E+00	0.7821361E-03	0.2580881E-01	0.1278505E+04	1103	1
0.5940068E+01	0.9681721E+00	0.7773379E-03	0.2565718E-01	0.1286441E+04	1104	1
0.5945067E+01	0.9682989E+00	0.7725458E-03	0.2550554E-01	0.1294421E+04	1105	1
0.5950066E+01	0.9684246E+00	0.7678100E-03	0.2535582E-01	0.1302405E+04	1106	1
0.5955065E+01	0.9685499E+00	0.7630971E-03	0.2520609E-01	0.1310449E+04	1107	1
0.5960064E+01	0.9686741E+00	0.7584379E-03	0.2505875E-01	0.1318499E+04	1108	1
0.5965063E+01	0.9687974E+00	0.7538239E-03	0.2491188E-01	0.1326570E+04	1109	1
0.5970062E+01	0.9689195E+00	0.7492674E-03	0.2476740E-01	0.1334637E+04	1110	1
0.5975061E+01	0.9690410E+00	0.7447463E-03	0.2462339E-01	0.1342746E+04	1111	1
0.5980060E+01	0.9691613E+00	0.7402746E-03	0.2448130E-01	0.1350850E+04	1112	1
0.5985059E+01	0.9692808E+00	0.7358456E-03	0.2434063E-01	0.1358980E+04	1113	1
0.5990058E+01	0.9694002E+00	0.7314347E-03	0.2419996E-01	0.1367176E+04	1114	1
0.5995057E+01	0.9695188E+00	0.7270572E-03	0.2406073E-01	0.1375407E+04	1115	1
0.6000056E+01	0.9696375E+00	0.7226378E-03	0.2392149E-01	0.1383713E+04	1116	1
0.6005055E+01	0.9697547E+00	0.7183896E-03	0.2378416E-01	0.1392002E+04	1117	1
0.6010054E+01	0.9698718E+00	0.7141004E-03	0.2364731E-01	0.1400363E+04	1118	1
0.6015053E+01	0.9699879E+00	0.7098601E-03	0.2351189E-01	0.1408728E+04	1119	1
0.6020052E+01	0.9701038E+00	0.7056340E-03	0.2337646E-01	0.1417165E+04	1120	1
0.6025051E+01	0.9702194E+00	0.7014291E-03	0.232400E-01	0.1425660E+04	1121	1
0.6030050E+01	0.9703341E+00	0.6972684E-03	0.2310896E-01	0.1434168E+04	1122	1
0.6035049E+01	0.9704600E+00	0.6931475E-03	0.2297688E-01	0.1442694E+04	1123	1
0.6040048E+01	0.9705855E+00	0.6890842E-03	0.2284670E-01	0.1451201E+04	1124	1
0.6045047E+01	0.9707124E+00	0.6850502E-03	0.2271748E-01	0.1459747E+04	1125	1
0.6050046E+01	0.9708381E+00	0.6810352E-03	0.2258873E-01	0.1468353E+04	1126	1
0.6055045E+01	0.9709636E+00	0.6770543E-03	0.2246094E-01	0.1476986E+04	1127	1
0.6060044E+01	0.9710893E+00	0.6731122E-03	0.2233410E-01	0.1485636E+04	1128	1
0.6065043E+01	0.9712148E+00	0.6692007E-03	0.2220821E-01	0.1494320E+04	1129	1
0.6070042E+01	0.9713331E+00	0.6653152E-03	0.2208376E-01	0.1503047E+04	1130	1
0.6075041E+01	0.9714408E+00	0.6614316E-03	0.2195883E-01	0.1511872E+04	1131	1
0.6080040E+01	0.9715483E+00	0.6576139E-03	0.2183580E-01	0.1520649E+04	1132	1
0.6085039E+01	0.9716549E+00	0.6538110E-03	0.2171326E-01	0.1529493E+04	1133	1
0.6090038E+01	0.9717613E+00	0.6500462E-03	0.2159214E-01	0.1538352E+04	1134	1
0.6095037E+01	0.9718674E+00	0.6462960E-03	0.2147102E-01	0.1547278E+04	1135	1
0.6100036E+01	0.9719725E+00	0.6425686E-03	0.2135086E-01	0.1556254E+04	1136	1
0.6105035E+01	0.9720776E+00	0.6388631E-03	0.2123213E-01	0.1565234E+04	1137	1

Compute (by hand) the nodal forces for nodes 1 and 5, i.e.

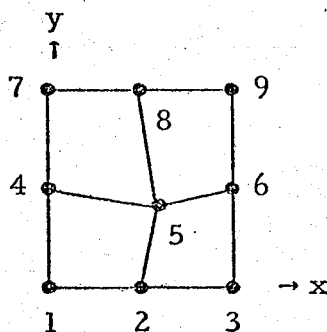
$$f_{21}^e = \int_{\Omega^e} f_2 N_1 d\Omega$$

$$f_{25}^e = \int_{\Omega^e} f_2 N_5 d\Omega$$

(Note that $f_{1a}^e = 0$, $1 \leq a \leq 8$, and, by symmetry, $f_{21}^e = f_{22}^e = f_{23}^e = f_{24}^e$ and $f_{25}^e = f_{26}^e = f_{27}^e = f_{28}^e$.)

The results of this simple computation should be somewhat surprising. Comment.

2. Consider the "patch" of elements shown below.



The nodal coordinates are:

a	x_a	y_a
1	0	0
2	1	0
3	2	0
4	0	1
5	1.1	.8
6	2	1
7	0	2
8	1	2
9	2	2

Use LEARN to execute the following

displacement boundary value problems ($1 \leq a \leq 4$, $6 \leq a \leq 9$):

test #	d_{1a}	d_{2a}
1	1	0
2	0	1
3	x_a	0
4	0	x_a
5	y_a	0
6	0	y_a

Use the following data:

- plane strain option
- $E = 1.$, $\nu = .3$

If the first letter of your last name is

$\left\{ \begin{array}{l} \text{A-H} \\ \text{I-R} \\ \text{S-Z} \end{array} \right\}$	use the	$\left\{ \begin{array}{l} 2 \times 2 \text{ quad.} \end{array} \right\}$	element.
		$\left\{ \begin{array}{l} 1\text{-pt. } \lambda \text{ term} \\ 2 \times 2 \text{ } \mu \text{ term} \end{array} \right\}$	
		$\left\{ \begin{array}{l} \text{incomp. modes} \end{array} \right\}$	

Based upon the results you obtain, determine whether or not your element passes the patch test.

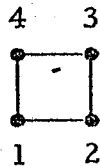
(Consider results "exact", if they are correct to 8, or more, significant digits. Due to round-off errors, one cannot expect truly exact results on digital computers.)

- The "rank check" option in the LEARN can be used to determine the number of zero eigenvalues of the global stiffness matrix $\tilde{K}$. That is, the diagonal elements of the diagonal matrix $\tilde{D}$ in the factorization

$K = U^T D U$, where U is an upper triangular matrix, are printed.

The number of zeroes in D equals the number of zero eigenvalues of K , which corresponds to the number of zero-energy modes of the structure in question. Normally, there are, at most, 3 zero-energy modes -- the rigid body modes.

Consider the single element shown below



a	x_a	y_a
1	0	0
2	1	0
3	1	1
4	0	1

Assume the data is the same as in problem 2, and there are no boundary conditions.

Perform rank check analyses for

- (a) the element you used in problem 2. $1 \times 1 \lambda \quad 2 \times 2 \mu$
- (b) 1-pt. on λ term and 1-pt. on μ term element. $1 \times 1 \lambda \quad 1 \times 1 \mu$

Comment on your results.

$$1. \quad \tilde{f} = \begin{Bmatrix} 0 \\ -g \end{Bmatrix}$$

$$f_p^e = \int N_a f_i d\Omega_e$$

$$p = \text{mod}(a-1) + i \quad \text{mod} = 2 \quad i=1,2 \quad 1 \leq a \leq 8$$

$$\text{since } f_1 = 0 \Rightarrow \forall \text{ all odd } p \quad f_p^e = 0$$

$$f_{2a}^e = \int f_2 N_a d\Omega_e = \int_{-1}^1 \int_{-1}^1 f_2 N_a \cdot \frac{h^2}{4} d\xi d\eta$$

$$f_{2a}^e = -\frac{gh^2}{4} \int_{-1}^1 \int_{-1}^1 N_a d\xi d\eta$$

$$N_1 = -\frac{1}{4}(1-\xi)(1-\eta)[1+\xi+\eta]$$

$$N_5 = \frac{1}{2}(1-\xi^2)(1-\eta)$$

$$\text{where } N_i = N_j - \frac{1}{2}(N_{i+1} + N_{i+2}) \quad \text{4 nodes quad} \quad i=1,2,3,4$$

$$N_2 = -\frac{1}{4}(1+\xi)(1-\eta)[1-\xi+\eta]$$

$$N_6 = \frac{1}{2}(1+\xi)(1-\eta^2)$$

$$N_3 = -\frac{1}{4}(1+\xi)(1+\eta)[1-\xi-\eta]$$

$$N_7 = \frac{1}{2}(1-\xi^2)(1+\eta)$$

$$N_4 = -\frac{1}{4}(1-\xi)(1+\eta)[1+\xi-\eta]$$

$$N_8 = \frac{1}{2}(1-\eta^2)(1+\xi)$$

$$\iint N_1 dA = -\frac{1}{4} \int_{-1}^1 \int_{-1}^1 [(1-\eta)\xi + (1-\eta^2)] d\eta$$

$$\left(\xi \cdot 2\eta \Big|_{\eta=-1}^1 + 2(\eta - \eta^3) \Big|_{\eta=-1}^1 \right) = -\frac{1}{4} \int_{-1}^1 (1-\xi) \cdot 2(\xi + \frac{2}{3}) d\xi = \frac{1}{2} \int_{-1}^1 (\xi - 1)(\xi + \frac{2}{3}) d\xi$$

$$= \frac{1}{2} 2 \left(\frac{\xi^3}{3} - \frac{2}{3}\xi \right) \Big|_{\xi=-1}^1 = -\frac{1}{3}$$

$$\therefore f_{2a}^e = \frac{gh^2}{12} \quad \text{for } a=1,2,3,4$$

if for N_2 we let $-\xi' = +\xi$ then we get same integral; for N_3 let $-\xi' = \xi, -\eta' = \eta$; for N_4 let $-\eta' = \eta$

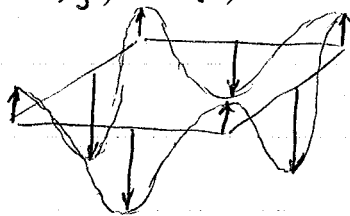
$$\iint N_5 = \frac{1}{2} \int_{-1}^1 d\xi (1-\xi^2) \int_{-1}^1 (1-\eta) d\eta = \frac{1}{2} \cdot 2 \left(\xi - \frac{\xi^3}{3} \right) \Big|_{\xi=-1}^1 \cdot 2 \cdot \eta \Big|_{\eta=-1}^1 = 2 \cdot \frac{2}{3} = \frac{4}{3}$$

$$\therefore f_{2a}^e = -\frac{gh^2}{3} \quad a=5,6,7,8$$

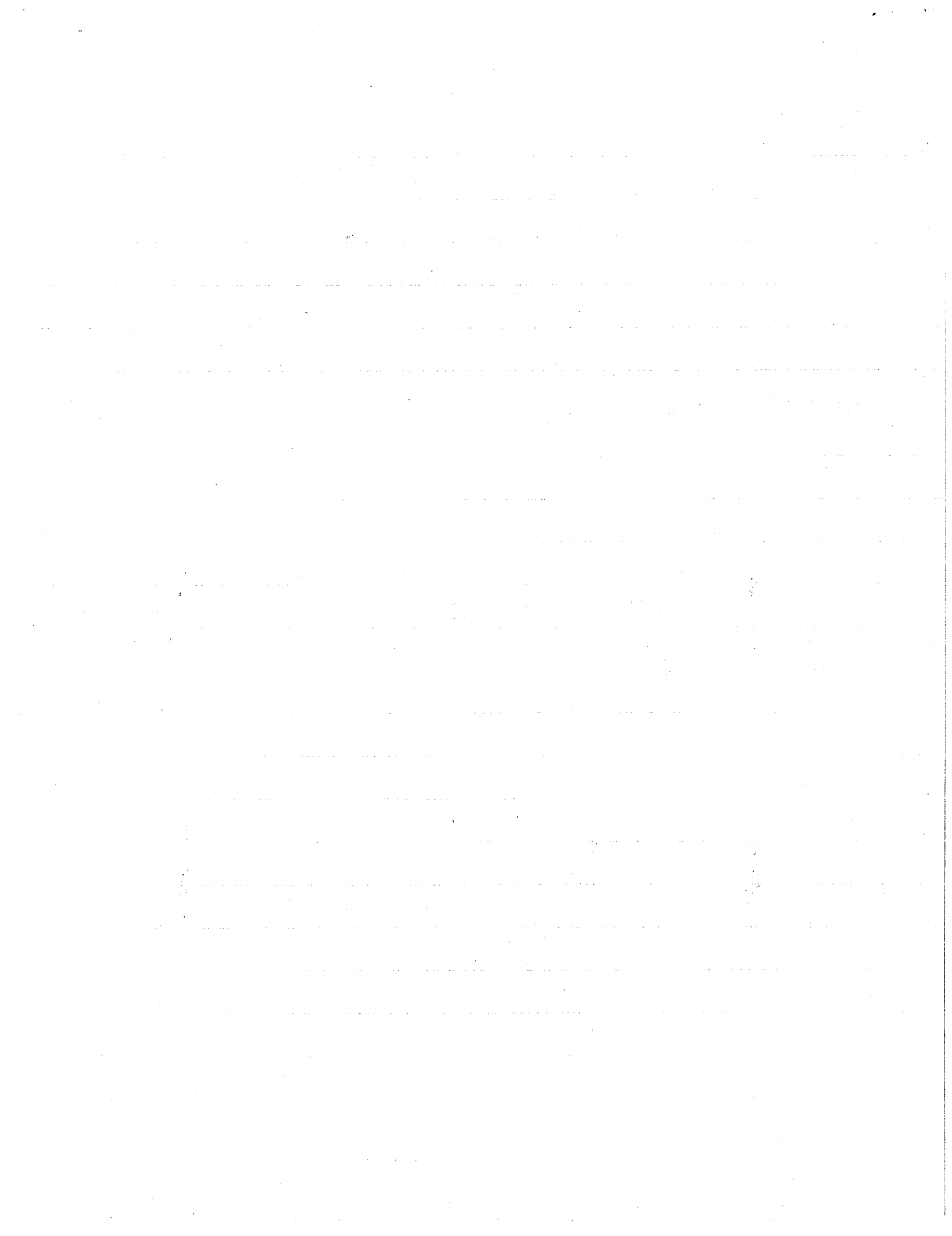
Again same trick as before shows equality of integrals. Now $\sum_{a=1}^8 N_a = 1$

$$\therefore \sum f_{2a}^e = \sum \int f_2 N_a d\Omega_e = \int f_2 \cdot 1 d\Omega_e = f_2 \Omega_e = -gh^2$$

$$\text{Now } 4 \cdot \frac{gh^2}{12} + 4 \left(-\frac{gh^2}{3} \right) = -\frac{gh^2}{3} \quad \checkmark = \sum f_{2a}^e$$



expect a distribution all of which $= |gh^2/8|$ are negative (true for linear case but not quadratic).



2. The patch test was run for the $1 \times 1 \lambda$ $2 \times 2 \mu$ element using as data $E=1.0$ $\nu=.3$ with the plane strain option. Below are the results of tests 1 thru 6.

TEST 1 $D_{1a}=1$ $D_{2a}=0$

D I S P L A C E M E N T S

NODE NUMBER	DOF 1	DOF 2
1	1.00000000D+00	0.0
2	1.00000000D+00	0.0
3	1.00000000D+00	0.0
4	1.00000000D+00	0.0
5	1.00000000D+00	-4.01290290D-19
6	1.00000000D+00	0.0
7	1.00000000D+00	0.0
8	1.00000000D+00	0.0
9	1.00000000D+00	0.0

TEST 2 $D_{1a}=0$ $D_{2a}=1$

D I S P L A C E M E N T S

NODE NUMBER	DOF 1	DOF 2
1	0.0	1.00000000D+00
2	0.0	1.00000000D+00
3	0.0	1.00000000D+00
4	0.0	1.00000000D+00
5	1.90819582E-17	1.00000000D+00
6	0.0	1.00000000D+00
7	0.0	1.00000000D+00
8	0.0	1.00000000D+00
9	0.0	1.00000000D+00

TEST 3 $D_{1a}=\chi_a$ $D_{2a}=0$

D I S P L A C E M E N T S

NODE NUMBER	DOF 1	DOF 2
1	0.0	0.0
2	1.00000000D+00	0.0
3	2.00000000D+00	0.0
4	0.0	0.0
5	1.10000000D+00	-3.77212872D-17
6	2.00000000D+00	0.0
7	0.0	0.0
8	1.00000000D+00	0.0
9	2.00000000D+00	0.0

TEST 4 $D_{1a} = 0$ $D_{2a} = X_a$

DISPLACEMENTS

NODE NUMBER	DOF1	DOF2
1	0.0	0.0
2	0.0	1.00000000D+00
3	0.0	2.00000000D+00
4	0.0	0.0
5	-6.93889390E-18	1.10000000D+00
6	0.0	2.00000000D+00
7	0.0	0.0
8	0.0	1.00000000D+00
9	0.0	2.00000000D+00

TEST 5 $D_{1a} = y_a$ $D_{2a} = 0$

DISPLACEMENTS

NODE NUMBER	DOF1	DOF2
1	0.0	0.0
2	0.0	0.0
3	0.0	0.0
4	1.00000000D+00	0.0
5	8.00000000D-01	2.40774174D-17
6	1.00000000D+00	0.0
7	2.00000000D+00	0.0
8	2.00000000D+00	0.0
9	2.00000000D+00	0.0

TEST 6 $D_{1a} = 0$ $D_{2a} = y_a$

DISPLACEMENTS

NODE NUMBER	DOF1	DOF2
1	0.0	0.0
2	0.0	0.0
3	0.0	0.0
4	0.0	1.00000000D+00
5	4.33680869E-17	8.00000000D-01
6	0.0	1.00000000D+00
7	0.0	2.00000000D+00
8	0.0	2.00000000D+00
9	0.0	2.00000000D+00

we note that the element passed the patch test with flying "colors".

3. Below are the elements of D for the $1 \times 1 \lambda$, $2 \times 2 \mu$ case as well as those for the $1 \times 1 \lambda \neq \mu$ without and with the hourglass modes.

$1 \times 1 \lambda$
 $2 \times 2 \mu$

ARRAY D (NEQ) OF FACTORIZATION

1	5.28846154D-01
2	4.19580420D-01
3	2.88461538D-01
4	2.88461538D-01
5	2.84900285D-01
6	-2.70400564D-33
7	1.11289174D-03
8	-3.97598621D-14

Based on the results of this run we show that there are at most two energy modes (the 10^{-33} and 10^{-14} entries)

$1 \times 1 \lambda \neq \mu$
no hourglass

ARRAY D (NEQ) OF FACTORIZATION

1	4.32692308D-01
2	2.99145299D-01
3 *	2.74725275D-01
4 *	6.93889390D-17
5	-1.73472348D-17
6	1.94289029D-17
7	-2.04836768D-17
8	-2.89698820D-17

Based on the results of this run we show what appears to be 5 zero eigenvalues (all entries that are 10^{-17}). Because the hourglass options were not used we question the results. If we only use the negative entries then we show the three energy modes, but this method of picking the eigenvalues is dangerous.

$1 \times 1 \lambda \neq \mu$
w/ hourglass

ARRAY D (NEQ) OF FACTORIZATION

1	5.24267399D-01
2	4.14047379D-01
3	2.87885767D-01
4	2.77915633D-01
5 *	2.74725275D-01
6 *	6.93889390D-17
7	-1.06783072D-16
8	5.45391847D-17

Based on the results of this run we show at most three zero entries (10^{-16} , 10^{-17}). However note the starred entries. These entries lead me to believe that such ^{numbers} ~~coincidence~~ may not be accidental and that 6.94×10^{-17} is a non zero entry, thus reducing the no. of zero entries to two (same as the $1 \times 1 \lambda$ $2 \times 2 \mu$ case). Also note entries here as very similar to the first run.

Hence we can say that the $1 \times 1 \lambda + \mu$ can require the use of the hourglass option and that the higher modes agree better with the $1 \times 1 \lambda$, $2 \times 2 \mu$ case. The lowest mode is somehow lost here and requires another case to be run in order to find it.

Tell me who
you are → name _____

ME 235b
Finite Element Methods
Spring Quarter 1981

Stanford University

Instructor: T. J. R. Hughes

Midterm Exam

Time = 1 Hour
Total Points = 100

1 20
2 20
3 40 10
4 20

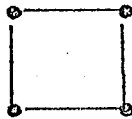
80 70

20

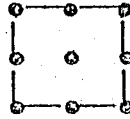
1. (20 points) A series of convergence rate studies on linear elastic boundary-value problems is performed. All results indicate that $\underline{e} = \underline{u}^h - \underline{u}$ behaves as follows:

$$\|\underline{e}\|_1 = O(h^2)$$

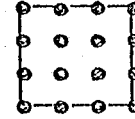
Which of the following elements was used in the studies?



4-node



9-node



16-node

Justify your answer.

mine node element - since we had done problems in class with
sobolev norms. If $\underline{x}(\underline{\xi})$ is linear we ^{were told} showed that,

$$\|\underline{e}\|_1 \leq \text{const.} \|\underline{u} - \underline{u}^h\| \leq ch^k \|\underline{u}\|_{k+1} \quad \text{where } k \text{ was the degree of completeness of polynomial}$$

then for the nine pt node we can show that a quadratic polyn.
is complete $\therefore k=2$.

approx.

$$\|u\|_m^2 \leq C_3^2 h^{2\alpha} \|u\|_5^2$$

$$s \geq 1$$

are connected

$$k=2$$

$$\|u\|_3$$

the degree of complete polyn

5x4

interp formula having 3 points

07

$$-1 \leq \frac{1}{1-\lambda \delta t}$$

$$-1 + \lambda \delta t \leq 1$$

$$\lambda \delta t \leq 2$$

$$\delta t \geq \frac{2}{\lambda}$$

$$f_n \lambda < 0$$

$$\lambda \delta t \leq 0 \leq 2$$

JOB 3525
COMMAND? loc
JOB 3525 HPFFVAR AWAITING FR...
COMMAND? fetch 3525 clr
COMMAND? list 1840/las
1840.
STATISTICS

$$u^h - u = e.$$

$$c_1 \|e\|_m^2 \leq a(e, e) \leq a(u^h - u, u^h - u) \leq c_2 \|u\|_0$$

we had shown that

$$\therefore \|e\|_m \leq \sqrt{\frac{c_2 c_3}{c_1}} h^\alpha \|u\|_5$$

$$\stackrel{m=1}{\|e\|_1} \leq \text{const } h^\alpha \|u\|_5$$

$$\begin{aligned} &\text{if } X(\xi) \text{ is linear} \quad \text{sd } \alpha \\ &\leq \text{const } h^k \|u\|_{k+1} \quad \text{if } \alpha \end{aligned}$$

ie $P(x, y) = \alpha_0 + \alpha_1 x + \alpha_2 x^2 + \alpha_3 y + \alpha_4 y^2 + \alpha$

since highest degree is 2 $\Rightarrow$ Lagrange

on a side $\therefore 9$ nodes.

2. (20 points) To determine if a new finite element for 2-dimensional heat conduction is convergent, a patch test is performed. The results are as follows:

d_a

a	x_a	y_a	Prob. 1	Prob. 2	Prob. 3
1	1.0	0.0	0.0	1.0	1
2	2.0	0.0	0.0	2.0	1
3	3.0	0.0	0.0	3.0	1
4	1.0	1.0	1.0	1.0	1
5	2.1	1.2	1.2	2.0	1
6	3.0	1.0	1.0	3.0	1
7	1.0	3.0	3.0	1.0	1
8	2.0	3.0	3.0	2.0	1
9	3.0	3.0	3.0	3.0	1

interior node —→ 5

Nodes 1-4 and 6-9 are boundary nodes.

Does the element pass the patch test? Justify your answer.

no, since x_5 for problem 2 $\neq$ x coord of node 5

2

~~20~~ 10

3. (30 points) Consider the Lagrange quadrilaterals in two-dimensional elasticity. Take the case of the "general" element with n_{es} nodes per side. (Thus the element has n_{es}^2 nodes) Determine formulas for the constraint ratio, assuming

i) selective reduced integration

ii) exact integration

What conclusions may be drawn from these results?

for selective

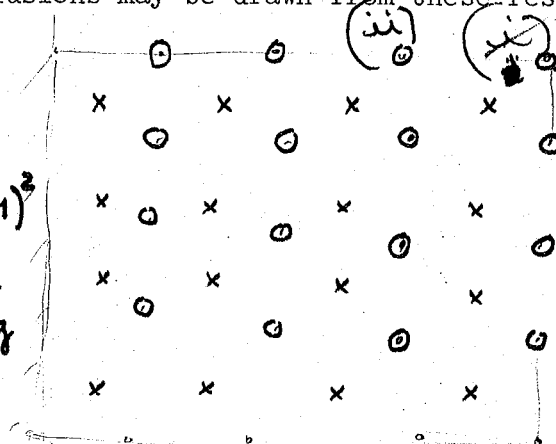
0 - 2 eqns/pt

$$\therefore n_{eq} = 2(n_{es}-1)^2$$

x - constraints for selective integ

$$n_c = (n_{es}-1)^2$$

$$CR = 2$$



for n_{es} nodes $\Rightarrow$ poly of degree $2n_{es}-1$

$\therefore$ no. of points for integ is $\frac{n_{es}}{2}$

thus we will have $(\frac{n_{es}}{2})^2 - 1$ constraints

$$2(n_{es}-1)^2 \text{ eqns}$$

there are $n_{es}^2 - 1$ constraints

$\therefore$ for regular the

$$CR = \frac{2(n_{es}-1)^2}{n_{es}^2 - 1}$$

$$CR = 2$$

for reduced the

$$CR = 2(n_{es}-1)^2$$

(i) ~~(i)~~
(ii) ~~(ii)~~

for selective integ

for n_{es} nodes there are $2(n_{es}-1)^2$ eqns

there are $(n_{es}-1)$ integ points side

1-less

$$\therefore CR = \frac{2(n_{es}-1)^2}{2(n_{es}-1)} = n_{es}-1$$

$$\therefore n_c = (n_{es}-1)^2$$

degree $2(n_{es}-1)-1$ = Polyn for selec integ $\Rightarrow$ it has $2(n_{es}-1)$ coeff has $2n_{es}-2-1$ constraints

selective integration has a better constraint ratio

-10

for exact integ

0 - 2 eq/pt

$$\therefore n_{eq} = 2(n_{es}-1)^2$$

$$n_{es} = n_{es}^2 - 1$$

$$\therefore CR = \frac{2(n_{es}-1)^2}{n_{es}^2 - 1}$$

$$= 2(n_{es}-1)$$

for n_{es} points

lagr. poly is of degree $n_{es}-1$ & has n_{es} terms

$\therefore$ poly of $(n_{es}-1)$ in x, y has n_{es}^2 terms

$$\text{and } n_c = n_{es}^2 - 1 \Leftarrow \nabla \cdot u = 0$$

4. (30 points) Consider the following "fractional-step" algorithm for the model heat equation:

$$\frac{d_{n+1/2} - d_n}{(\Delta t/2)} + \lambda d_n = 0$$

$$\frac{d_{n+1} - d_{n+1/2}}{(\Delta t/2)} + \lambda d_{n+1} = 0$$

- Obtain an expression for the amplification factor.
- Determine the conditions under which the algorithm is stable.
- What is the rate-of-convergence of this algorithm with respect to solutions of $\dot{d} + \lambda d = 0$.

$$d_{n+1/2} - d_n + \lambda \Delta t/2 d_n = 0$$

$$d_{n+1} - d_{n+1/2} + \lambda \Delta t/2 d_{n+1} = 0$$

$$d_{n+1} + \lambda \Delta t/2 d_{n+1} - d_n + \lambda \Delta t/2 d_n = 0$$

$$d_{n+1} (1 + \lambda \Delta t/2) = d_n (1 - \lambda \Delta t/2)$$

$$\therefore (i) \quad A = \frac{1 - \lambda \Delta t/2}{1 + \lambda \Delta t/2}$$

(ii) if $|A| = \left| \frac{1 - \lambda \Delta t/2}{1 + \lambda \Delta t/2} \right| < 1$ then this is unconditionally stable if $\lambda, \Delta t > 0$

$$-1 < \frac{1 - \frac{2\lambda \Delta t}{2}}{1 + \lambda \Delta t/2} < 1$$

(ii) ✓

$$\therefore -2\lambda \Delta t/2 < (1 + \lambda \Delta t/2) \cdot 0 = 0 \quad \text{satisfied if } \Delta t, \lambda > 0 \text{ or if } \Delta t, \lambda < 0$$

$$-1 < 1 - \frac{2\lambda \Delta t/2}{1 + \lambda \Delta t/2} \Rightarrow 1 > \frac{\lambda \Delta t/2}{1 + \lambda \Delta t/2} \quad \text{or} \quad 1 + \lambda \Delta t/2 > \lambda \Delta t/2$$

$\therefore |A| < 1$ then we get stability; for our normal system $\lambda, \Delta t > 0$ as "conditions"

say unconditionally stable

(iii) This system is equivalent to the $\alpha = 1/2$ condition for $\dot{d} + \lambda d = 0$

so what is ans.?

(generalized trap. approx. of)

For $\alpha = 1/2 \quad |T| < C \Delta t^{k+1} \quad k=2$

Midterm Solutions

$$1. \quad \|e\|_1 = \|u^h - u\|_1 \leq C h^k \|u\|_{k+1}$$

$$k=2 \Rightarrow 9\text{-node element}$$

$$\text{N.B. } e = u^h - u \quad \underline{\text{not}} \quad U^h - u$$

$$\|u^h - u\|_1 \leq \text{const } \|U^h - u\|_1 \leq \text{const } h^k \|u\|_{k+1}$$

best approx
||e||₁

$$2. \quad \text{Prob. 1} \quad d_a = y_a \quad \text{OK}$$

$$\text{Prob. 2} \quad d_a = \lambda a \quad \text{no good, i.e. } d_5 = 2.1 \neq 2.0$$

$$\text{Prob. 3} \quad d_a = 1 \quad \text{OK}$$

ans. NO

$$3. (i) \quad n_{eq} = 2(n_{ex} - 1)^2$$

$$n_c = (n_{ex} - 1)^2 \quad \text{"standard" reduced quadrature of } \lambda\text{-term}$$

(see handout sheet)

$$\therefore CR = 2 \quad (\text{selective reduced integration})$$

$$(ii) \quad n_c = n_{ex}^2 - 1 \quad (\text{exact int.})$$

$$CR = \frac{2(n_{ex} - 1)^2}{n_{ex}^2 - 1} = \frac{2(n_{ex} - 1)}{n_{ex} + 1}$$

3. (cont'd)

For exact integration, CR improves the higher the order of the element. For example,

n_{es}	CR
2	2/3
3	1
4	6/5
$\vdots$	$\vdots$
∞	2

Thus the necessity of reduced integration decreases as the order of the element increases. (This is observed in practice.)

$$4. \quad d_{n+1/2} - d_n + \frac{\lambda \Delta t}{2} d_n = 0 \quad (1)$$

$$d_{n+1} - d_{n+1/2} + \frac{\lambda \Delta t}{2} d_{n+1} = 0 \quad (2)$$

(ii) solve (1) for $d_{n+1/2}$ and eliminate from (2) to arrive at:

$$d_{n+1} = A d_n ; \quad A = \frac{1 - \lambda \Delta t / 2}{1 + \lambda \Delta t / 2}$$

4. (cont'd) Observe that the foregoing A is identical to the one obtained for $\alpha = 1/2$ (i.e. trapezoidal rule) in the generalized trapezoidal family (see class notes)

(ii) unconditionally stable

(iii) $k=2$ (rate of conv.)

11

ME 235B
Finite Element Methods
Spring Quarter 1981

Stanford University
Instructor: T. Hughes

Final Exam
Time = 3 Hours
Total Points = 300

1	50
2	60
3	75
4	10
5	60
	255

Course A+

Instructions: Open notes and homework are allowed, You may use all results previously obtained in the course lectures and homework, so be as brief as possible.

CESAR LEVY
Name

1. (60 points) The initial-boundary-value problem for an elastic membrane on a Winkler foundation is stated as follows:

Given $f: \Omega \times]0, T[\rightarrow \mathbb{R}$, $g: \Gamma_g \times]0, T[\rightarrow \mathbb{R}$, $h: \Gamma_h \times]0, T[\rightarrow \mathbb{R}$, and $w_0: \Omega \rightarrow \mathbb{R}$, and $\dot{w}_0: \Omega \rightarrow \mathbb{R}$; find $w: \Omega \times]0, T[\rightarrow \mathbb{R}$ such that

$$\begin{aligned} w_{,ii} - \alpha w + f &= w_{,tt} \text{ on } \Omega \times]0, T[\\ w &= g & x \in \Gamma_g &, \quad t \in]0, T[\\ w_{,n} &= h & x \in \Gamma_h &, \quad t \in]0, T[\\ w &= w_0 & x \in \Omega &, \quad t = 0 \\ w_{,t} &= \dot{w}_0 & x \in \Omega &, \quad t = 0 \end{aligned}$$

where $w_{,n} = \partial w / \partial n$ is the derivative in the direction of the outward unit normal vector and $\alpha > 0$.

Set up the Following finite element paraphernalia:

- A weak form of the problem (i.e. "W") in which the h -boundary condition is "natural."
- The galerkin form of the problem (i.e. "G").
- The matrix initial-value ordinary differential equation problem (i.e. "M").

2. (70 points) The problem under consideration consists of

$$\tilde{M} \dot{\tilde{d}} + \tilde{K} \tilde{d} = \tilde{F}$$

$$\tilde{d}(0) = \tilde{d}_0$$

where $\tilde{M}$ and $\tilde{K}$ are symmetric and positive definite matrices

$\tilde{F} = \tilde{F}(t)$ and $\tilde{d}_0$ are given. Consider the following one-parameter

(α) family of predictor-corrector algorithms:

$$\tilde{M} \tilde{v}_{n+1} + \tilde{K} \tilde{d}_{n+1} = \tilde{F}_{n+1}$$

$$\tilde{d}_{n+1} = \tilde{d}_n + \Delta t(1 - \alpha) \tilde{v}_n$$

$$\tilde{d}_{n+1} = \tilde{d}_{n+1} + \Delta t \alpha \tilde{v}_{n+1}$$

Assume $\alpha \in]0, 1]$.

- Determine an expression for the amplification factor.
- Determine under what circumstances the algorithms are stable.
- Obtain an expression for the local truncation error.
- Determine the rate of convergence.

Hint The equations may be combined to form

$$(\tilde{M} - \alpha \Delta t \tilde{K}) \tilde{v}_{n+1} + \tilde{K} \tilde{d}_{n+1} = \tilde{F}_{n+1}$$

which is useful in the analysis.

3. Consider the second-order ordinary differential equation

$$\ddot{d} + \omega^2 d = 0 \quad (1)$$

(a) (25 points) The forward Euler method for (1) consists of the following equations:

$$a_n + \omega^2 d_n = 0$$

$$d_{n+1} = d_n + \Delta t v_n$$

$$v_{n+1} = v_n + \Delta t a_n$$

Set up the 2 x 2 amplification matrix for this method and determine the stability behavior.

(b) (50 points) The backward Euler method for (1) consists of:

$$a_{n+1} + \omega^2 d_{n+1} = 0$$

$$d_{n+1} = d_n + \Delta t v_{n+1}$$

$$v_{n+1} = v_n + \Delta t a_{n+1}$$

Determine the stability behavior and order-of-accuracy of this method.

4. (35 points) Consider a 4-node bilinear element which is used in an elastodynamics calculation. Assume that the Jacobian determinant is constant and equal to $A^e/4$, where A^e is the element area. The quadrature rule used to evaluate the element mass matrix is the "product trapezoidal rule," namely

ℓ	$\tilde{\xi}_\ell$	$\tilde{\eta}_\ell$	W_ℓ
1	-1	-1	1
2	1	-1	1
3	1	1	1
4	-1	1	1

Obtain exact expressions for all entries in the element mass matrix.

5. (60 points) Consider the 4-node, bilinear, isoparametric quadrilateral element. Show that $u_{i,i}^h$ at $\xi = \eta = 0$ is the mean value of $u_{i,i}^h$ over the element. That is, show

$$0 = \int_{-1}^{+1} \int_{-1}^{+1} \left(u_{i,i}^h(\xi, \eta) - u_{i,i}^h(0, 0) \right) j(\xi, \eta) d\xi d\eta$$

20
1. (a) Given f, g, h, w_0 and $\dot{w}_0$ as in (5) Find: w s.t. $\forall u \in V$ (ie $w = v + g$)
where $w: [0, T] \times \bar{\Omega} \rightarrow S_{gh}$ and $u = 0$ on $\Gamma_g \forall t \in [0, T]$

$$\int_{\Omega} (u \nabla^2 w - \alpha w u + f u - w_{,tt} u) d\Omega = 0$$

$$\int_{\Omega} [(u w_{,i})_i - u_{,i} w_{,i}] d\Omega = \int_{\Omega} \alpha w u d\Omega - \int_{\Omega} f u d\Omega + \int_{\Omega} u w_{,tt} d\Omega$$

$$\int_{\Gamma} (n \cdot \nabla w) u d\Gamma - \int_{\Omega} u_{,i} w_{,i} d\Omega = \dots$$

$$\int_{\Gamma} w_{,n} u d\Gamma - \int_{\Omega} w_{,i} u_{,i} d\Omega = \dots$$

now $w_{,n} = h$ on Γ_h

$$\therefore \int_{\Omega} w_{,i} u_{,i} d\Omega = \int_{\Gamma_h} h u d\Gamma - \int_{\Omega} \alpha w u d\Omega + \int_{\Omega} f u d\Omega - \int_{\Omega} u w_{,tt} d\Omega$$

$$\rightarrow \left\{ (u, \ddot{w}) + a(w, u) + (\alpha w, u) = (u, h)_{\Gamma_h} + (u, f) \right\} \checkmark$$

$$\text{where } a(w, u) = \int_{\Omega} w_{,i} u_{,i} d\Omega \quad (u, h)_{\Gamma_h} = \int_{\Gamma_h} h u d\Gamma \quad (A, B) = \int_{\Omega} A B d\Omega$$

$$\left\{ \begin{array}{l} \text{also } (u, w(0) - w_0) = 0 \\ \text{also } (u, \dot{w}(0) - \dot{w}_0) = 0 \end{array} \right\} \checkmark$$

20
1. (b) Given $f, g, h, w_0, \dot{w}_0$ as in (5)

Find a $v^h: [0, T] \rightarrow V^h \subset V$ where $w^h = v^h + g^h, g^h: [0, T] \rightarrow S^h$

s.t. $\forall u^h \in V^h$ then

$$(\alpha v^h, u) + (u^h, \ddot{v}^h) + a(v^h, u) = (u^h, h)_{\Gamma_h} + (u^h, f) - (\alpha g^h, u) - (u^h, \ddot{g}^h) - a(g^h, u)$$

$$\text{and } (u^h, v(0)^h) = (u^h, w_0) - (u^h, g(0)) \checkmark$$

$$(u^h, \dot{v}(0)^h) = (u^h, \dot{w}_0) - (u^h, \dot{g}(0)) \checkmark$$

10
1. (c) let $g^h(x, t) = \sum_{\alpha \in \eta_g} N_{\alpha}(x) g_{\alpha}(t) \in S^h; v^h(x, t) = \sum_{\alpha \in \eta_g} N_{\alpha}(x) d_{\alpha}(t) \in V^h$ find $d: [0, T] \rightarrow \mathbb{R}$
then letting $\tilde{M} = \tilde{A}^{\eta_g} (m_{pq}^e)$ $m_{pq}^e = [m_{pq}^e]$ $m_{pq}^e = \int_{\Omega^e} N_{\alpha} N_{\beta} d\Omega$ $p = \text{node } i$ a
 $K = \tilde{A}^{\eta_g} (k_{\alpha}^e)$ $k_{\alpha}^e = [k_{pq}^e]$ $k_{pq}^e = \int_{\Omega^e} \frac{B^T D B}{\Delta} d\Omega$ $\leftarrow N_{\alpha,i} N_{\beta,i}$ $q = \text{node } j$ b



$$\underline{F}(t) = \underline{A}^{nel} \left(\underline{\ddot{f}}^e(t) \right) \quad \underline{f}^e = [f_p^e] \quad \underline{f}_p^e = \int_{\Omega^e} N_a f d\Omega + \int_{\Gamma_e} N_a h d\Gamma - \sum_{q=1}^{n_{el}} \left(k_{pq}^e g_q^e + c_{pq}^e \dot{g}_q^e + m_{pq}^e \ddot{g}_q^e \right)$$

where $\underline{C}^e = [c_{pq}^e]$ $\underline{C} = \underline{A}^{nel}(\underline{C}^e)$ $c_{pq}^e = \int_{\Omega^e} N_a N_b d\Omega$

if $\alpha = \text{const}$ then $c_{pq}^e = \alpha m_{pq}^e$ ✓

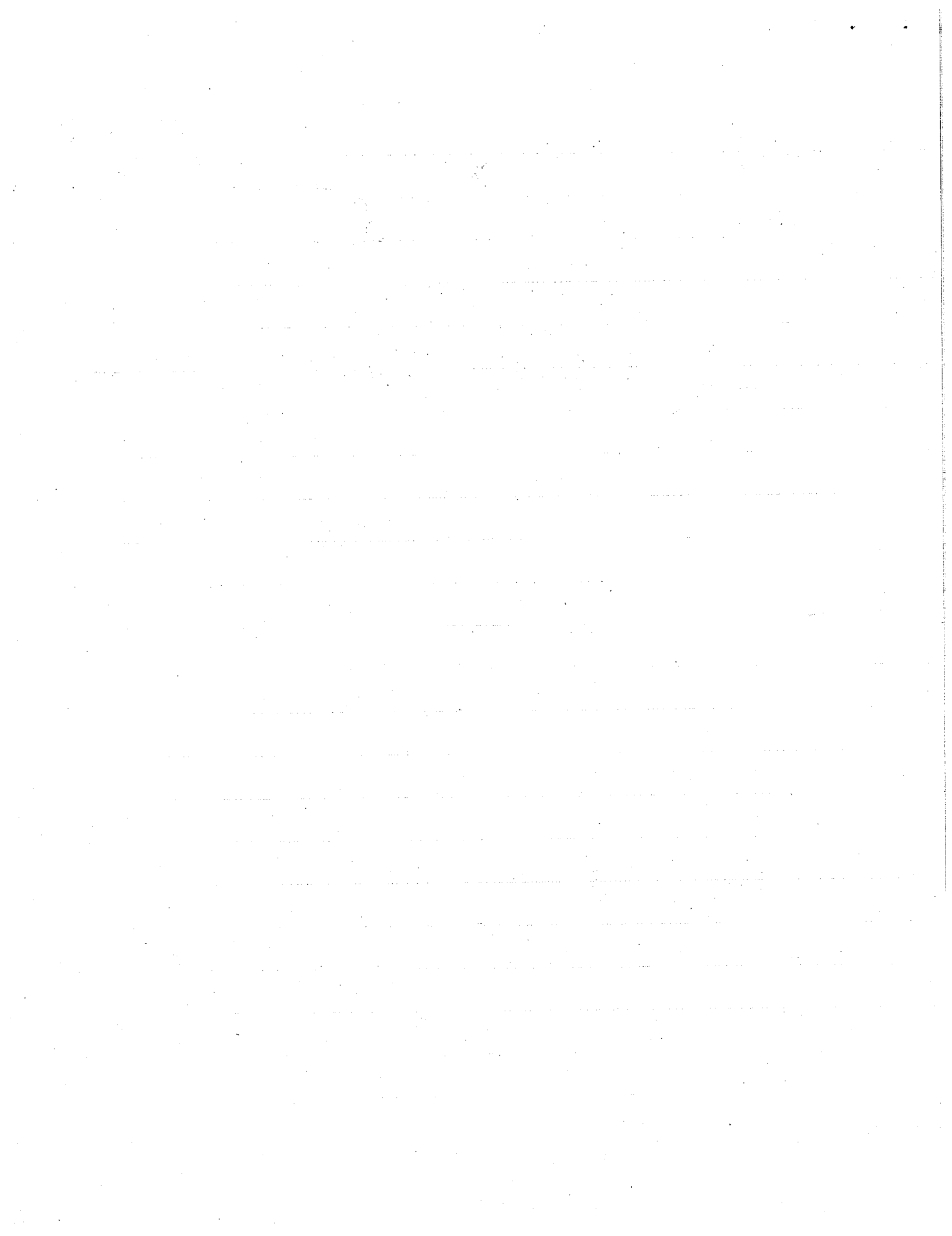
∴ our eqns become

$$\underline{M} \underline{\ddot{d}} + (\underline{C} + \underline{K}) \underline{\dot{d}} = \underline{F} \quad \checkmark$$

and $\underline{d}_0 = \underline{\bar{M}}^{-1} \left\{ \underline{A}^{nel}(\underline{\bar{d}}^e) \right\} \quad \underline{\bar{M}} = \underline{A}(\underline{\bar{m}}^e) \quad m_{pq}^e = \int_{\Omega^e} N_a N_b d\Omega$

$$\underline{\bar{d}}^e = \{ \bar{d}_p^e \} \quad \bar{d}_p^e = \int_{\Omega^e} N_a u_{0,i} d\Omega - \sum_q \bar{m}_{pq}^e g_q^e(0)$$

$$\underline{\dot{d}}_0 = \underline{\bar{M}}^{-1} \left\{ \underline{A}^{nel}(\underline{\bar{d}}^e) \right\} \quad \text{where} \quad \bar{d}_p^e = \int_{\Omega^e} N_a \dot{u}_{0,i} d\Omega - \sum_q \bar{m}_{pq}^e \dot{g}_q^e(0)$$



$$2. \quad (M - \alpha \Delta t K) \underline{v}_{n+1} + K \underline{d}_{n+1} = \underline{F}_{n+1}$$

now $M \underline{v}_n + K \underline{\tilde{d}}_n = \underline{F}_n$

$$M \underline{v}_n + K [\underline{d}_n + \Delta t (1-\alpha) \underline{v}_n] = \underline{F}_n$$

$$[M + K \Delta t (1-\alpha)] \underline{v}_n + K \underline{d}_n = \underline{F}_n \quad (2)$$

and $\underline{d}_{n+1} - \Delta t \alpha \underline{v}_{n+1} = \underline{d}_n + \Delta t (1-\alpha) \underline{v}_n$

$$\Delta t \alpha \underline{v}_{n+1} = \underline{d}_{n+1} - \underline{d}_n - \Delta t (1-\alpha) \underline{v}_n$$

$$(M - \alpha \Delta t K) [\underline{d}_{n+1} - \underline{d}_n - \Delta t (1-\alpha) \underline{v}_n] + K \Delta t \alpha \underline{v}_{n+1} = \underline{F}_{n+1} \Delta t \alpha \quad (1)$$

applying the fact that $\underline{\psi}_\ell^T K \underline{\psi}_m = \lambda_\ell \delta_{\ell m}$ & $\underline{\psi}_\ell^T M \underline{\psi}_m = \delta_{\ell m}$

$$\underline{d}_{n+1} = \sum_m \underline{d}_{n+1}(t) \underline{\psi}_m \quad \underline{d}_n = \sum_m \underline{d}_n(t) \underline{\psi}_m \quad \dot{\underline{d}}_n = \sum_m \dot{\underline{d}}_n(t) \underline{\psi}_m$$

then we get $[\underline{d}_{\ell,n+1} (1 - \alpha \Delta t \lambda_\ell) - \underline{d}_{\ell,n} (1 - \alpha \Delta t \lambda_\ell) - \Delta t (1-\alpha) \dot{\underline{d}}_{\ell,n} (1 - \alpha \Delta t \lambda_\ell)]$

$$+ \Delta t \alpha \lambda_\ell \underline{d}_{\ell,n+1} = \underline{F}_{\ell,n+1} \Delta t \alpha \quad \text{for } (1)$$

and $[1 + \lambda_\ell \Delta t (1-\alpha)] \dot{\underline{d}}_{\ell,n} + \lambda_\ell \underline{d}_{\ell,n} = \underline{F}_{\ell,n} \quad \text{for } (2)$

then $\dot{\underline{d}}_{\ell,n} = \frac{\underline{F}_{\ell,n} - \lambda_\ell \underline{d}_{\ell,n}}{1 + \lambda_\ell \Delta t (1-\alpha)}$ solving for $\dot{\underline{d}}_{\ell,n}$

Put into (1)

$$\therefore \underline{d}_{\ell,n+1} (1 - \alpha \Delta t \lambda_\ell + \Delta t \alpha \lambda_\ell) - \underline{d}_{\ell,n} (1 - \alpha \Delta t \lambda_\ell) - \Delta t (1-\alpha) (1 - \alpha \Delta t \lambda_\ell) \frac{(\underline{F}_{\ell,n} - \lambda_\ell \underline{d}_{\ell,n})}{1 + \lambda_\ell \Delta t (1-\alpha)} = \underline{F}_{\ell,n+1} \Delta t \alpha$$

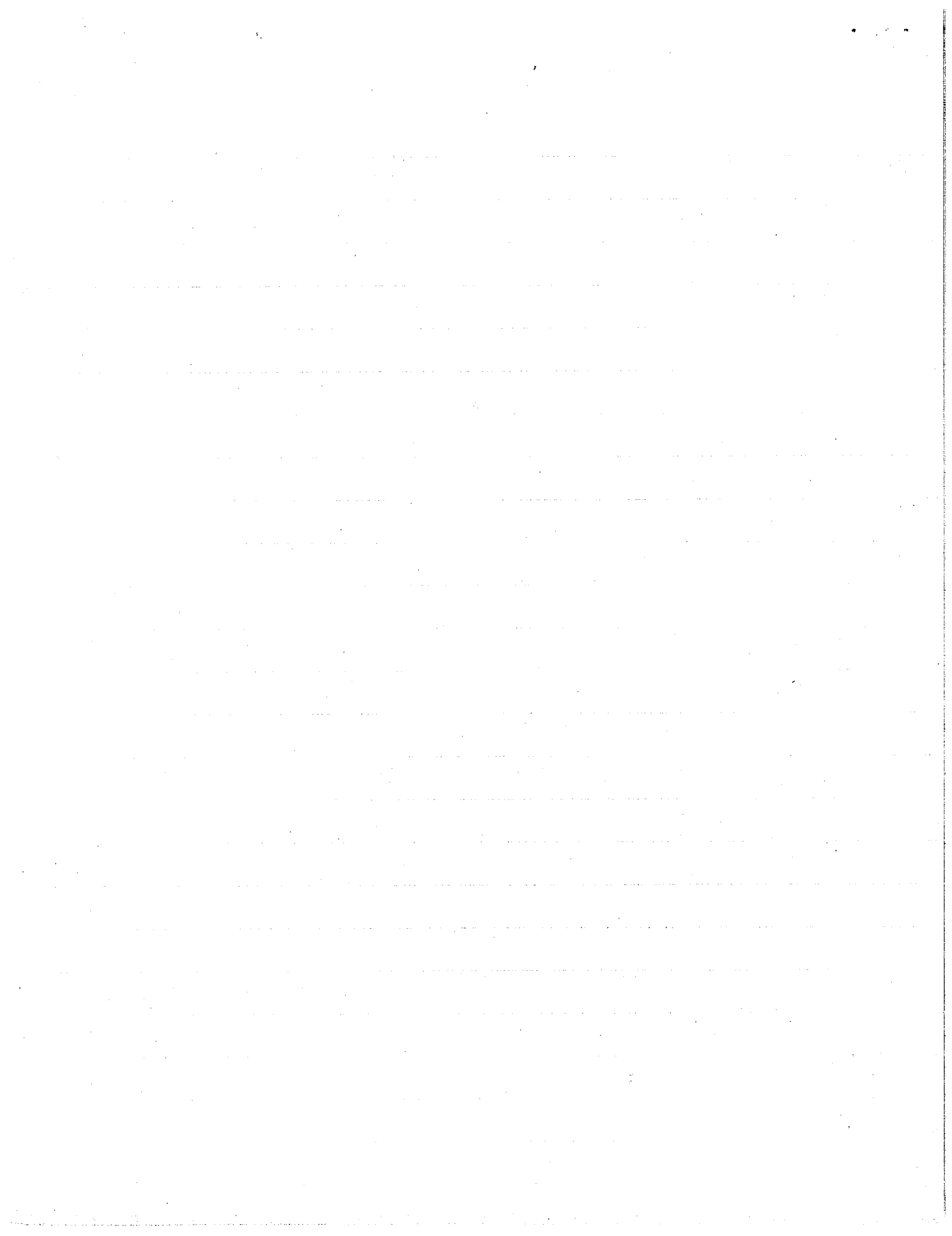
$$\underline{d}_{\ell,n+1} = \underline{d}_{\ell,n} \left\{ (1 - \alpha \Delta t \lambda_\ell) - \frac{\lambda_\ell \Delta t (1-\alpha) (1 - \alpha \Delta t \lambda_\ell)}{1 + \lambda_\ell \Delta t (1-\alpha)} \right\} + \underline{F}_{\ell,n+1} \Delta t \alpha + \underline{F}_{\ell,n} \frac{\Delta t (1-\alpha) (1 - \alpha \Delta t \lambda_\ell)}{1 + \lambda_\ell \Delta t (1-\alpha)}$$

$$\underline{d}_{\ell,n+1} = \underline{d}_{\ell,n} \left\{ \frac{(1 - \alpha \Delta t \lambda_\ell) [1 + \lambda_\ell \Delta t (1-\alpha)] - \lambda_\ell \Delta t (1-\alpha)}{1 + \lambda_\ell \Delta t (1-\alpha)} \right\} + \underline{L}_n$$

$$\underline{d}_{\ell,n+1} = \underline{d}_{\ell,n} \left[\frac{1 - \alpha \Delta t \lambda_\ell}{1 - \alpha \lambda_\ell \Delta t} \cdot (1 + \lambda_\ell \Delta t) \right] + \underline{L}_n$$

(a) ³⁰ $A = 1 - \lambda \Delta t$ ✓

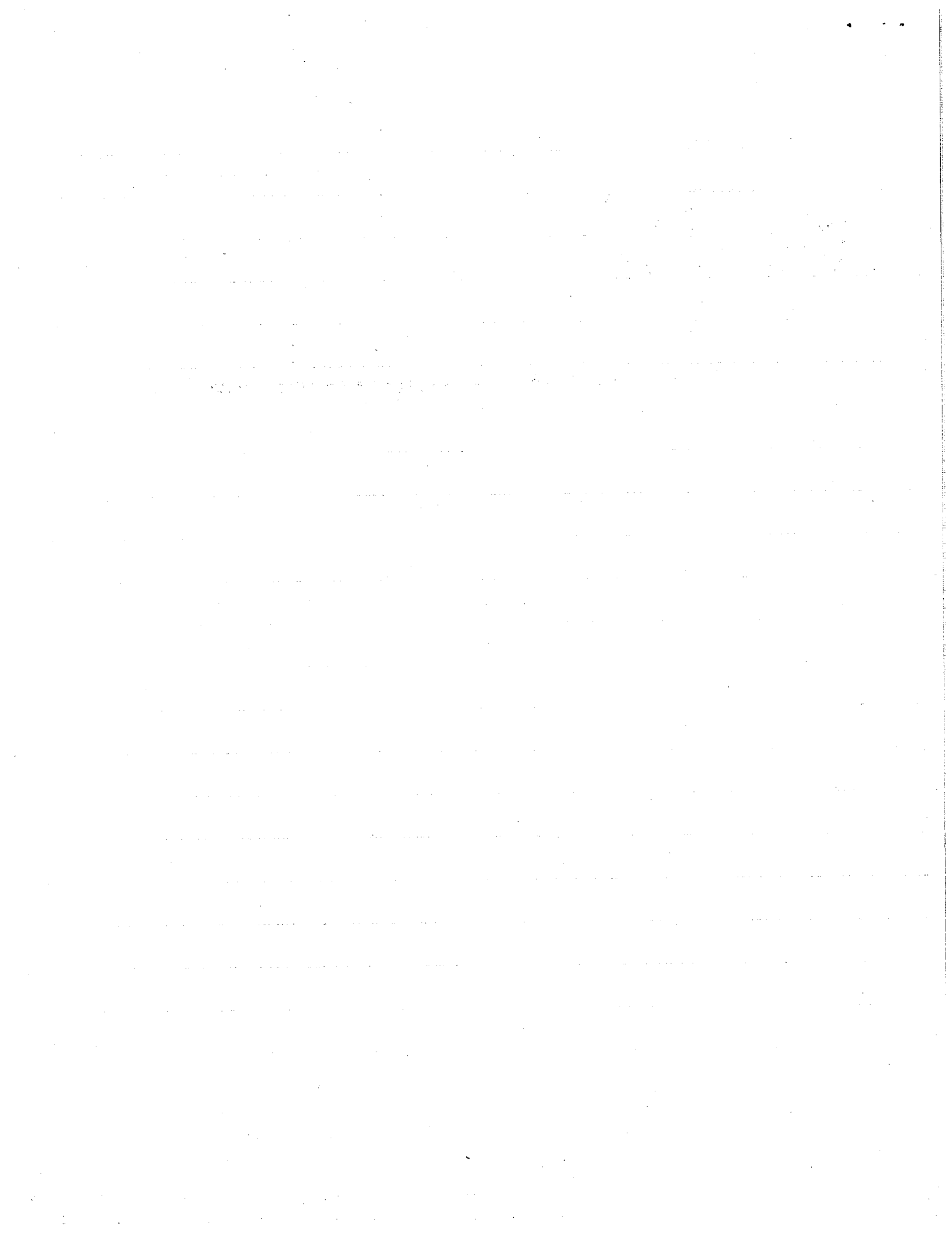
(b) ¹⁵ if $|A| < 1$ i.e. $-1 < 1 - \lambda \Delta t < 1 \Rightarrow \Delta t < \frac{2}{\lambda}$ for $\lambda, \Delta t > 0$



(c) if $\lambda \Delta t > 0$ $\lambda \Delta t < 2$ or $\boxed{\Delta t < \frac{2}{\lambda}}$ ie $\alpha = 0$

~~(d)~~ (c) $|\tau| \leq c \Delta t^2$ (d) $k=1$ ¹⁰
5 ✓ ✓

wanted an expression -10



75

3. $a_n + \omega^2 d_n = 0 \Rightarrow a_n = -\omega^2 d_n$

(a) 75
$$\begin{pmatrix} d_{n+1} \\ v_{n+1} \end{pmatrix} = \begin{pmatrix} d_n + \Delta t v_n \\ v_n - \Delta t \omega^2 d_n \end{pmatrix} = \begin{pmatrix} 1 & \Delta t \\ -\Delta t \omega^2 & 1 \end{pmatrix} \begin{pmatrix} d_n \\ v_n \end{pmatrix}$$

$$A = \begin{pmatrix} 1 & \Delta t \\ -\Delta t \omega^2 & 1 \end{pmatrix}$$

$$A_1 = \frac{1}{2} \text{tr } A = \frac{1}{2}(2) = 1$$

$$A_2 = \det A = 1 + \Delta t^2 \omega^2$$

$$\lambda_{\pm} = A_1 \pm \sqrt{A_1^2 - A_2} = 1 \pm \sqrt{1 - 1 + \Delta t^2 \omega^2} = 1 \pm i \omega \Delta t \quad \checkmark$$

$|\lambda|^{\frac{1}{2}} = (1 + \omega^2 \Delta t^2)^{\frac{1}{2}} > 1 \quad \checkmark \quad \Delta t \neq 0 \quad \omega \neq 0 \quad \therefore \text{system blow-up}$
not stable. even though $A_1 \leq (A_2 + 1)/2 \quad A_2 \neq 1.$

(b) 50 $a_{n+1} + \omega^2 d_{n+1} = 0 \quad a_{n+1} = -\omega^2 d_{n+1}$

$$d_{n+1} = d_n + \Delta t v_{n+1}$$

$$v_{n+1} = v_n + (\Delta t \omega^2 d_{n+1})$$

$$\Rightarrow \begin{pmatrix} 1 - \Delta t \omega^2 & \Delta t \\ \Delta t \omega^2 & 1 \end{pmatrix} \begin{pmatrix} d_{n+1} \\ v_{n+1} \end{pmatrix} = \begin{pmatrix} d_n \\ v_n \end{pmatrix}$$

$$A^{-1} = \begin{pmatrix} 1 & -\Delta t \omega^2 \\ \Delta t \omega^2 & 1 \end{pmatrix}$$

$$A = \frac{\begin{pmatrix} 1 & \Delta t \\ -\omega^2 \Delta t & 1 \end{pmatrix}}{1 + \Delta t^2 \omega^2}$$

$$A_1 = \frac{1}{2} \text{tr } A = \frac{1}{1 + \Delta t^2 \omega^2}$$

$$A_2 = \det A = 1$$

$$\therefore \lambda_{\pm} = \frac{1}{1 + \Delta t^2 \omega^2} \pm \sqrt{\frac{1 - (1 + \Delta t^2 \omega^2)^{-2}}{(1 + \Delta t^2 \omega^2)^2}}$$

$$\frac{1}{1 + \Delta t^2 \omega^2} \pm i \sqrt{\frac{(\Delta t^2 \omega^2)(\omega^2 \Delta t^2 + 2)}{(1 + \Delta t^2 \omega^2)^2}} \quad \text{Complex}$$

$\checkmark$ since $-1 < A_1 < 1$ and $A_2 = 1$ we have spectral stability
 since backward euler $\alpha = 1 \quad \therefore |\tau(t)| \leq C \Delta t^{k+1}$ where $k=1$

10



-10

$$m_e = \frac{1}{4} \int_{-1}^1 \int_{-1}^1 N_A N_B d\xi d\eta = \frac{A}{4} \left\{ N_A N_B \text{ evaluated at the 4 points} \times \text{weights} \right\}$$

-15 δ_{ij}

$$\text{where } N_1 = \frac{(1+\xi)(1+\eta)}{4}$$

$$N_2 = \frac{(1+\xi)(1-\eta)}{4}$$

$$N_3 = \frac{(1-\xi)(1+\eta)}{4}$$

$$N_4 = \frac{(1-\xi)(1-\eta)}{4}$$

$$\begin{matrix} (-1, -1) \\ 1 \end{matrix}$$

$$\begin{matrix} (-1, 1) \\ 0 \end{matrix}$$

$$\begin{matrix} (1, -1) \\ 0 \end{matrix}$$

$$\begin{matrix} (1, 1) \\ 0 \end{matrix}$$

$$0$$

$$1$$

$$0$$

$$0$$

$$0$$

$$0$$

$$1$$

$$0$$

$$0$$

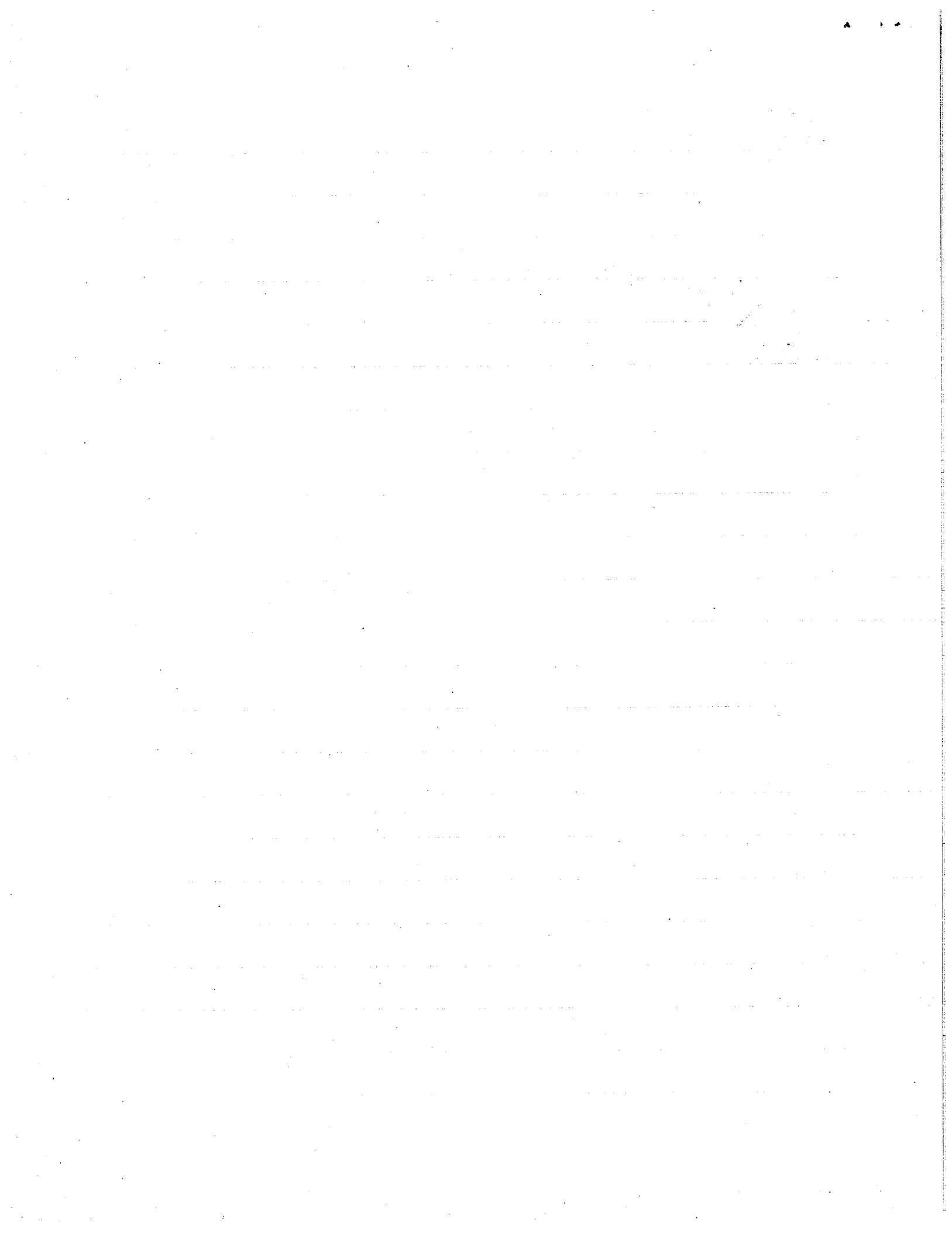
$$0$$

$$0$$

$$1$$

$$m_e = \frac{A}{4} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

didn't understand problem ^{well} needed
to ask question on it



5. 60 $u_{1,1} = \sum_{a=1}^4 N_{a,x} dx_a$ $u_{2,2} = \sum_{a=1}^4 N_{a,y} dy_a$
 $= \sum_{a=1}^4 [N_{a,\xi} y_{2,\eta} - N_{a,\eta} y_{2,\xi}] / j dx_a$ $\sum_{a=1}^4 - (N_{a,\xi} x_{1,\eta} - N_{a,\eta} x_{1,\xi}) / j dy_a$

$$\therefore \iint (u_{1,1}^n + u_{2,2}^n) f \, d\xi \, d\eta = \iint \sum_{a=1}^4 N_{a,\xi} (x_{a,\eta} - x_{\xi,\eta}) + \sum N_{a,\eta} (x_{a,\xi} - x_{\xi,\eta})_a \, d\xi \, d\eta$$

but $y_{\eta} = \sum_{a=1}^4 N_a(\xi)_{,\eta} y_a$ $x_{,\eta} = \sum N_a(\xi)_{,\eta} x_a$ etc.

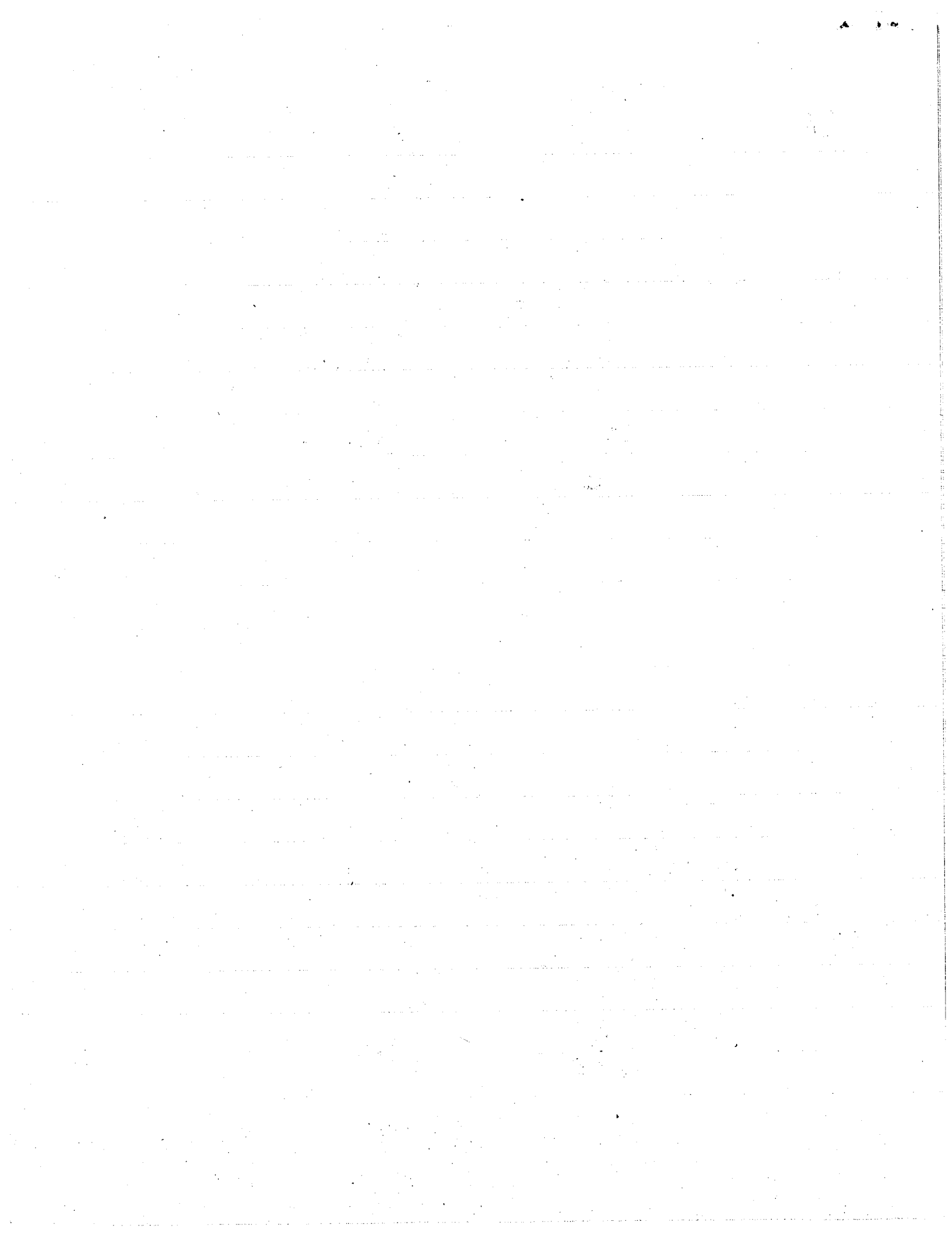
$$\iint \nabla_{\vec{r}} \cdot \vec{j} d\vec{r} d\vec{r}' = \iint d\vec{r} d\vec{r}' \sum_{a=1}^A N_{a,\beta} \sum_{b=1}^A (d_{xa} N_{b,\gamma} Y_b - d_{ya} N_{b,\beta} X_b) + \sum N_{a,\gamma} \sum (d_{ya} N_{b,\beta} X_b - d_{xa} N_{b,\gamma} Y_b)$$

Now since $N_{\alpha}(\xi)$ are linear in ξ . Then $N_{\alpha,\beta}$ will be linear in η and vice versa $N_{\alpha,\eta}$ will be linear in ξ . When integrated from $(-1, 1)$ the linear term is quadratic and gives no contrib. $\therefore$ this is same as setting $\xi, \eta = 0$. ✓

$$\therefore \int_{-1}^1 \int_{-1}^1 [u_{ij,i}^h(\xi, \eta) - u_{ij,i}^h(0,0)] j(\xi, \eta) d\xi d\eta = 0, \text{ all other terms are}$$

$\delta y_a, y_a, x_b, x_a, y_b, \delta x_a$ are constants and do not change overall result

$$1 \pm \eta \quad \eta \pm \eta^2$$



Final Exam Solutions

60)

- 1 a. Given $f: \Omega \times]0, T[\rightarrow \mathbb{R}$, $g: \Gamma_g \times]0, T[\rightarrow \mathbb{R}$,
 $h: \Gamma_h \times]0, T[\rightarrow \mathbb{R}$, $w_0: \Omega \rightarrow \mathbb{R}$, and
 $\dot{w}_0: \Omega \rightarrow \mathbb{R}$

as usual

Find $w: \bar{\Omega} \times [0, T] \rightarrow \mathbb{R}$ ($w \in \mathcal{S}$)
 $\ni \forall \bar{w}: \bar{\Omega} \times [0, T] \rightarrow \mathbb{R}$ ($\bar{w} \in \mathcal{V}$) $\ni$

$$(\bar{w}, \ddot{w}) + a(\bar{w}, w) = (\bar{w}, f) + (\bar{w}, h)_{\Gamma}$$

$$(\bar{w}, w(0) - w_0) = 0$$

$$(\bar{w}, \dot{w}(0) - \dot{w}_0) = 0$$

where: $(\bar{w}, \ddot{w}) = \int_{\Omega} \bar{w} \ddot{w} d\Omega$

$$a(\bar{w}, w) = \int_{\Omega} \{ \bar{w}_{,i} w_{,i} + \alpha \bar{w} w \} d\Omega$$

$$(\bar{w}, f) = \int_{\Omega} \bar{w} f d\Omega$$

$$(\bar{w}, h)_{\Gamma} = \int_{\Gamma_h} \bar{w} h d\Gamma, \text{ etc.}$$

- b. Given ... ; Find $w^h = w^h + g^h \in \mathcal{S}^h$
 $\ni \forall \bar{w}^h \in \mathcal{V}^h$ (as usual)

$$(\bar{w}^h, \ddot{w}^h) + a(\bar{w}^h, w^h) = (\bar{w}^h, f) + (\bar{w}^h, h)_{\Gamma}$$

1b (cont'd)

$$(\bar{u}^h, \ddot{u}^h) + \alpha (\bar{u}^h, u^h) = (\bar{u}^h, f) + (\bar{u}^h, h)_\Gamma - (\bar{u}^h, \ddot{g}^h) - \alpha (\bar{u}^h, g^h)$$

(Likewise for initial conditions)

1c
$$\underline{M} \underline{\ddot{d}} + \underline{K} \underline{d} = \underline{F} \quad ; \quad \underline{d}(0) = \underline{d}_0 \quad ; \quad \underline{\dot{d}}(0) = \underline{\dot{d}}_0$$

where

$$m_{ab}^e = \int_{\Omega^e} N_a N_b d\Omega$$

$$k_{ab}^e = \int_{\Omega^e} N_{a,i} N_{b,i} + \alpha N_a N_b d\Omega$$

$$f_a^e = \int_{\Omega^e} N_a f d\Omega + \int_{\Gamma_{ch}^e} N_a h d\Gamma - \sum_{b=1}^{n_{ch}} (m_{ab}^e \ddot{g}_b + k_{ab}^e g_b)$$

(usual assembly algo. ; + initial conditions)

2

The model says:

$$y_{n+1} + \lambda \tilde{d}_{n+1} = F_{n+1} \quad (1)$$

$$\tilde{d}_{n+1} = d_n + \Delta t (1-\alpha) y_n \quad (2)$$

$$d_{n+1} = \tilde{d}_{n+1} + \Delta t \alpha y_{n+1} \quad (3)$$

$$(1) \text{ and } (3) \Rightarrow (1-\alpha\Delta t\lambda)y_{n+1} + \lambda d_{n+1} = F_{n+1} \quad (4)$$

$$\text{likewise at } n: (1-\alpha\Delta t\lambda)y_n + \lambda d_n = F_n \quad (5)$$

$$(2) \text{ and } (3) \Rightarrow d_{n+1} = d_n + \Delta t (1-\alpha) y_n + \Delta t \alpha y_{n+1} \quad (6)$$

mult. (6) by $(1-\alpha\Delta t\lambda)$:

$$(1-\alpha\Delta t\lambda)d_{n+1} = (1-\alpha\Delta t\lambda)d_n + \Delta t (1-\alpha) \overbrace{(1-\alpha\Delta t\lambda)y_n}^{(5)} + \Delta t \alpha \underbrace{(1-\alpha\Delta t\lambda)y_{n+1}}_{(4)}$$

$$\cancel{(1-\alpha\Delta t\lambda)}d_{n+1} = \cancel{(1-\alpha\Delta t\lambda)}d_n + \Delta t (1-\alpha) [F_n - \lambda d_n] + \Delta t \alpha [F_{n+1} - \cancel{\lambda d_{n+1}}]$$

$$d_{n+1} = (1-\lambda\Delta t\epsilon)d_n + \Delta t \underbrace{[(1-\alpha)F_n + \alpha F_{n+1}]}_{F_{n+\alpha}}$$

$$(30) \quad a. A = 1 - \lambda\Delta t\epsilon$$

$$(15) \quad b. |A| < 1 \Rightarrow \lambda\Delta t\epsilon < 2.$$

$$(15) \quad c. T(\epsilon_n) = d(\epsilon_{n+1}) - (1-\lambda\Delta t\epsilon) d(\epsilon_n) - \Delta t F_{n+\alpha} \\ = \Delta t \left(\frac{d(\epsilon_n)}{1} + \lambda d(\epsilon_n) - F(\epsilon_n) \right) \\ + \Delta t^2 \left(\frac{\ddot{d}(\epsilon_n)}{2} - \alpha \frac{F(\epsilon_n)}{1} \right) + \dots \\ \therefore |T(\epsilon)| \leq C \Delta t^2$$

$$(10) \quad d. k=1 \text{ (first-order)}$$

$$\stackrel{(25)}{\Rightarrow} \begin{Bmatrix} d_{n+1} \\ v_{n+1} \end{Bmatrix} = \hat{A} \begin{Bmatrix} d_n \\ v_n \end{Bmatrix}$$

$$\hat{A} = \begin{bmatrix} 1 & \Delta t \\ -\Delta t \omega^2 & 1 \end{bmatrix}$$

$$\rho(\hat{A}) = (1 + \Omega^2)^{1/2} > 1 \quad \forall \Omega > 0 \quad ; \quad \Omega = \omega \Delta t$$

\$\therefore\$ unconditionally unstable

$$\stackrel{(50)}{b.} \quad \rho(\hat{A}) = 1 / (1 + \Omega^2)^{1/2} < 1, \quad \forall \Omega > 0$$

\$\therefore\$ unconditionally stable

first-order accurate [\$T = O(\Delta t^2)\$]

$$\stackrel{(35)}{4.} \quad m_{pq} = \delta_{ij} \int_{-1}^{+1} \int_{-1}^{+1} g N_a N_b d\Omega; \quad p = (a-1)n_{ad} + i \\ q = (b-1)n_{bd} + j \\ = \delta_{ij} \int_{-1}^{+1} \int_{-1}^{+1} g N_a N_b d\tilde{\xi} d\tilde{\eta} = \delta_{ij} \frac{\Delta^e}{4} \int_{-1}^{+1} \int_{-1}^{+1} g N_a N_b d\tilde{\xi} d\tilde{\eta} \\ \stackrel{\text{trap rule}}{=} \delta_{ij} \frac{\Delta^e}{4} (\delta_{a1} \delta_{b1} g(\tilde{\xi}_1, \tilde{\eta}_1) + \dots + \delta_{a4} \delta_{b4} g(\tilde{\xi}_4, \tilde{\eta}_4)) \\ = \delta_{ij} \frac{\Delta^e}{4} (\delta_{a1} \delta_{b1} g(-1, -1) + \dots + \delta_{a4} \delta_{b4} g(-1, 1))$$

$$[m_{pq}] = \frac{\Delta^e}{4} \begin{bmatrix} g(-1, -1) & & & & \\ & g(-1, -1) & & & \\ & & g(1, -1) & & \\ & & & g(1, -1) & \\ & & & & g(1, 1) \\ \text{diagonal} \rightarrow & & & & & g(1, 1) \\ \text{zeros elsewhere} & & & & & & g(1, 1) \\ & & & & & & & g(-1, 1) \end{bmatrix}$$

$$5. \quad \mu_i^h = \sum_{a=1}^4 N_a \, da$$

$$\begin{aligned} & \int_{-1}^{+1} \int_{-1}^{+1} (\mu_{i,i}^h(\xi, \zeta) - \mu_{i,i}^h(0,0)) j(\xi, \zeta) \, d\xi \, d\zeta \\ &= \sum_{a=1}^4 da \left\{ \int_{-1}^{+1} \int_{-1}^{+1} (N_{a,i} - N_{a,i}|_{(0,0)}) j \, d\xi \, d\zeta \right\} \end{aligned}$$

recall:

$$\begin{bmatrix} \xi_x & \xi_y \\ \zeta_x & \zeta_y \end{bmatrix} = \begin{bmatrix} x_\xi & x_\zeta \\ y_\xi & y_\zeta \end{bmatrix}^{-1} = j^{-1} \begin{bmatrix} y_\zeta & -x_\zeta \\ -y_\xi & x_\xi \end{bmatrix}$$

consider:

$$\begin{aligned} N_{a,x} &= (N_{a,\xi} N_{b,\zeta} - N_{a,\zeta} N_{b,\xi}) y_b / j \\ &\quad \text{(sum } b) \\ &= [\xi_a (1 + \zeta_a \zeta) \zeta_b (1 + \xi_b \xi) \\ &\quad - \zeta_a (1 + \xi_a \xi) \xi_b (1 + \zeta_b \zeta)] y_b / (16j) \\ &\quad \text{(sum } b) \\ &= (\xi_a \zeta_b - \zeta_a \xi_b + \text{dim \& bilin* terms} \\ &\quad \text{in } \xi, \zeta) y_b / (16j) \end{aligned}$$

$$\therefore \int_{-1}^{+1} \int_{-1}^{+1} N_{a,x} j \, d\xi \, d\zeta = (\xi_a \zeta_b - \zeta_a \xi_b) y_b / 4 \quad \text{(sum } b)$$

* The bilin. terms actually cancel, but this is not needed in the proof.

5. (cont'd) The preceding result derives from the fact that the j 's cancel and lin & bilin terms integrate to zero.

$$N_{a,x} \Big|_{(0,0)} = \frac{(\xi_a \zeta_b - \zeta_a \xi_b) y_b}{\text{sum } b}$$

$$\sum_{-1}^{+1} \sum_{-1}^{+1} N_{a,x} \Big|_{(0,0)} j d\xi d\zeta = N_{a,x} \Big|_{(0,0)} \sum_{-1}^{+1} \sum_{-1}^{+1} j d\xi d\zeta$$

Thus to prove $0 = \sum_{-1}^{+1} \sum_{-1}^{+1} (N_{a,x} - N_{a,x} \Big|_{(0,0)}) j d\xi d\zeta$
it suffices to show that

$$\sum_{-1}^{+1} \sum_{-1}^{+1} j d\xi d\zeta = 4 j(0,0)$$

$$\begin{aligned} j &= x_{,\xi} y_{,\zeta} - x_{,\zeta} y_{,\xi} \\ &= (N_{a,\xi} N_{b,\zeta} - N_{a,\zeta} N_{b,\xi}) x_a y_b \quad (\text{sum } a,b) \\ &= (\xi_a \zeta_b - \zeta_a \xi_b + \text{lin \& bilin}^* \text{ terms in } \xi, \zeta) x_a y_b \quad (\text{sum } a,b) \end{aligned}$$

$$\therefore \sum_{-1}^{+1} \sum_{-1}^{+1} j d\xi d\zeta = 4 (\xi_a \zeta_b - \zeta_a \xi_b) x_a y_b$$

$$\text{and } j(0,0) = (\xi_a \zeta_b - \zeta_a \xi_b) x_a y_b. \quad \blacksquare$$

Is there a simpler way to see this?

* See footnote on p. 5.

TECHNIQUES FOR DEVELOPING 'SPECIAL' FINITE ELEMENT SHAPE FUNCTIONS WITH PARTICULAR REFERENCE TO SINGULARITIES

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SUMMARY

A concise and efficient algorithm is presented for deriving finite element shape functions from an arbitrary set of independent functions. Based on special one-dimensional interpolatory schemes derived via the algorithm, a variety of two- and three-dimensional interpolatory schemes are developed which are useful for modelling singular behaviour. The methodology presented is general and may be fruitfully applied to the development of 'special' finite element shape functions for a variety of other situations.

INTRODUCTION

In many areas of finite element analysis, elements with special properties are required to achieve maximal accuracy. As examples, we may mention 'upwind' finite elements for treating convection operators;¹⁻³ infinite elements for the representation of spatial domains which extend to infinity;^{4,5} boundary-layer elements;⁶ and singular elements for modelling point and line singularities engendered by geometric features such as re-entrant corners and cracks.⁷⁻⁹ In this paper we are concerned with techniques for developing shape (i.e. interpolation) functions for special elements in general, but with particular emphasis on singularities.

Geometric features of the spatial domain may give rise to highly singular response in the solution of a field problem.¹⁰ Eigenfunction analyses are frequently available for purposes of characterizing the asymptotic strength of the singularity.¹¹ However, the actual solution depends upon the nature of the loading and there are situations when the singularity does not contribute significantly.¹² For this reason it has been suggested that, ideally, singular finite elements should strike some balance between being able to represent smooth behaviour and potential singular behaviour.¹² In this way the element will remain effective under all loading conditions. Translated into analytical terms, the finite element shape functions should be designed to represent the standard low-order polynomials as well as the singular functions emanating from asymptotic analysis. To construct shape functions of this type gives rise to an interpolation problem which involves a system of linear simultaneous equations. In special circumstances explicit representations are available for the solution of the interpolation problem. A most noteworthy example is the Lagrange interpolation formula (see, e.g., Davis¹³)

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which is the solution of the interpolation problem for polynomials passing through distinct points. It appears, unfortunately, that no such representation exists when different classes of functions are merged, as seems appropriate for singular elements. Although, in principle, the interpolation problem could be solved numerically for each singular element during formation, this would be highly inefficient due to the increased number of computations required. Furthermore, numerical sensitivity may manifest itself in certain configurations. Consequently, it appears desirable to develop techniques for systematically defining shape functions for singularity modelling (and for developing special elements in general) which circumvent the interpolation problem. This is further attested to by the fact that very few elements developed so far achieve the desired properties delineated above (examples of ones that do may be found in References 7 and 14), and those that do have been developed by *ad hoc* means; thus generalization to higher order elements, or more complex situations, is by no means obvious.

In the next section we present a highly concise algorithm for generating shape functions from an arbitrary starting set of independent functions. A novel feature of the procedure is that *no* coefficient matrix is involved (i.e. no system of simultaneous equations need be solved).

Then we develop some useful one-dimensional interpolatory schemes for modelling singularities by employing the interpolation algorithm. These schemes are the basis for constructing two- and three-dimensional elements for modelling point and line singularities. A number of examples are presented which illustrate the methodology of constructing multi-dimensional elements. Several of these examples appear to be new and are of interest in their own right.

Conclusions are presented in the final section.

PRELIMINARY RESULTS

Let $\{\xi_a | a = 1, 2, \dots, n\}$ denote a set of distinct points (i.e. 'nodes') in the interval $[-1, 1]$. A generic point in $[-1, 1]$ is denoted ξ . Let $r = r(\xi) = (1 + \xi)/2$ and $r_a = r(\xi_a)$. Note that $r(\xi) \in [0, 1]$ for all $\xi \in [-1, 1]$.† The physical co-ordinate is denoted x .

In multi-dimensional applications, we use bold-faced notations analogous to the preceding ones. For example, in three dimensions $\xi = (\xi, \eta, \zeta)$; $\mathbf{r} = (r, s, t)$; where $s = (1 + \eta)/2$; $t = (1 + \zeta)/2$; $\eta, \zeta \in [-1, 1]$, and $\mathbf{x} = (x, y, z)$.

Lagrange interpolation formula

The *Lagrange polynomials* are defined by

$$l_a(f) = \frac{\prod_{\substack{b=1 \\ b \neq a}}^m (f - f_b)}{\prod_{\substack{b=1 \\ b \neq a}}^m (f_a - f_b)}, \quad a = 1, 2, \dots, m \quad (1)$$

where the f_a 's are distinct points in $\mathbb{R}$. It may be seen from (1) that $l_a(f)$ is an $m-1$ st order polynomial in f which satisfies the 'interpolation property' (i.e. $l_a(f_b) = \delta_{ab}$, the Kronecker delta). Note that f may be taken to be ξ, r, r^α , or any other convenient function of ξ .

† Both ' r ' and ' ξ ' co-ordinate descriptions are commonly used in the formulation of finite-element properties. The former seems preferable in modelling singularities, whereas the latter is often favoured because quadrature rules are generally tabulated with respect to the interval $[-1, 1]$. In this paper we find it convenient to employ the r -description. Results may be translated to the ξ -description via change of variables.

Algorithm for constructing interpolation functions

Consider an n -node finite element. We assume we are given a set of 'preliminary' shape functions, N_a , $a = 1, 2, \dots, n$, which satisfy the interpolation property on the first m nodes only; viz. $N_a(\mathbf{r}_b) = \delta_{ab}$, where $a, b = 1, 2, \dots, m < n$. The idea is to systematically modify the N_a 's such that the interpolation property is satisfied on all n nodes. The procedure is given as follows:†

$$\text{Step 1} \quad N_{m+1}(\mathbf{r}) \leftarrow \frac{N_{m+1}(\mathbf{r}) - \sum_{a=1}^m N_{m+1}(\mathbf{r}_a) N_a(\mathbf{r})}{N_{m+1}(\mathbf{r}_{m+1}) - \sum_{a=1}^m N_{m+1}(\mathbf{r}_a) N_a(\mathbf{r}_{m+1})} \quad (2)$$

$$\text{Step 2} \quad N_a(\mathbf{r}) \leftarrow N_a(\mathbf{r}) - N_a(\mathbf{r}_{m+1}) N_{m+1}(\mathbf{r}), \quad a = 1, 2, \dots, m \quad (3)$$

Step 3 If $m+1 < n$, replace m by $m+1$ and repeat steps 1–3.
 If $m+1 = n$, stop.

Remarks

I. After completing steps 1 and 2, the shape functions satisfy the interpolation property on the first $m+1$ nodes. Clearly, after completing the entire procedure, the desired properties are achieved.

II. In many cases, the algorithm enables the explicit hand calculation of shape functions. *However, there is no need to do this in practice, as the algorithm itself may be programmed as part of a shape function subroutine.‡*

III. The derivatives of shape functions may be constructed in similar fashion. The analogues of steps 1 and 2 are (respectively)

$$\partial N_{m+1}(\mathbf{r}) \leftarrow \frac{\partial N_{m+1}(\mathbf{r}) - \sum_{a=1}^m N_{m+1}(\mathbf{r}_a) \partial N_a(\mathbf{r})}{N_{m+1}(\mathbf{r}_{m+1}) - \sum_{a=1}^m N_{m+1}(\mathbf{r}_a) N_a(\mathbf{r}_{m+1})} \quad (4)$$

and

$$\partial N_a(\mathbf{r}) \leftarrow \partial N_a(\mathbf{r}) - N_a(\mathbf{r}_{m+1}) \partial N_{m+1}(\mathbf{r}) \quad (5)$$

In (4) and (5), ' ∂ ' represents the partial differentiation operator with respect to any of the co-ordinates (e.g. in three dimensions, ∂ may represent $\partial/\partial r$, $\partial/\partial s$ or $\partial/\partial t$).

IV. Interpolation problems are generally formulated in terms of a system of simultaneous linear equations. It may be noted that, if $m = 1$, the algorithm encompassed by steps 1–3 solves the interpolation problem *without* engendering any coefficient matrix. Thus no 'storage' is required (beyond that for the N_a 's) in programming the procedure. Another nice feature of the algorithm is that the interpolation property possessed by the preliminary shape functions is fully exploited in that only $n - m$ passes through steps 1–3 need be performed.

In practice, the Lagrange polynomial formula may be employed to generate a partial set of preliminary shape functions which satisfies the interpolation property on the first m nodes. This will be illustrated in the following examples.

EXAMPLES

1. Consider a one-dimensional three-node element with nodes given by $r_1 = 0$, $r_2 = \frac{1}{2}$ and $r_3 = 1$. We wish to construct shape functions capable of exactly representing 1, r and r^α . (If $\alpha = 2$, we

† The arrow, $\leftarrow$, reads: 'is replaced by'.

‡ By a 'shape function subroutine', we mean a subprogram which, upon given a value of $\mathbf{r}$, returns the N_a 's and/or derivatives of the N_a 's, evaluated at $\mathbf{r}$.

have the usual three-node quadratic element. Other values of α enable the modelling of singularities.) We begin with the following preliminary shape functions:

$$N_1(r) = 1 - 2r; \quad N_2(r) = 2r; \quad N_3(r) = r^\alpha \quad (6)$$

Observe that $N_a(r_b) = \delta_{ab}$, $1 \leq a, b \leq 2$. In terms of our algorithm, $m = 2$ and $n = 3$, so only one pass through steps 1 and 2 is required:

$$N_3(r) \leftarrow \frac{N_3(r) - N_3(\frac{1}{2})N_2(r)}{N_3(1) - N_3(\frac{1}{2})N_2(1)} = \frac{r^\alpha - 2(\frac{1}{2})^\alpha r}{1 - 2(\frac{1}{2})^\alpha} \quad (7)$$

$$N_1(r) \leftarrow N_1(r) - N_1(1)N_3(r) = 1 - 2r + \left[\frac{r^\alpha - 2(\frac{1}{2})^\alpha r}{1 - 2(\frac{1}{2})^\alpha} \right] \quad (8)$$

$$N_2(r) \leftarrow N_2(r) - N_2(1)N_3(r) = 2r - 2 \left[\frac{r^\alpha - 2(\frac{1}{2})^\alpha r}{1 - 2(\frac{1}{2})^\alpha} \right] \quad (9)$$

Likewise, the derivatives are given by (in this case ∂N is taken to mean $\partial N/\partial r$):

$$\partial N_3(r) \leftarrow \frac{\partial N_3(r) - N_3(\frac{1}{2}) \partial N_2(r)}{N_3(1) - N_3(\frac{1}{2})N_2(1)} = \frac{\alpha r^{\alpha-1} - 2(\frac{1}{2})^\alpha}{1 - 2(\frac{1}{2})^\alpha} \quad (10)$$

$$\partial N_1(r) \leftarrow \partial N_1(r) - N_1(1) \partial N_3(r) = -2 + \left[\frac{\alpha r^{\alpha-1} - 2(\frac{1}{2})^\alpha}{1 - 2(\frac{1}{2})^\alpha} \right] \quad (11)$$

$$\partial N_2(r) \leftarrow \partial N_2(r) - N_2(1) \partial N_3(r) = 2 - 2 \left[\frac{\alpha r^{\alpha-1} - 2(\frac{1}{2})^\alpha}{1 - 2(\frac{1}{2})^\alpha} \right] \quad (12)$$

2. Consider a one-dimensional four-node element for which the nodes are given by $r_1 = 0$, $r_2 = \frac{1}{3}$, $r_3 = \frac{2}{3}$ and $r_4 = 1$. We wish to construct shape functions capable of exactly representing 1 , r^α , $r^{2\alpha}$ and r^β . The following special cases are of practical importance:

$\alpha = 1, \beta = 3$: $1, r, r^2, r^3$ cubic polynomial in r

$\alpha = 1, \beta$ arbitrary: $1, r, r^2, r^\beta$ quadratic polynomial in r plus linear polynomial in r^β

α arbitrary, $\beta = 1$: $1, r, r^\alpha, r^{2\alpha}$ linear polynomial in r plus quadratic polynomial in r^α

α arbitrary, $\beta = 3\alpha$: $1, r^\alpha, r^{2\alpha}, r^{3\alpha}$ cubic polynomial in r^α

To obtain a starting set of shape functions, we may employ the Lagrange interpolation formula with $f = r^\alpha$ and $m = 3$, namely

$$N_1(r) = l_1(r^\alpha) = \frac{(r^\alpha - r_2^\alpha)(r^\alpha - r_3^\alpha)}{(r_1^\alpha - r_2^\alpha)(r_1^\alpha - r_3^\alpha)} = \frac{(3^\alpha r^\alpha - 1)(3^\alpha r^\alpha - 2^\alpha)}{2^\alpha} \quad (13)$$

$$N_2(r) = l_2(r^\alpha) = \frac{(r^\alpha - r_1^\alpha)(r^\alpha - r_3^\alpha)}{(r_2^\alpha - r_1^\alpha)(r_2^\alpha - r_3^\alpha)} = \frac{3^\alpha r^\alpha (3^\alpha r^\alpha - 2^\alpha)}{(1 - 2^\alpha)} \quad (14)$$

$$N_3(r) = l_3(r^\alpha) = \frac{(r^\alpha - r_1^\alpha)(r^\alpha - r_2^\alpha)}{(r_3^\alpha - r_1^\alpha)(r_3^\alpha - r_2^\alpha)} = \frac{3^\alpha r^\alpha (3^\alpha r^\alpha - 1)}{2^\alpha (2^\alpha - 1)} \quad (15)$$

It may be verified that $N_a(r_b) = \delta_{ab}$ for all $a, b = 1, 2, 3$. The last preliminary shape function is taken to be

$$N_4(r) = r^\beta \quad (16)$$

With regard to the algorithm, $m = 3$ and $n = 4$, so only one pass through steps 1 and 2 is needed, namely

$$N_4(r) \leftarrow \frac{N_4(r) - \sum_{a=2}^3 N_4(r_a) N_a(r)}{1 - \sum_{a=2}^3 N_4(r_a) N_a(1)} \quad (17)$$

$$N_a(r) \leftarrow N_a(r) - N_a(1) N_4(r), \quad a = 1, 2, 3 \quad (18)$$

Equations (17) and (18) may be programmed in a shape function subroutine, along with corresponding expressions for derivatives.

3. As a final one-dimensional example, we consider a five-node element for which the nodes are given by $r_1 = 0, r_2 = \frac{1}{4}, r_3 = \frac{1}{2}, r_4 = \frac{3}{4}$ and $r_5 = 1$. The shape functions are to be capable of exactly representing $1, r^\alpha, r^{2\alpha}, r^\beta$ and r^γ . The following special cases are of practical importance:

$\alpha = 1, \beta = 3, \gamma = 4$:	$1, r, r^2, r^3, r^4$	quartic polynomial in r
$\alpha = 1, \beta = 3, \gamma$ arbitrary:	$1, r, r^2, r^3, r^\gamma$	cubic polynomial in r plus linear polynomial in r^γ .
$\alpha = 1, \beta$ arbitrary, $\gamma = 2\beta$:	$1, r, r^2, r^\beta, r^{2\beta}$	quadratic polynomial in r plus quadratic polynomial in r^β
α arbitrary, $\beta = 3\alpha, \gamma = 1$:	$1, r, r^\alpha, r^{2\alpha}, r^{3\alpha}$	linear polynomial in r plus cubic polynomial in r^α
α arbitrary, $\beta = 3\alpha, \gamma = 4\alpha$:	$1, r^\alpha, r^{2\alpha}, r^{3\alpha}, r^{4\alpha}$	quartic polynomial in r^α

As in the previous example, the most appropriate choice for a set of starting shape functions is given by (13)–(16) and

$$N_5(r) = r^\gamma \quad (19)$$

This time $m = 3$ and $n = 5$, so two passes through steps 1–3 are required. The following relations need be executed, in turn, for $m = 3$ and $m = 4$:

$$N_{m+1}(r) \leftarrow \frac{N_{m+1}(r) - \sum_{a=2}^m N_{m+1}(r_a) N_a(r)}{N_{m+1}(r_{m+1}) - \sum_{a=2}^m N_{m+1}(r_a) N_a(r_{m+1})} \quad (20)$$

$$N_a(r) \leftarrow N_a(r) - N_a(r_{m+1}) N_{m+1}(r), \quad a = 1, 2, \dots, m \quad (21)$$

The above expressions, and their counterparts for derivatives, may be programmed in shape function subroutines.

Remarks

V. The preceding one-dimensional examples cover most cases of practical interest. If the need arises, higher order shape functions may be constructed along similar lines. The basic philosophy is to maximize m so as to minimize calculations in the shape function routines.

VI. For completeness, we wish to mention the simplest one-dimensional shape functions which allow the modelling of singularities; namely the two-node element with $r_1 = 0, r_2 = 1, N_1(r) = 1 - r^\alpha$ and $N_2(r) = r^\alpha$.

VII. Let $u(r) = \sum_{a=1}^n N_a(r) d_a$, where (by construction) $d_a = u(r_a)$. The function u is capable of exactly representing the powers of r which were used in deriving the N_a 's. (The d_a 's need only be set accordingly.) In practice, it is usually desired that the same powers of the physical

co-ordinate, $x = x(r)$, be representable. This will be attained if $x = x(r)$ is affine (i.e. of the form $x(r) = c_1 + c_2 r$). There are two ways of achieving this in practice:

The *first* depends upon 1 and r being present among the powers of r used in deriving the shape functions. In this case the isoparametric concept may be invoked; i.e. we take $x(r) = \sum_{a=1}^n N_a(r)x_a$, and equally space the nodal co-ordinates (i.e. x_a 's in x -space).

The *second* procedure does not require that 1 and r be present, but involves an additional set of shape function calculations, so it is less efficient. In this case, we take $x(r) = \sum_{a=1}^n P_a(r)x_a$, where the P_a 's are the standard polynomial shape functions, and the x_a 's are, again, equally spaced.

To see the necessity of adopting the second procedure when 1 and r are absent from the N_a 's, we need only consider the two-node element of Remark VI. If the isoparametric concept is adopted, then u varies linearly with x (rather than x^a). To see this we note that if $x_2 \neq x_1$ then we can write $d_a = c_1 + c_2 x_a$, where the c 's are constants. Consequently

$$\begin{aligned} u &= \sum_{a=1}^2 N_a d_a \\ &= \sum_{a=1}^2 N_a (c_1 + c_2 x_a) \\ &= c_1 \left(\sum_{a=1}^2 N_a \right) + c_2 \left(\sum_{a=1}^2 N_a x_a \right) \\ &= c_1 + c_2 x \quad \blacksquare \end{aligned} \quad (22)$$

We shall now employ the above results in deriving two-dimensional finite element shape functions for modelling line and point singularities.

Examples

4. Consider the four-node quadrilateral illustrated in Figure 1(a). The shape functions may be constructed from products of linear polynomial shape functions in the s -direction and the

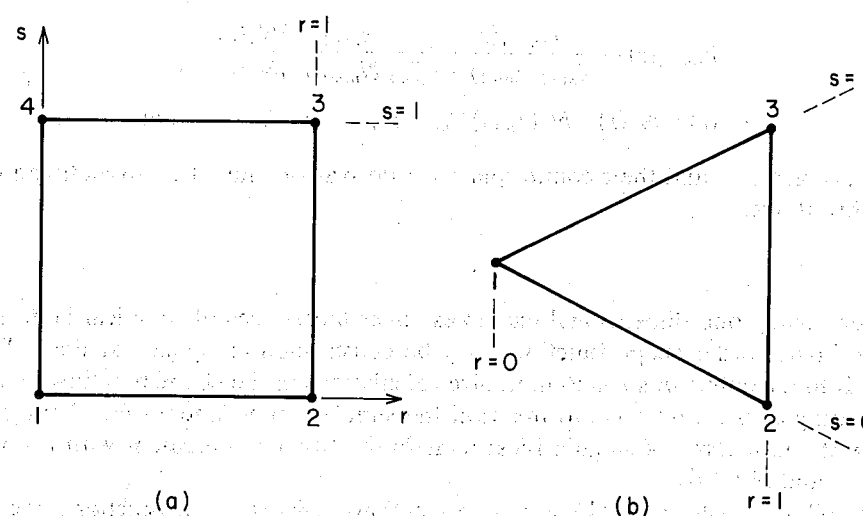


Figure 1. (a) Four-node quadrilateral elements. (b) Three-node triangular element formed by degenerating the four-node quadrilateral

shape functions of Remark VI in the r -direction, namely

$$N_1(r, s) = (1 - r^\alpha)(1 - s) \quad (23)$$

$$N_2(r, s) = r^\alpha(1 - s) \quad (24)$$

$$N_3(r, s) = r^\alpha s \quad (25)$$

$$N_4(r, s) = (1 - r^\alpha)s \quad (26)$$

Thus $u(r, s) = \sum_{a=1}^4 N_a(r, s)d_a$ is capable of exactly representing a line singularity of order r^α along edge 14. By virtue of Remark VII the standard polynomial shape functions need be employed to define the geometry. That is $\mathbf{x}(r, s) = \sum_{a=1}^4 P_a(r, s)\mathbf{x}_a$, where

$$P_1(r, s) = (1 - r)(1 - s) \quad (27)$$

$$P_2(r, s) = r(1 - s) \quad (28)$$

$$P_3(r, s) = rs \quad (29)$$

$$P_4(r, s) = (1 - r)s \quad (30)$$

This element was first proposed by Akin.¹⁵

5. A three-node triangular element, see Figure 1(b), may be developed by combining the shape functions associated with nodes 3 and 4 in the previous example, namely

$$N_1(r, s) \leftarrow N_1(r, s) + N_4(r, s) = 1 - r^\alpha \quad (31)$$

$$P_1(r, s) \leftarrow P_1(r, s) + P_4(r, s) = 1 - r \quad (32)$$

The shape functions of nodes 2 and 3 remain the same. This element is capable of exactly representing a point singularity of order r^α . Similar interpolations were used in References 9 and 16.

The procedure employed in this example is often referred to as *element degeneration*.

6. Consider the nine-node Lagrange-type quadrilateral illustrated in Figure 2(a). The shape functions shall be constructed from products of the three-node one-dimensional shape functions

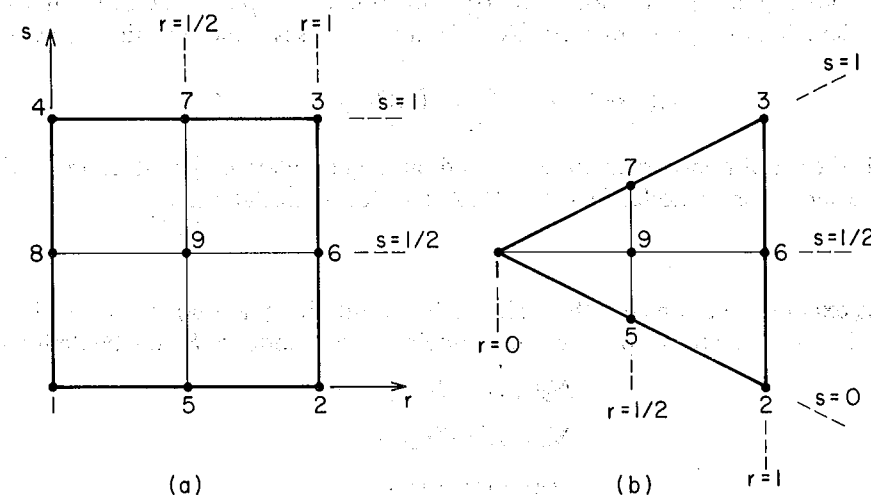


Figure 2. (a) Nine-node Lagrange quadrilateral element. Node 9 is omitted for the eight-node serendipity quadrilateral. (b) Seven-node triangular element formed by degenerating the nine-node Lagrange quadrilateral. In the six-node triangle, node 9 is omitted

of Example 1. That is, we define $N_a(r)$, $a = 1, 2, 3$, by (7)–(9); and $P_a(s)$ by the same expressions with r replaced by s and $\alpha = 2$. The two-dimensional shape functions may then be defined as follows:

$$N_1(r, s) = N_1(r)P_1(s) \quad (33)$$

$$N_2(r, s) = N_3(r)P_1(s) \quad (34)$$

$$N_3(r, s) = N_3(r)P_3(s) \quad (35)$$

$$N_4(r, s) = N_1(r)P_3(s) \quad (36)$$

$$N_5(r, s) = N_2(r)P_1(s) \quad (37)$$

$$N_6(r, s) = N_3(r)P_2(s) \quad (38)$$

$$N_7(r, s) = N_2(r)P_3(s) \quad (39)$$

$$N_8(r, s) = N_1(r)P_2(s) \quad (40)$$

$$N_9(r, s) = N_2(r)P_2(s) \quad (41)$$

In this case $u(r, s) = \sum_{a=1}^9 N_a(r, s)d_a$ is capable of exactly representing a line singularity of order r^α along edge 14. Furthermore, the following monomials may be exactly represented: $1, r, s, r^\alpha, rs, s^2, r^\alpha s, s^2 r$ and $r^\alpha s^2$. By virtue of the presence of $1, r$ and s , either the isoparametric concept may be employed to define the geometry (i.e. $\mathbf{x}(r, s) = \sum_{a=1}^9 N_a(r, s)\mathbf{x}_a$), or recourse may be made to the standard Lagrange polynomials (i.e. $\mathbf{x}(r, s) = \sum_{a=1}^9 P_a(r, s)\mathbf{x}_a$, where the $P_a(r, s)$'s may be obtained from (33)–(41) by replacing $N_a(r)$ by $P_a(r)$, $a = 1, 2, 3$).

Remark

VIII. Serendipity-type nodal patterns for quadrilaterals are the same as Lagrange nodal patterns, except that all internal nodes are omitted. Our procedure for developing shape functions for serendipity-type elements is to first establish shape functions for the non-corner nodes which satisfy the interpolation property (it turns out that this may easily be done), then establish a preliminary set of shape functions for the four corner nodes which are systematically modified to satisfy the interpolation property. For example, suppose the preliminary corner-node shape functions are given by (23)–(26). Then the necessary modifications may be written

$$N_a(\mathbf{r}) \leftarrow N_a(\mathbf{r}) - \sum_{b \in B_a} N_a(\mathbf{r}_b)N_b(\mathbf{r}), \quad a \in C \quad (42)$$

where B_a is the set of non-corner nodes located on the boundary edges which meet at node a , and C is the set of corner nodes (e.g. $C = \{1, 2, 3, 4\}$ for quadrilaterals).

Examples

7. In this example, we consider the eight-node serendipity quadrilateral in which node 9 of Figure 2(a) is omitted. The shape functions associated with nodes 5–8 may be defined by

$$N_5(r, s) = N_2(r)(1 - s) \quad (43)$$

$$N_6(r, s) = r^\alpha P_2(s) \quad (44)$$

$$N_7(r, s) = N_2(r)s \quad (45)$$

$$N_8(r, s) = (1 - r^\alpha)P_2(s) \quad (46)$$

In (43)–(46), $N_2(r)$ and $P_2(s)$ are defined as in the previous example. The shape functions for nodes 1–4 may be defined by (42), where the preliminary corner-node shape functions are given by (23)–(26); explicitly:

$$N_1(r, s) \leftarrow N_1(r, s) - \frac{1}{2}N_8(r, s) - \left(1 - \frac{1}{2^\alpha}\right)N_5(r, s) \quad (47)$$

$$N_2(r, s) \leftarrow N_2(r, s) - \frac{1}{2^\alpha}N_5(r, s) - \frac{1}{2}N_6(r, s) \quad (48)$$

$$N_3(r, s) \leftarrow N_3(r, s) - \frac{1}{2}N_6(r, s) - \frac{1}{2^\alpha}N_7(r, s) \quad (49)$$

$$N_4(r, s) \leftarrow N_4(r, s) - \left(1 - \frac{1}{2^\alpha}\right)N_7(r, s) - \frac{1}{2}N_8(r, s) \quad (50)$$

The serendipity shape functions defined above are capable of exactly representing the same monomials as the nine-node Lagrange shape functions of Example 6, except for rs^2 , which is no longer present.

An alternative definition of serendipity shape functions may be constructed by replacing r^α by r in (23)–(26), (44) and (46). In this case the preliminary corner-node shape functions reduce to the standard bilinear shape functions and (42) yields in place of (47)–(50):

$$N_1(r, s) \leftarrow N_1(r, s) - [N_8(r, s) + N_5(r, s)]/2 \quad (51)$$

$$N_2(r, s) \leftarrow N_2(r, s) - [N_5(r, s) + N_6(r, s)]/2 \quad (52)$$

$$N_3(r, s) \leftarrow N_3(r, s) - [N_6(r, s) + N_7(r, s)]/2 \quad (53)$$

$$N_4(r, s) \leftarrow N_4(r, s) - [N_7(r, s) + N_8(r, s)]/2 \quad (54)$$

This alternative definition of serendipity shape functions is capable of exactly representing the same monomials as the nine-node Lagrange shape functions, except for $r^\alpha s^2$.

This example is typical in that the definition of serendipity shape functions from a given set of one-dimensional shape functions is non-unique. The choice must be made on the basis of the appropriateness of the functions for the intended application.

Again, either the $N_a(r, s)$'s, or the standard polynomial serendipity shape functions, may be used for purposes of defining the geometry. (The polynomial serendipity shape functions may be obtained from (43)–(46) and (51)–(54) by replacing $N_2(r)$ by $P_2(r)$ in (43) and (45) and by replacing r^α by r in (44) and (46).)

8. Seven-node and six-node triangular elements, see e.g. Figure 2(b), for modelling point singularities, may be obtained from the quadrilaterals of Examples 6 and 7, respectively, by combining the shape functions associated with nodes 1, 4 and 8, that is

$$N_1(r, s) \leftarrow N_1(r, s) + N_4(r, s) + N_8(r, s) \quad (55)$$

(If the standard polynomial shape functions are used to define the geometry, they must be similarly combined.) The shape functions associated with the remaining nodes are the same as in Examples 6 and 7. Shape functions N_4 and N_8 play no further role in the calculations, but typically would be set to zero in a shape function subroutine.

9. A 16-node Lagrange quadrilateral, see Figure 3(a), may be formed by taking products of the four-node one-dimensional shape functions of Example 2. In this case, we define $N_a(r)$, $a = 1, 2, 3, 4$, by (13)–(18); and $P_a(s)$ by the same expressions in which $\alpha = 1$ and $\beta = 3$. The results

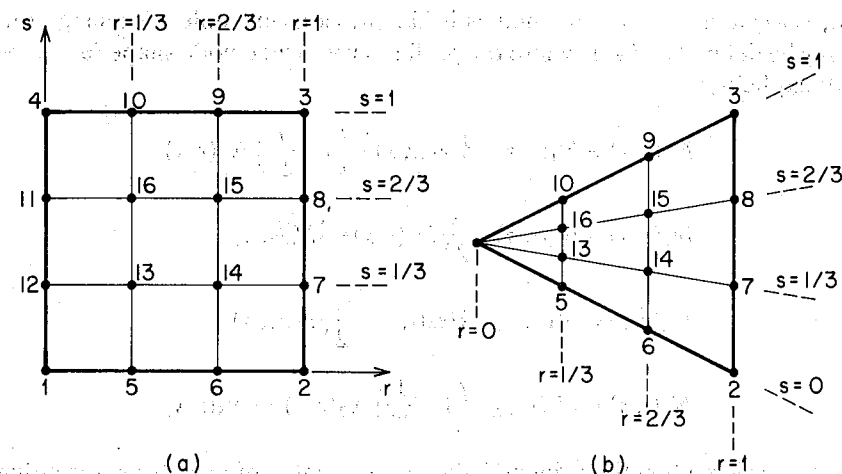


Figure 3: (a) 16-node Lagrange quadrilateral element. Nodes 13–16 are omitted for the 12-node serendipity quadrilateral. (b) 13-node triangular element formed by degenerating the 16-node Lagrange quadrilateral. In the nine-node serendipity triangle, nodes 13–16 are omitted

are

$$N_1(r, s) = N_1(r)P_1(s) \quad (56)$$

$$N_2(r, s) = N_4(r)P_1(s) \quad (57)$$

$$N_3(r, s) = N_4(r)P_4(s) \quad (58)$$

$$N_4(r, s) = N_1(r)P_4(s) \quad (59)$$

$$N_5(r, s) = N_2(r)P_1(s) \quad (60)$$

$$N_6(r, s) = N_3(r)P_1(s) \quad (61)$$

$$N_7(r, s) = N_4(r)P_2(s) \quad (62)$$

$$N_8(r, s) = N_4(r)P_3(s) \quad (63)$$

$$N_9(r, s) = N_3(r)P_4(s) \quad (64)$$

$$N_{10}(r, s) = N_2(r)P_4(s) \quad (65)$$

$$N_{11}(r, s) = N_1(r)P_3(s) \quad (66)$$

$$N_{12}(r, s) = N_1(r)P_2(s) \quad (67)$$

$$N_{13}(r, s) = N_2(r)P_2(s) \quad (68)$$

$$N_{14}(r, s) = N_3(r)P_2(s) \quad (69)$$

$$N_{15}(r, s) = N_3(r)P_3(s) \quad (70)$$

$$N_{16}(r, s) = N_2(r)P_3(s) \quad (71)$$

The monomials present this time are illustrated in Figure 4 by way of a 'Pascal triangle'. As long as either α or β equals 1, the isoparametric concept may be employed to define the geometry. Otherwise, the standard 16-node Lagrange polynomials may be employed. (These may be obtained from (56)–(71) by replacing $N_a(r)$ by $P_a(r)$, $a = 1, 2, 3, 4$.)

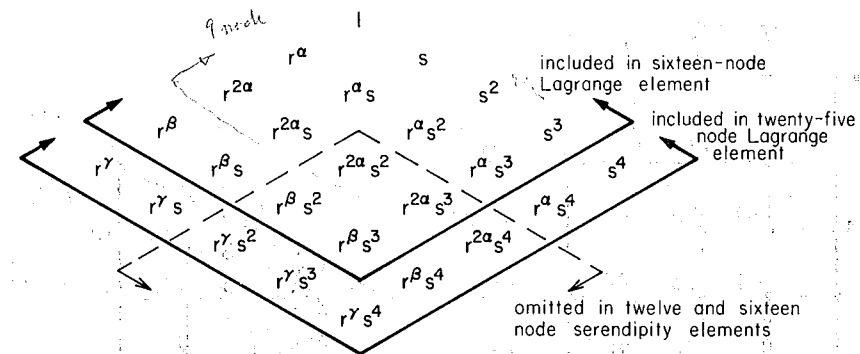


Figure 4. Pascal triangle for Examples 9, 10, 12 and 13

10. A 12-node serendipity quadrilateral in which nodes 13–16 of Figure 3(a) are omitted may be obtained as follows. Define the non-corner node shape functions via

$$N_5(r, s) = N_2(r)(1 - s) \quad (72)$$

$$N_6(r, s) = N_3(r)(1 - s) \quad (73)$$

$$N_7(r, s) = r^\alpha P_2(s) \quad (74)$$

$$N_8(r, s) = r^\alpha P_3(s) \quad (75)$$

$$N_9(r, s) = N_3(r)s \quad (76)$$

$$N_{10}(r, s) = N_2(r)s \quad (77)$$

$$N_{11}(r, s) = (1 - r^\alpha)P_3(s) \quad (78)$$

$$N_{12}(r, s) = (1 - r^\alpha)P_2(s) \quad (79)$$

In (72)–(79), $N_2(r)$, $N_3(r)$, $P_2(s)$ and $P_3(s)$ are defined as in the previous example. The shape functions associated with nodes 1–4 may be obtained from (42). The monomials present may be gleaned from Figure 4.

As in Example 7, the one-dimensional shape functions under consideration may be combined differently to form other serendipity-type elements. For example, we may replace r^α by r^β in (23)–(26), (74), (75), (78) and (79). In this case, the monomials present may be deduced from Figure 4 by interchanging r^α and r^β (the $r^{2\alpha}$ terms remain unchanged).

11. Triangles, see Figure 3(b), may likewise be constructed from the results of Examples 9 and 10 by coalescing the nodes along edge 14. The shape functions need be combined in the usual way:

$$N_1(r, s) \leftarrow N_1(r, s) + N_4(r, s) + N_{11}(r, s) + N_{12}(r, s) \quad (80)$$

and likewise for the polynomial shape functions, if employed to define the geometry. The remaining shape functions are as given in Examples 9 and 10.

Remark

IX. At this point, the algorithm for defining serendipity-type shape functions should be clear from Examples 7 and 10. Likewise, the procedures for defining Lagrange-type quadrilaterals, and for degenerating quadrilaterals into triangles, should also be clear. Consequently, we shall

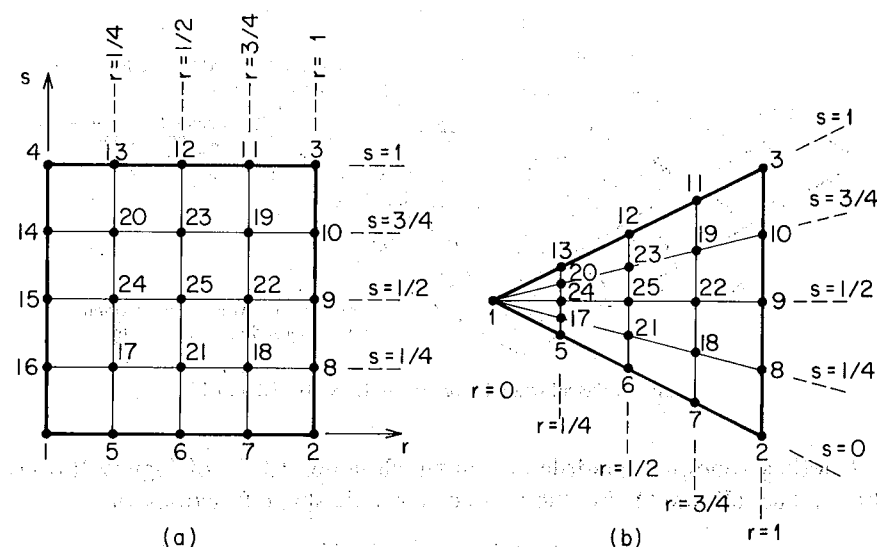


Figure 5. (a) 25-node Lagrange quadrilateral element. Nodes 17–25 are omitted for the 16-node serendipity quadrilateral. (b) 21-node triangular element formed by degenerating the 25-node Lagrange quadrilateral. In the 12-node serendipity triangle, nodes 17–25 are omitted

present without further elaboration the results for the 25-node Lagrange quadrilateral, 16-node serendipity quadrilateral, and corresponding triangular elements (see Figure 5). In these cases, the one-dimensional shape functions of Example 3 are employed.

Examples

12. 25-node Lagrange quadrilateral

$$N_1(r, s) = N_1(r)P_1(s) \quad (81)$$

$$N_2(r, s) = N_5(r)P_1(s) \quad (82)$$

$$N_3(r, s) = N_5(r)P_5(s) \quad (83)$$

$$N_4(r, s) = N_1(r)P_5(s) \quad (84)$$

$$N_5(r, s) = N_2(r)P_1(s) \quad (85)$$

$$N_6(r, s) = N_3(r)P_1(s) \quad (86)$$

$$N_7(r, s) = N_4(r)P_1(s) \quad (87)$$

$$N_8(r, s) = N_5(r)P_2(s) \quad (88)$$

$$N_9(r, s) = N_5(r)P_3(s) \quad (89)$$

$$N_{10}(r, s) = N_5(r)P_4(s) \quad (90)$$

$$N_{11}(r, s) = N_4(r)P_5(s) \quad (91)$$

$$N_{12}(r, s) = N_3(r)P_5(s) \quad (92)$$

$$N_{13}(r, s) = N_2(r)P_5(s) \quad (93)$$

$$N_{14}(r, s) = N_1(r)P_4(s) \quad (94)$$

$$N_{15}(r, s) = N_1(r)P_3(s) \quad (95)$$

$$N_{16}(r, s) = N_1(r)P_2(s) \quad (96)$$

$$N_{17}(r, s) = N_2(r)P_2(s) \quad (97)$$

$$N_{18}(r, s) = N_4(r)P_2(s) \quad (98)$$

$$N_{19}(r, s) = N_4(r)P_4(s) \quad (99)$$

$$N_{20}(r, s) = N_2(r)P_4(s) \quad (100)$$

$$N_{21}(r, s) = N_3(r)P_2(s) \quad (101)$$

$$N_{22}(r, s) = N_4(r)P_3(s) \quad (102)$$

$$N_{23}(r, s) = N_3(r)P_4(s) \quad (103)$$

$$N_{24}(r, s) = N_2(r)P_3(s) \quad (104)$$

$$N_{25}(r, s) = N_3(r)P_3(s) \quad (105)$$

In the above, the $N_a(r)$'s are defined by (13)–(16) and (19)–(21); the $P_a(s)$'s are the standard quartic polynomials, which may be obtained from the $N_a(r)$'s by replacing r with s , and setting $\alpha = 1$, $\beta = 3$ and $\gamma = 4$.

13. 16-node serendipity quadrilateral

$$N_5(r, s) = N_2(r)(1-s) \quad (106)$$

$$N_6(r, s) = N_3(r)(1-s) \quad (107)$$

$$N_7(r, s) = N_4(r)(1-s) \quad (108)$$

$$N_8(r, s) = r^\alpha P_2(s) \quad (109)$$

$$N_9(r, s) = r^\alpha P_3(s) \quad (110)$$

$$N_{10}(r, s) = r^\alpha P_4(s) \quad (111)$$

$$N_{11}(r, s) = N_4(r)s \quad (112)$$

$$N_{12}(r, s) = N_3(r)s \quad (113)$$

$$N_{13}(r, s) = N_2(r)s \quad (114)$$

$$N_{14}(r, s) = (1-r^\alpha)P_4(s) \quad (115)$$

$$N_{15}(r, s) = (1-r^\alpha)P_3(s) \quad (116)$$

$$N_{16}(r, s) = (1-r^\alpha)P_2(s) \quad (117)$$

The $N_a(r)$'s and $P_a(s)$'s in (106)–(117) are the same as used in the previous example. Equation (42) may be employed to define $N_a(r, s)$, $a = 1, 2, 3, 4$. Again we wish to point out that alternative definitions of the serendipity shape functions may also be constructed (cf. Examples 7 and 10).

14. 21- and 12-node triangles

The shape functions of Examples 12 and 13 are combined in the usual way

$$N_1(r, s) \leftarrow N_1(r, s) + N_4(r, s) + N_{14}(r, s) + N_{15}(r, s) + N_{16}(r, s) \quad (118)$$

The remaining functions are defined in Examples 12 and 13. ■

We shall now illustrate, by way of examples, how to use the previously defined two-dimensional shape functions to develop three-dimensional functions for modelling line singularities. We shall assume that the singularity lies along the r -axis.

Examples

15. Consider the wedge-shaped six-node element in Figure 6. Let $N_a(r, s)$, $a = 1, 2, 3$, denote the three-node triangle functions defined in Example 5. Shape functions for the wedge may be

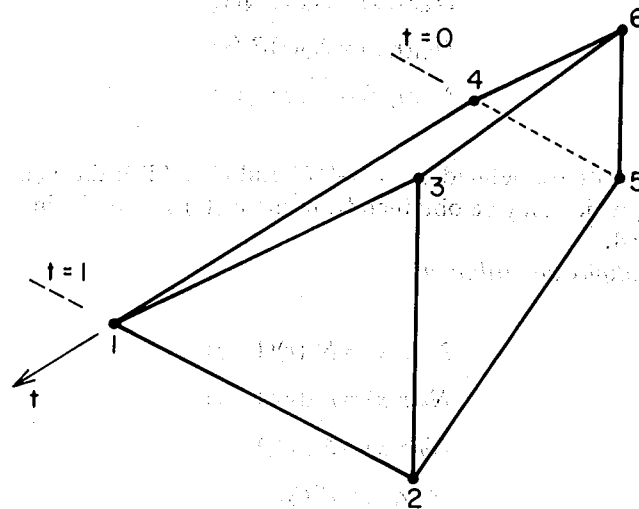


Figure 6. Six-node wedge element

defined as follows:

$$N_1(r, s, t) = N_1(r, s)t \quad (119)$$

$$N_2(r, s, t) = N_2(r, s)t \quad (120)$$

$$N_3(r, s, t) = N_3(r, s)t \quad (121)$$

$$N_4(r, s, t) = N_1(r, s)(1-t) \quad (122)$$

$$N_5(r, s, t) = N_2(r, s)(1-t) \quad (123)$$

$$N_6(r, s, t) = N_3(r, s)(1-t) \quad (124)$$

For the geometric mapping, the $N_a(r, s)$'s in (119)–(124) should be replaced by the standard polynomial shape functions (i.e. $P_a(r, s)$'s in Example 5).

16. Consider the 18-node wedge illustrated in Figure 7(a). Let $\tilde{N}_a(r, s)$, $a = 1, 2, \dots, 6$, denote the shape functions of the six-node triangle which were defined in Example 8, and are

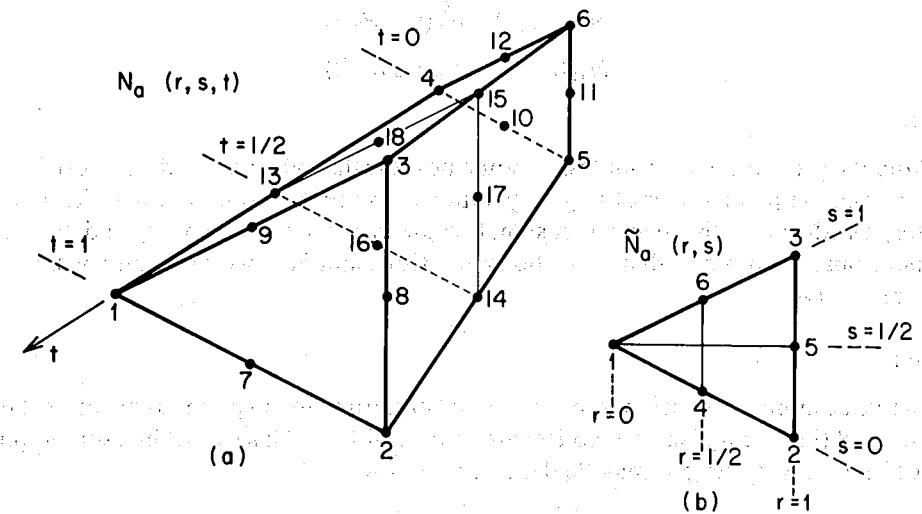


Figure 7. (a) 18-node wedge element. Nodes 16–18 are omitted for the 15-node wedge. (b) Six-node triangular element

renumbered here for convenience—see Figure 7(b). Let $P_a(t)$, $a = 1, 2, 3$, be the standard one-dimensional polynomial shape functions of the three-node element. The three-dimensional shape functions may be defined as follows:

$$N_1(r, s, t) = \tilde{N}_1(r, s)P_3(t) \quad (125)$$

$$N_2(r, s, t) = \tilde{N}_2(r, s)P_3(t) \quad (126)$$

$$N_3(r, s, t) = \tilde{N}_3(r, s)P_3(t) \quad (127)$$

$$N_4(r, s, t) = \tilde{N}_1(r, s)P_1(t) \quad (128)$$

$$N_5(r, s, t) = \tilde{N}_2(r, s)P_1(t) \quad (129)$$

$$N_6(r, s, t) = \tilde{N}_3(r, s)P_1(t) \quad (130)$$

$$N_7(r, s, t) = \tilde{N}_4(r, s)P_3(t) \quad (131)$$

$$N_8(r, s, t) = \tilde{N}_5(r, s)P_3(t) \quad (132)$$

$$N_9(r, s, t) = \tilde{N}_6(r, s)P_3(t) \quad (133)$$

$$N_{10}(r, s, t) = \tilde{N}_4(r, s)P_1(t) \quad (134)$$

$$N_{11}(r, s, t) = \tilde{N}_5(r, s)P_1(t) \quad (135)$$

$$N_{12}(r, s, t) = \tilde{N}_6(r, s)P_1(t) \quad (136)$$

$$N_{13}(r, s, t) = \tilde{N}_1(r, s)P_2(t) \quad (137)$$

$$N_{14}(r, s, t) = \tilde{N}_2(r, s)P_2(t) \quad (138)$$

$$N_{15}(r, s, t) = \tilde{N}_3(r, s)P_2(t) \quad (139)$$

$$N_{16}(r, s, t) = \tilde{N}_4(r, s)P_2(t) \quad (140)$$

$$N_{17}(r, s, t) = \tilde{N}_5(r, s)P_2(t) \quad (141)$$

$$N_{18}(r, s, t) = \tilde{N}_6(r, s)P_2(t) \quad (142)$$

Remark

X. From the previous two examples, it should be clear that when the nodal pattern is the same for each constant- t plane (i.e. is of Lagrange type), then the three-dimensional shape functions are simply products of the two-dimensional shape functions in r and s , and the standard one-dimensional polynomial shape functions in t . Consequently, we shall not consider further patterns of this type.

Examples

17. In this example we consider the same nodal pattern as in the previous example; however, nodes 16 to 18 (i.e. the midsurface nodes) are to be omitted. (This results in a serendipity-type element.) The shape functions may be defined as follows:

$$N_7(r, s, t) = \tilde{N}_4(r, s)t \quad (143)$$

$$N_8(r, s, t) = \tilde{N}_5(r, s)t \quad (144)$$

$$N_9(r, s, t) = \tilde{N}_6(r, s)t \quad (145)$$

$$N_{10}(r, s, t) = \tilde{N}_4(r, s)(1-t) \quad (146)$$

$$N_{11}(r, s, t) = \tilde{N}_5(r, s)(1-t) \quad (147)$$

$$N_{12}(r, s, t) = \tilde{N}_6(r, s)(1-t) \quad (148)$$

$$N_{13}(r, s, t) = \tilde{N}_1(r, s)P_2(t) \quad (149)$$

$$N_{14}(r, s, t) = \tilde{N}_2(r, s)P_2(t) \quad (150)$$

$$N_{15}(r, s, t) = \tilde{N}_3(r, s)P_2(t) \quad (151)$$

In the above, the $N_a(r, s)$'s, $\tilde{N}_a(r, s)$'s and $P_2(t)$ are the same as defined in the previous two examples. A set of preliminary shape functions for the corner nodes are given by

$$N_1(r, s, t) = N_1(r, s)t \quad (152)$$

$$N_2(r, s, t) = N_2(r, s)t \quad (153)$$

$$N_3(r, s, t) = N_3(r, s)t \quad (154)$$

$$N_4(r, s, t) = N_1(r, s)(1-t) \quad (155)$$

$$N_5(r, s, t) = N_2(r, s)(1-t) \quad (156)$$

$$N_6(r, s, t) = N_3(r, s)(1-t) \quad (157)$$

The preceding definitions need to be modified to satisfy the interpolation property. To this end the $N_a(r, s, t)$'s are used in (42), in which $C = \{1, 2, \dots, 6\}$.

We may again mention that the definition of serendipity-type elements is not unique in the present context and other possibilities arise.

18. Consider the 24-node serendipity-type wedge depicted in Figure 8(a). The three-dimensional shape functions are to be generated from the two-dimensional ones associated with the nodal pattern of Figure 8(b) (cf. Example 11). The non-corner shape functions are defined

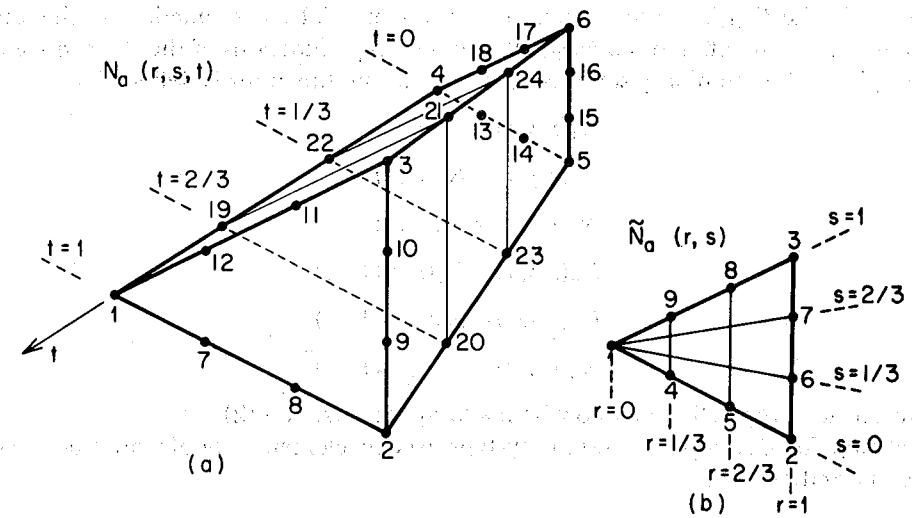


Figure 8. (a) 24-node wedge element. (b) Nine-node triangular element

by:

$$N_7(r, s, t) = \tilde{N}_4(r, s)t \quad (158)$$

$$N_8(r, s, t) = \tilde{N}_5(r, s)t \quad (159)$$

$$N_9(r, s, t) = \tilde{N}_6(r, s)t \quad (160)$$

$$N_{10}(r, s, t) = \tilde{N}_7(r, s)t \quad (161)$$

$$N_{11}(r, s, t) = \tilde{N}_8(r, s)t \quad (162)$$

$$N_{12}(r, s, t) = \tilde{N}_9(r, s)t \quad (163)$$

$$N_{13}(r, s, t) = \tilde{N}_4(r, s)(1-t) \quad (164)$$

$$N_{14}(r, s, t) = \tilde{N}_5(r, s)(1-t) \quad (165)$$

$$N_{15}(r, s, t) = \tilde{N}_6(r, s)(1-t) \quad (166)$$

$$N_{16}(r, s, t) = \tilde{N}_7(r, s)(1-t) \quad (167)$$

$$N_{17}(r, s, t) = \tilde{N}_8(r, s)(1-t) \quad (168)$$

$$N_{18}(r, s, t) = \tilde{N}_9(r, s)(1-t) \quad (169)$$

$$N_{19}(r, s, t) = N_1(r, s)P_3(t) \quad (170)$$

$$N_{20}(r, s, t) = N_2(r, s)P_3(t) \quad (171)$$

$$N_{21}(r, s, t) = N_3(r, s)P_3(t) \quad (172)$$

$$N_{22}(r, s, t) = N_1(r, s)P_2(t) \quad (173)$$

$$N_{23}(r, s, t) = N_2(r, s)P_2(t) \quad (174)$$

$$N_{24}(r, s, t) = N_3(r, s)P_2(t) \quad (175)$$

In (158)–(175), the $P_a(t)$'s are the standard cubic polynomial shape functions of the four-node one-dimensional element, and the $N_a(r, s)$'s are the shape functions of the three-node triangle (i.e. Example 5). A preliminary set of corner-node shape functions is defined by

$$N_1(r, s, t) = N_1(r, s)t \quad (176)$$

$$N_2(r, s, t) = N_2(r, s)t \quad (177)$$

$$N_3(r, s, t) = N_3(r, s)t \quad (178)$$

$$N_4(r, s, t) = N_1(r, s)(1-t) \quad (179)$$

$$N_5(r, s, t) = N_2(r, s)(1-t) \quad (180)$$

$$N_6(r, s, t) = N_3(r, s)(1-t) \quad (181)$$

Modifications to (176)–(181) may now be made by recourse to (42).

The scheme for developing serendipity-type wedge elements should now be clear from Examples 17 and 18.

Remark on implementation

The simplest elements can have their shape functions developed in closed form, as illustrated in the previous section. The automatic generation of these shape functions is complicated by the normalizing terms in the denominators of (2) and (4).

A practical approach is to include one r^β term in addition to the Lagrange family of r^α terms. Then the final function would contain the $1, r^\alpha, \dots, r^{m\alpha}, r^\beta$ powers. This would allow for standard polynomials as well as the common practical cases cited earlier. Such a procedure yields more efficient computations and does not overly restrict the application of the procedure.

It may also be mentioned that the 'variable-number-of-nodes' concept (see e.g. Reference 17) may be fruitfully employed to unify the implementation of families of special elements.

CONCLUSIONS

In this paper we have constructed a concise and efficient algorithm for generating shape functions from an arbitrary set of independent functions. The algorithm does not involve a coefficient matrix and thus eliminates the burdens of equation solving usually associated with interpolation problems.

The algorithm has been employed to generate one-dimensional shape functions which serve as a basis for modelling singularities. A variety of two- and three-dimensional elements have been constructed, and methodology for developing more general elements of the type considered has been described.

The techniques developed herein may also be applied fruitfully to the development of 'special' finite element interpolatory schemes for other purposes. Areas which may be mentioned in this regard are 'upwind' elements, infinite elements, boundary-layer elements and other singular situations (e.g. ones requiring higher order continuity.)

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LEARN - a Linear Static Finite Element Analysis Program

by

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Introduction

This report describes "LEARN," a simple finite element code for linear static analysis. For instructional purposes it may be considered a "first example" in that taping (disk reading and writing) is minimized, and blocking (segmenting large jobs via taping) and overlay structures are avoided. Also the class of capabilities--linear static analysis--is the simplest. Nevertheless many sophisticated procedures are employed (namely compacted column solution and data structures, element routines, etc.) which are used in more advanced situations.

LEARN is used in courses Ae/AM108a-c, Finite Element Methods, and serves the following three purposes:

- (1) It is felt that the first step in learning to implement finite element methods is to add a new element to an existing program. In this regard, LEARN is written so that it is very easy to add new elements. (Instructions are given in the section entitled "Adding New Elements to LEARN.") Students in the class are assigned projects, depending on their interests and level of competence in programming, which usually involve adding a new element to LEARN. For this reason, only two sets of element routines are included with the listing so that many standard element types (e.g. 3D and axisymmetric elasticity elements; beam, plate, and shell elements; heat conduction elements; etc.) are available for project assignments. The course lectures provide much additional information regarding coding element routines.

(2) LEARN makes available to the student a good, first, finite element "software package." The subroutines are efficiently and simply coded and, as the student becomes more advanced, may be used as a set of building blocks for more sophisticated applications. With this idea in mind it was decided to include a very efficient equation solver in LEARN which employs compacted column storage mode (see Section 2).

(3) The style of coding in LEARN derives heavily from the "Berkeley school" of finite element programming which the author learned by studying the works of E. L. Wilson, R. L. Taylor and many former colleagues and students, too numerous to mention, in the Division of Structural Engineering and Structural Mechanics, Department of Civil Engineering, University of California, Berkeley. As this style of programming is employed in many large-scale production codes that are used extensively throughout the world, it is felt to be a very worthwhile style for the student to become familiar with.

The reader of this report is forewarned that the level of documentation is not considered self-sufficient (nor is it intended to be) for fully understanding all aspects of the program, but rather as supplementary to the course lectures.

If any errors are detected in the code or documentation, please report them to T. Hughes, 221 Thomas Laboratory, x1232.

Description of Coding Techniques Used in LEARN

To understand the workings of a computer program, and in order to work with a computer program, one must be thoroughly familiar with its data structure.

Definition. The data structure is the mode of storage, location and history of existence of the data employed.

In LEARN all but the element data is stored in core (central memory) in blank common. The basic features of the data structure are:

- compacted column architecture

This feature is engendered by the method of equation solving employed. The solver is called variously a "profile," "skyline" or "active column" solver. Since, in large problems, the dominant portion of storage is that devoted to the stiffness matrix, $\underline{K}$, the storage scheme for $\underline{K}$ is the most important feature of the storage.

- dynamic storage allocation

Because of the variable sizes of the arrays needed in different problems, it would be very inefficient to "fix" the dimensioning. In the dynamic storage allocation concept the dimensions of arrays are set in the program at the time of execution.

- element groups

Many features of finite element computer programs are the same even though the intended applications may be quite different. To some extent this is also true of the individual finite element subroutines. However, the data structure and options available may vary considerably from one element to another. To accommodate these differences, the elements are read in, stored and operated upon in groups. The data structure for the group is set up via dynamic storage allocation in individual element subroutines.

Example. Compacted column storage scheme.

Assume the symmetric stiffness matrix $\tilde{K}$ has dimension 8×8 and has zero and nonzero elements indicated by O and X, respectively, in the figure below.

$$\tilde{K} = [\tilde{K}_{ij}] = \begin{matrix} & \begin{matrix} 1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 \end{matrix} \\ \begin{matrix} 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \\ 8 \end{matrix} & \begin{bmatrix} X & X & O & O & O & O & O & X \\ & X & X & O & X & O & O & O \\ & & X & X & X & O & O & O \\ & & & X & O & X & O & X \\ & & & & X & X & O & O \\ & & & & & X & X & O \\ & & & & & & X & X \\ & & & & & & & X \end{bmatrix} \end{matrix}$$

Symm.

The dashed line enveloping the nonzero terms is sometimes called the "profile" or "skyline."

$\tilde{K}$ is stored in a one-dimensional array, A, column-wise beginning with the first nonzero element in each column and ending with the diagonal term. The locations of the diagonal terms in A are stored in a one-dimensional integer array IDIAG. The dimension of IDIAG is NEQ, the number of equations (e.g., for $\tilde{K}$ above NEQ=8). The dimension of A is NA, the sum of the "column heights." The column height is the number of terms in a column beginning with the first nonzero term and ending with the diagonal term (e.g., for the above matrix $\tilde{K}$, column 1 has height 1, column 4 has height 2, column 6 has height 3, etc.) NA for $\tilde{K}$ is 24. The arrays A and IDIAG corresponding to $\tilde{K}$ are given explicitly as follows:

$$\begin{array}{ll}
 A(1) = K_{11} & \left. \vphantom{\begin{array}{l} A(1) \\ A(2) \\ A(3) \end{array}} \right\} \text{col. 1} \\
 A(2) = K_{12} & \left. \vphantom{\begin{array}{l} A(2) \\ A(3) \end{array}} \right\} \text{col. 2} \\
 A(3) = K_{22} & \\
 A(4) = K_{23} & \left. \vphantom{\begin{array}{l} A(4) \\ A(5) \end{array}} \right\} \text{col. 3} \\
 A(5) = K_{33} & \\
 A(6) = K_{34} & \left. \vphantom{\begin{array}{l} A(6) \\ A(7) \end{array}} \right\} \text{col. 4} \\
 A(7) = K_{44} & \\
 A(8) = K_{25} & \left. \vphantom{\begin{array}{l} A(8) \\ A(9) \\ A(10) \\ A(11) \end{array}} \right\} \text{col. 5} \\
 A(9) = K_{35} & \\
 A(10) = K_{45} & \\
 A(11) = K_{55} & \\
 A(12) = K_{46} & \left. \vphantom{\begin{array}{l} A(12) \\ A(13) \\ A(14) \end{array}} \right\} \text{col. 6} \\
 A(13) = K_{56} & \\
 A(14) = K_{66} &
 \end{array}$$

$$\begin{array}{l}
 A(15)=K_{67} \\
 A(16)=K_{77}
 \end{array}
 \left. \vphantom{\begin{array}{l} A(15)=K_{67} \\ A(16)=K_{77} \end{array}} \right\} \text{col. 7}$$

$$\begin{array}{l}
 A(17)=K_{18} \\
 A(18)=K_{28} \\
 A(19)=K_{38} \\
 A(20)=K_{48} \\
 A(21)=K_{58} \\
 A(22)=K_{68} \\
 A(23)=K_{78} \\
 A(24)=K_{88}
 \end{array}
 \left. \vphantom{\begin{array}{l} A(17)=K_{18} \\ A(18)=K_{28} \\ A(19)=K_{38} \\ A(20)=K_{48} \\ A(21)=K_{58} \\ A(22)=K_{68} \\ A(23)=K_{78} \\ A(24)=K_{88} \end{array}} \right\} \text{col. 8}$$

$$\begin{array}{l}
 IDIAG(1)=1 \\
 IDIAG(2)=3 \\
 IDIAG(3)=5 \\
 IDIAG(4)=7 \\
 IDIAG(5)=11 \\
 IDIAG(6)=14 \\
 IDIAG(7)=16 \\
 IDIAG(8)=24
 \end{array}$$

The zero terms beneath the skyline must be stored because they become nonzero during the solution process. All information in the original matrix $\tilde{K}$ is now contained in the arrays A and IDIAG in compact fashion. For large finite element stiffness matrices the saving of storage is considerable. ■■■

The heights of the columns in $\tilde{K}$ are determined by subroutine COLHT from the node numbers on element data cards. The column

is LM

heights are initially stored in IDIAG. When the calculation of column heights is completed, subroutine DIAG determines the diagonal addresses and "over writes" them in IDIAG.

The right-hand side vector in

$$\begin{matrix} K & d & = & F \\ \sim & \sim & & \sim \end{matrix}$$

is stored in an array B, of dimension NEQ, in the obvious way:

$$\begin{matrix} B(1) & = & F_1 \\ B(2) & = & F_2 \\ & \cdot & \cdot \\ & \cdot & \cdot \\ & \cdot & \cdot \\ B(8) & = & F_8 \end{matrix}$$

The elimination scheme employed is based upon the following theorem.

Theorem. Let $\begin{matrix} K \\ \sim \end{matrix}$ be symmetric. Then there exists a nonsingular upper triangular matrix $\begin{matrix} U \\ \sim \end{matrix}$, with unit diagonal entries, and a diagonal matrix $\begin{matrix} D \\ \sim \end{matrix}$ such that

$$\begin{matrix} K \\ \sim \end{matrix} = \begin{matrix} U \\ \sim \end{matrix}^T \begin{matrix} D \\ \sim \end{matrix} \begin{matrix} U \\ \sim \end{matrix} \quad (\text{factorization})$$

Furthermore, the number of negative, zero and positive eigenvalues of $\begin{matrix} K \\ \sim \end{matrix}$ is equal to the number of negative, zero and positive diagonal entries of $\begin{matrix} D \\ \sim \end{matrix}$, respectively.

Remarks. 1. In the case when $\tilde{K}$ is positive definite, the diagonal entries of $\tilde{D}$ are all strictly positive. Thus examining the diagonal entries of $\tilde{D}$ enables us to determine the rank of $\tilde{K}$. This is the purpose of the "rank check option" in the program.

2. The matrix $\tilde{U}$ has the same profile as $\tilde{K}$. The fact that $\tilde{U}$ has diagonal entries equal to one allows us to save storage by placing the diagonal entries of $\tilde{D}$ in the diagonal entries of the array A. Thus after factorization the nonzero entries of both $\tilde{U}$ and $\tilde{D}$ are stored in array A. During the factorization procedure no auxiliary storage is needed.

Crout elimination

The procedure used for performing the factorization is called Crout elimination, a convenient variant of Gauss elimination. In the Crout algorithm one column at a time is completely factorized, beginning with the first column and not involving subsequent columns. (This feature proves convenient in nonlinear analysis in which only a small zone of the entire effective stiffness is nonlinear. The linear portion can be located in the first columns and factorized once and for all.)

For further information on the Crout elimination algorithm employed see R. L. Taylor, "Computer Procedures for Finite Element Analysis," Chapter 24, in O. C. Zienkiewicz, The Finite Element Method, Third Edition, McGraw-Hill, London, 1977.

Forward reduction and back substitution

After factorization the solution is carried out in three steps

$$\begin{matrix} U \\ \sim \end{matrix} \begin{matrix} T \\ \sim \end{matrix} \begin{matrix} x \\ \sim \end{matrix} = \begin{matrix} F \\ \sim \end{matrix} \quad (\text{forward reduction})$$

$$\begin{matrix} D \\ \sim \sim \end{matrix} y = \begin{matrix} x \\ \sim \end{matrix}$$

$$\begin{matrix} U \\ \sim \sim \end{matrix} d = \begin{matrix} y \\ \sim \end{matrix} \quad (\text{back substitution})$$

In each step the solution is an explicit process, the coefficient matrix being triangular or diagonal. The intermediate vectors $\begin{matrix} x \\ \sim \end{matrix}$ and $\begin{matrix} y \\ \sim \end{matrix}$ are in turn stored in the array B. Thus again no additional storage besides that for A and B (and IDIAG) is needed to carry out the factorization and solution process.

Storage in blank common

Except for element group data, which is stored on a disk file and read in and out of core as needed, all of the main arrays are stored in blank common. This includes both integer and floating point data.

In the main program and subroutine LEARN, blank common is denoted by the one-dimensional array A.

The locations and lengths of arrays in blank common are as follows:

Storage in Blank Common - Solution Phase

<u>First Word</u>	<u>Array</u>
N1 = 1	ID(NDOF, NUMNP)
N2 = N1 + NDOF*NUMNP	X(NSD, NUMNP)
N3 = N2 + NSD*NUMNP*IPREC	F(NDOF, NUMNP) (=D(NDOF, NUMNP))
N4 = N3 + NDOF*NUMNP*IPREC	IDIAG(NEQ)
N5* = N4 + NEQ	B(NEQ)
N6 = N5 + NEQ*IPREC	A(NA)
N7 = N6 + NA*IPREC	E(MAX)
N8 = N7 + MAX-1 (= total required storage; must be .LT. MTOT for execution.)	

* During input phase, N1 to N5 are the same as above and element group data begins in N5. Individual element routines determine necessary storage for element groups.

where

ID(M,N)	=	equation number for degree of freedom M, node N
X(M,N)	=	Mth coordinate of node N
F(M,N) D(M,N)	} =	Mth direction force/displacement component of node N
IDIAG(N)	=	location of Nth diagonal element in coefficient array A.*
B(N)	=	Nth force/displacement component
A(N)	=	Nth element of coefficient matrix stored in compact column form
E(N)	=	Nth (single precision) word of element group array
NUMNP	=	number of nodes
NSD	=	number of space dimensions
NDOF	=	number of degrees of freedom
NEQ	=	number of equations
NA	=	number of terms in A
MAX	=	number of single-precision words in longest element group data file

*

This array name is used for the coefficient matrix in various subroutines and should not be confused with the array A denoting blank common.

$$\begin{aligned} \text{IPREC}^* &= \begin{cases} 1 & \text{single-precision floating point data} \\ & \text{used throughout program} \\ 2 & \text{double-precision floating point data} \\ & \text{used throughout program} \end{cases} \\ \text{MTOT}^* &= \text{total length of blank common available} \end{aligned}$$

The above names are the ones most commonly used in individual subroutines for the arrays in question. During the input and solution phases different data than that indicated above may be present in some of the arrays to conserve storage.

*Set in the main program. Some FORTRAN compilers allow MTOT to be set internally by the program during execution. However, this is not a feature of standard FORTRAN.

Dynamic storage allocation

Examples

In subroutine LEARN (the driving program) storage is allocated for several of the main arrays. For example, on lines 014 to 016 the "pointers" are determined for the ID and X arrays. The calling statements

```
CALL COORD (A(N2),NSD,NUMNP)
```

and

```
CALL BC      (A(N1),NDOF,NUMNP,NEQ)
```

and corresponding subroutine statements

```
SUBROUTINE COORD(X,NSD,NUMNP)
```

and

```
SUBROUTINE BC (ID,NDOF,NUMNP,NEQ)
```

respectively, mean array X begins in address N2 of blank common and array ID begins in address N1(=1) of blank common.

In COORD and BC, the arrays X and ID are treated as two-dimensional, as evidenced by the DIMENSION statements, viz.

```
DIMENSION X (NSD, 1)
```

and

```
DIMENSION ID (NDOF, 1),
```

respectively. (It does not matter that X and ID correspond to portions of a one-dimensional array in the calling program.) The significance

of the dimensioning is as follows: The arrays X and ID are stored in blank common in consecutively addressed cells (words) by columns.

Consequently the addresses of X(M,N) and ID(M,N) in A are given by (resp.)

$$\text{IPREC} * (\text{NSD} * (\text{N}-1) + \text{M}-1) + \text{N2}$$

and

$$\text{NDOF} * (\text{N}-1) + \text{M}-1 + \text{N1}$$

Thus the only information necessary to determine the appropriate addresses in A are the starting addresses (i.e. N2 and N1, resp.), which are provided by the CALL statements, and the number of rows (i.e. NSD and NDOF, resp.), which are provided by the DIMENSION statements in the subroutines. This explains why it is not necessary to define the number of columns of two-dimensional arrays "passed" to subroutines. The 1's in the column slots of the DIMENSION statements could be replaced by any positive integer. This only serves notice that these arrays are to be regarded as two-dimensional in the subroutines-- the integer is never used.

In the case of three-dimensional arrays, the first two dimensions must be precisely defined, whereas the third one may be set to 1.

Storage of element group data

The individual element routines determine the data storage for element groups. A description of the storage for the truss and plane elements follows:

Storage Requirements of Truss Element (NTYPE=1)

<u>First Word</u>	<u>Array</u>
N101 = NF*	E(NUMAT)
N102 = N101 + NUMAT*IPREC	AREA(NUMAT)
N103 = N102 + NUMAT*IPREC	WT(NUMAT)
N104 = N103 + NUMAT*IPREC	GRAV(3)
N105 = N104 + 3*IPREC	MAT(NUMEL)
N106 = N105 + NUMEL	IEN(2,NUMEL)
N107 = N106 + 2*NUMEL	LM(3, 2,NUMEL)
NL = N107 + 6*NUMEL	
LENGTH = NL - NF (=total required storage for element group)	

where

$E(N)$ = Young's modulus for material number N

$AREA(N)$ = cross-section area for material number N

*NF = N5 during input phase and NF = N7 during solution phase.

WT(N) = weight per unit length for material number N

GRAV(N) = multiplier of gravity load in Nth coordinate
direction

MAT(N) = material number for element number N

IEN(M,N) = Mth node of element number N

LM(L,M,N) = equation number for Lth degree of freedom,
Mth node, element N

NUMAT = number of geometric/material property sets
for this group

Storage Requirements of Plane Stress/Strain Element (NTYPE=3)

<u>First Word</u>	<u>Array</u>
N101 = NF*	E(NUMAT)
N102 = N101 + NUMAT*IPREC	POIS(NUMAT)
N103 = N102 + NUMAT*IPREC	WT(NUMAT)
N104 = N103 + NUMAT*IPREC	TH(NUMAT)
N105 = N104 + NUMAT*IPREC	GRAV(2)
N106 = N105 + 2*IPREC	MAT(NUMEL)
N107 = N106 + NUMEL	IEN(4,NUMEL)
N108 = N107 + 4*NUMEL	LM(2,4,NUMEL)
N109 = N108 + 8*NUMEL	IELNO(NUMPR)
N110 = N109 + NUMPR	ISIDE(NUMPR)
N111 = N110 + NUMPR	PRES(NUMPR)
NL = N111 + NUMPR*IPREC	
LENGTH = NL - NF (=total required storage for element group)	

The meaning of the arrays here may be deduced from the description in the section entitled "User's Manual."

*NF = N5 during input phase and NF = N7 during solution phase.

Adding an element to the program

Example Suppose we wish to add a shell element to the program. Let us call the element routine which sets up element group storage "SHELL."

The steps are as follows:

0. Study and understand thoroughly the subroutines TRUSS, TRUSSI, PLANE and PLANE1 before beginning.
1. Code a new version of subroutine ELMLIB to include a call to SHELL. The necessary modifications are indicated below:

GO TO (1, 2, 3, 4), I

1 CALL TRUSS
RETURN

2 CALL BEAM
RETURN

3 CALL PLANE
RETURN

4	CALL SHELL RETURN
---	----------------------

C
END

2. Write a description for the User's Manual to be included in Section 6.0, ELEMENT DATA. Mimic the style of Sections 6.1 and 6.2, and use the same names of variables as far as possible.

3. Code SHELL in the style of TRUSS and PLANE. (These are the routines which set the element dynamic storage allocation.)
4. Code subroutine SHELL1 in the style of TRUSS1 and PLANE1. Take care to code attractive and complete output data formats.
5. Code any additional subroutines needed which are not already provided in the program.

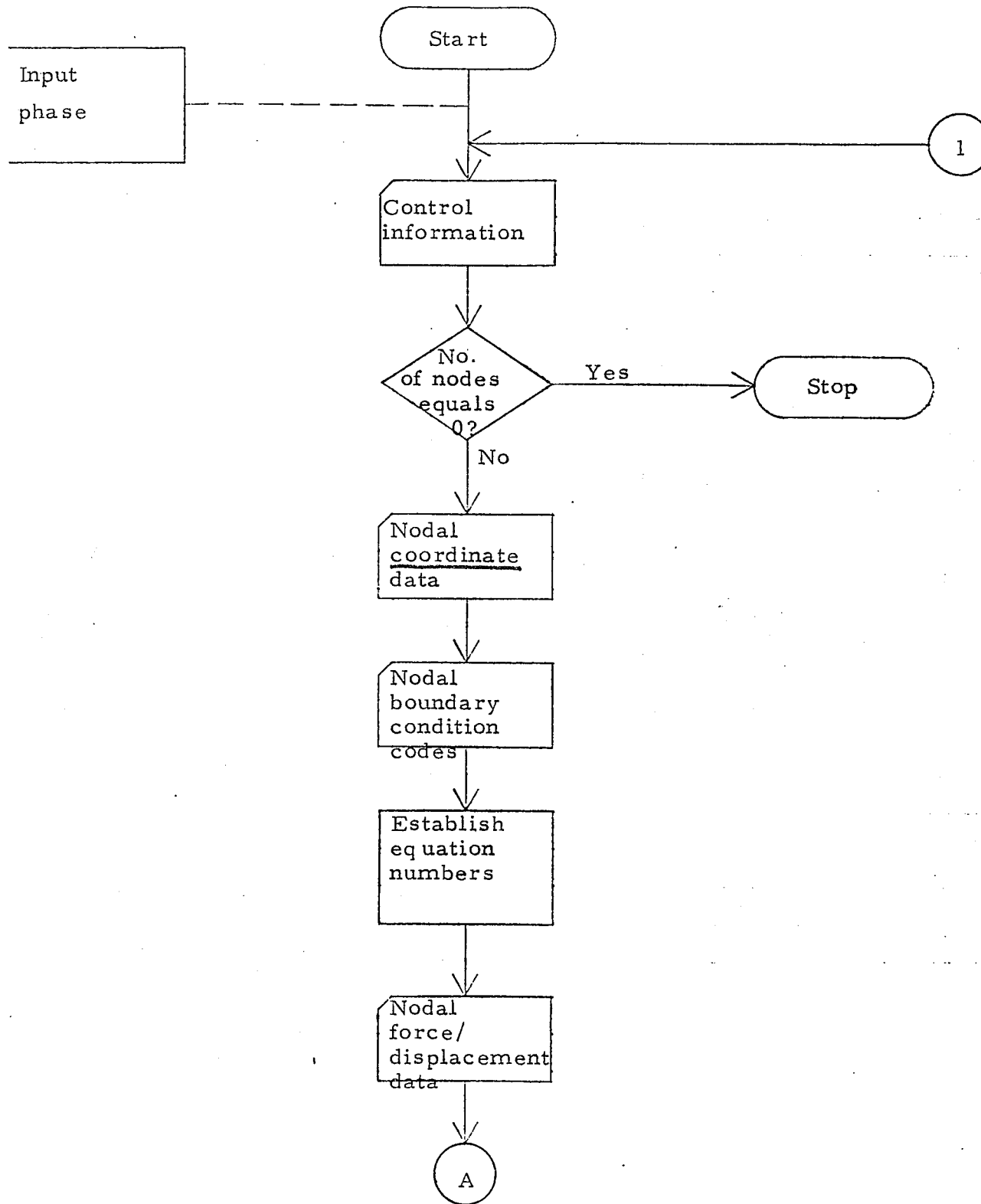
Instructions for debugging* and running the new element

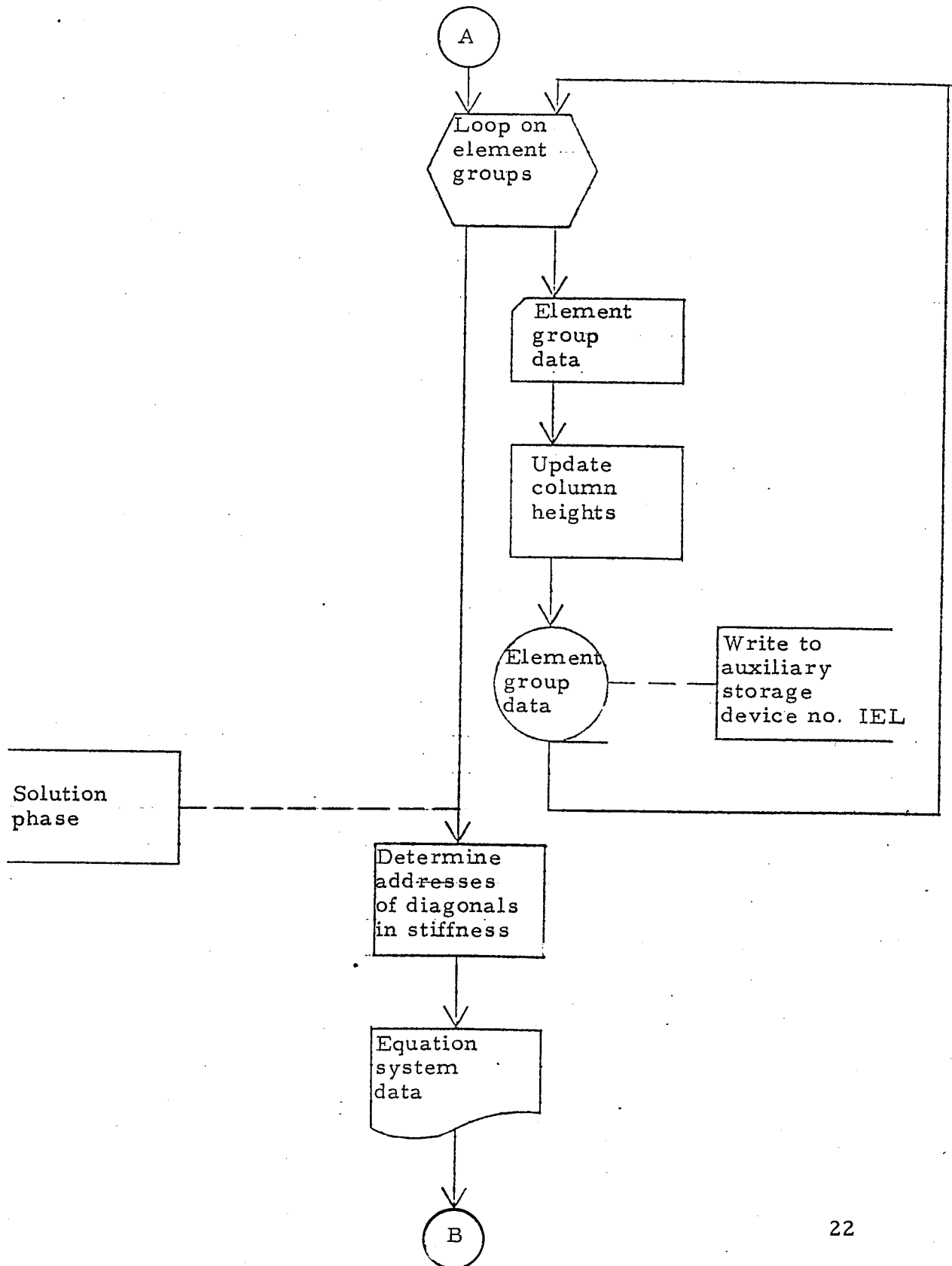
1. Compile all your new routines separately before including them in LEARN.
2. After all routines have compiled, run by loading your new routines with the data for an analysis. You can use the same job control cards as used in running LEARN. (In this way you will not change or destroy the program on permanent file.)
3. Initially you should run simple, one-element problems (e.g. homogeneous deformation states) which are easy to check. Check out all options (e.g., gravity loads, etc.) on the one element problems. Next proceed to "patch tests," which also should be easy to check. Finally, after you are convinced your element is working properly, try some interesting test problems for which you have exact analytical, or reliable experimental, data. Assess the accuracy of the element you have programmed.

* Debugging is boring and frustrating. When you are feeling particularly depressed by this experience, remember everyone who has even programmed has experienced what you are feeling now. Misery loves company.

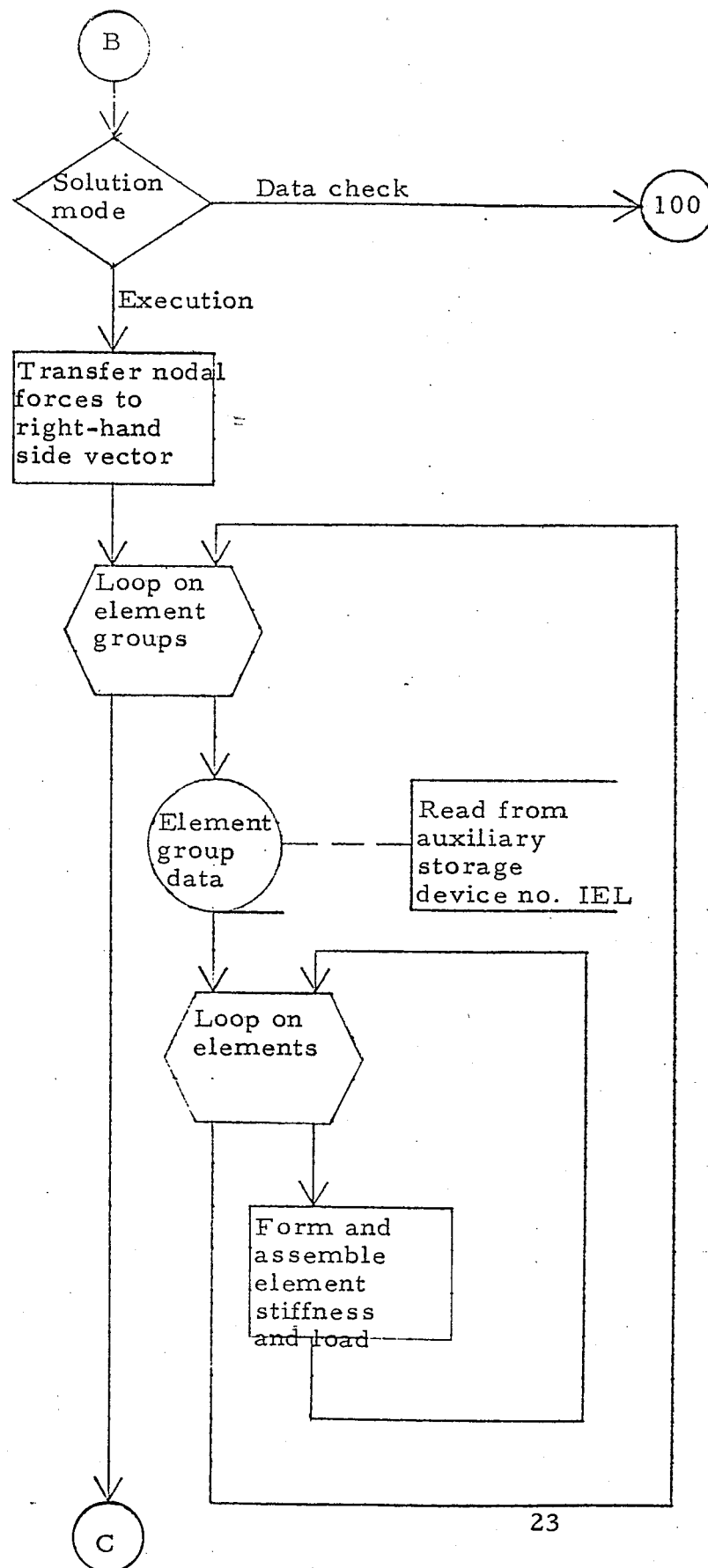
User's Manual

Flowchart

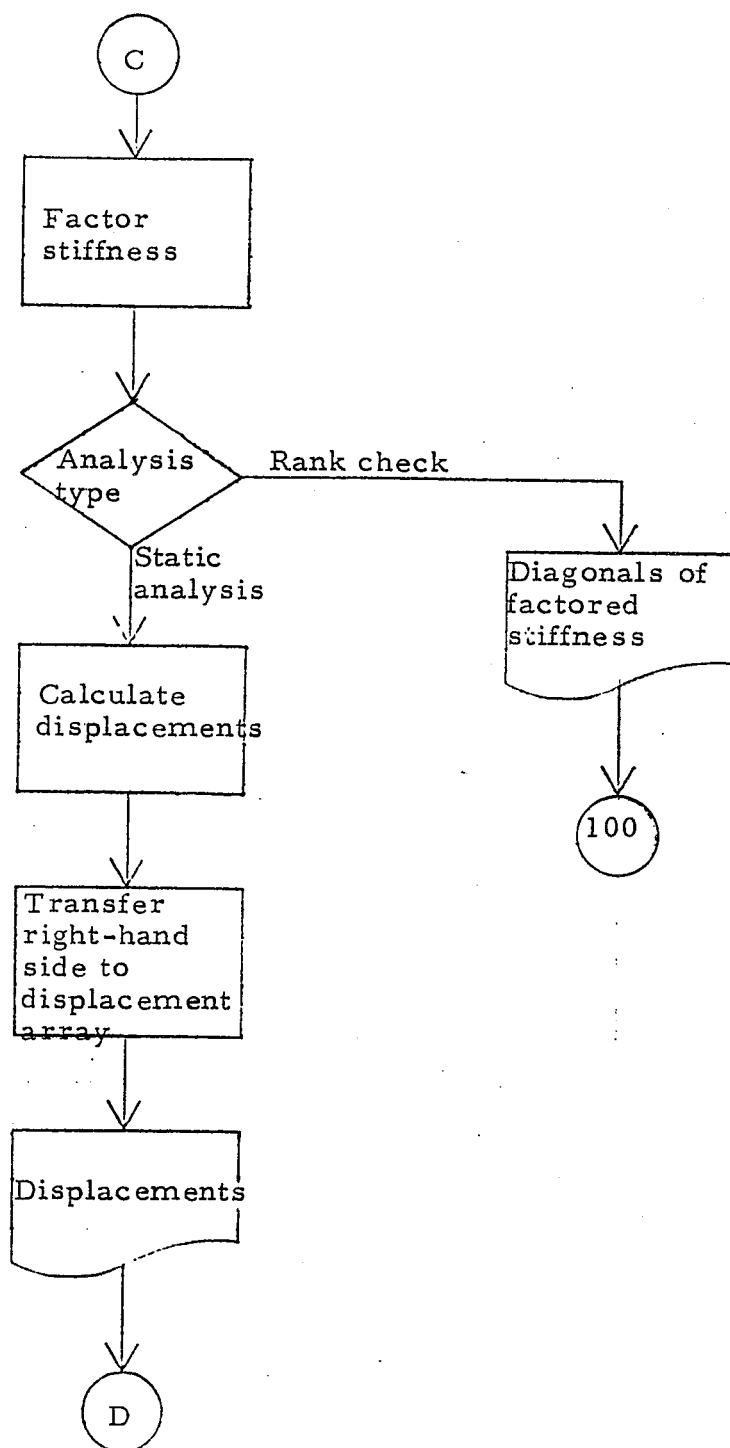




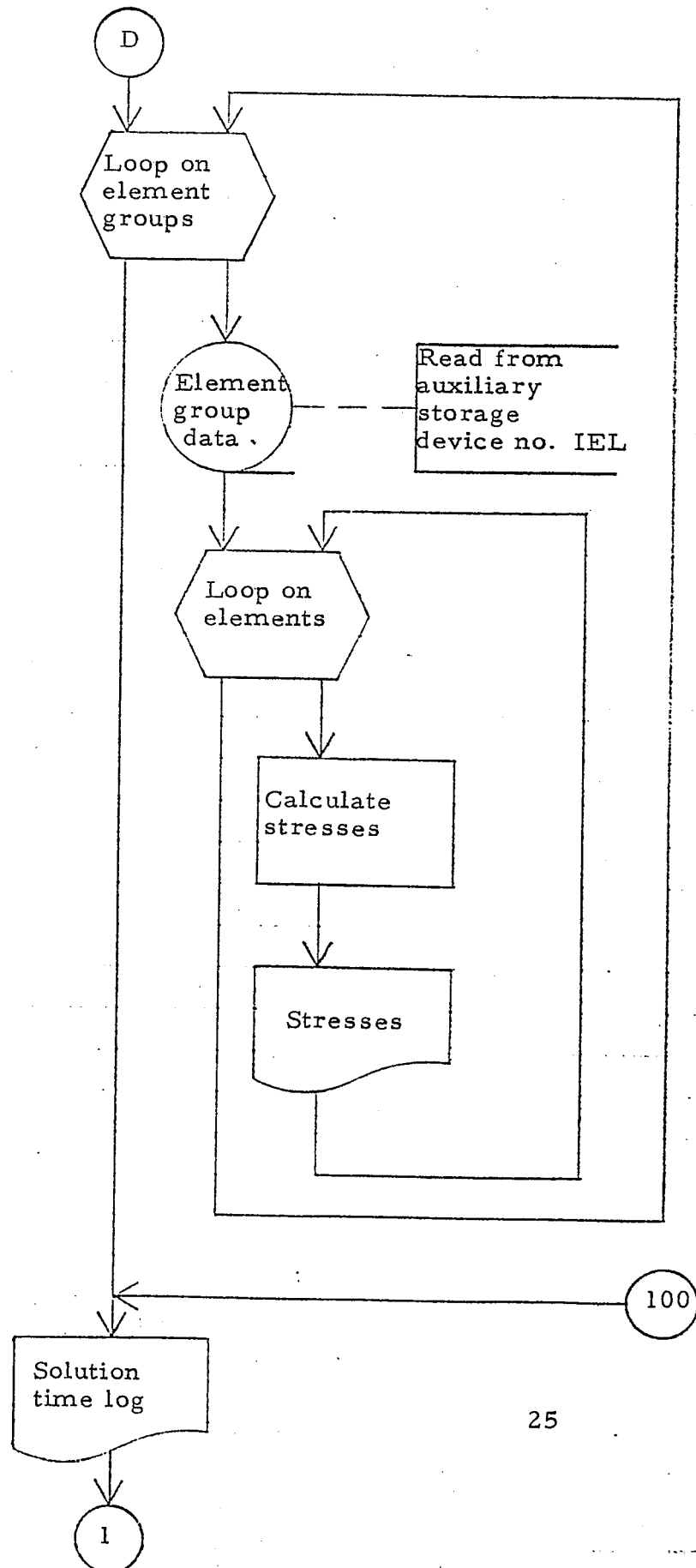
Flowchart (cont'd)



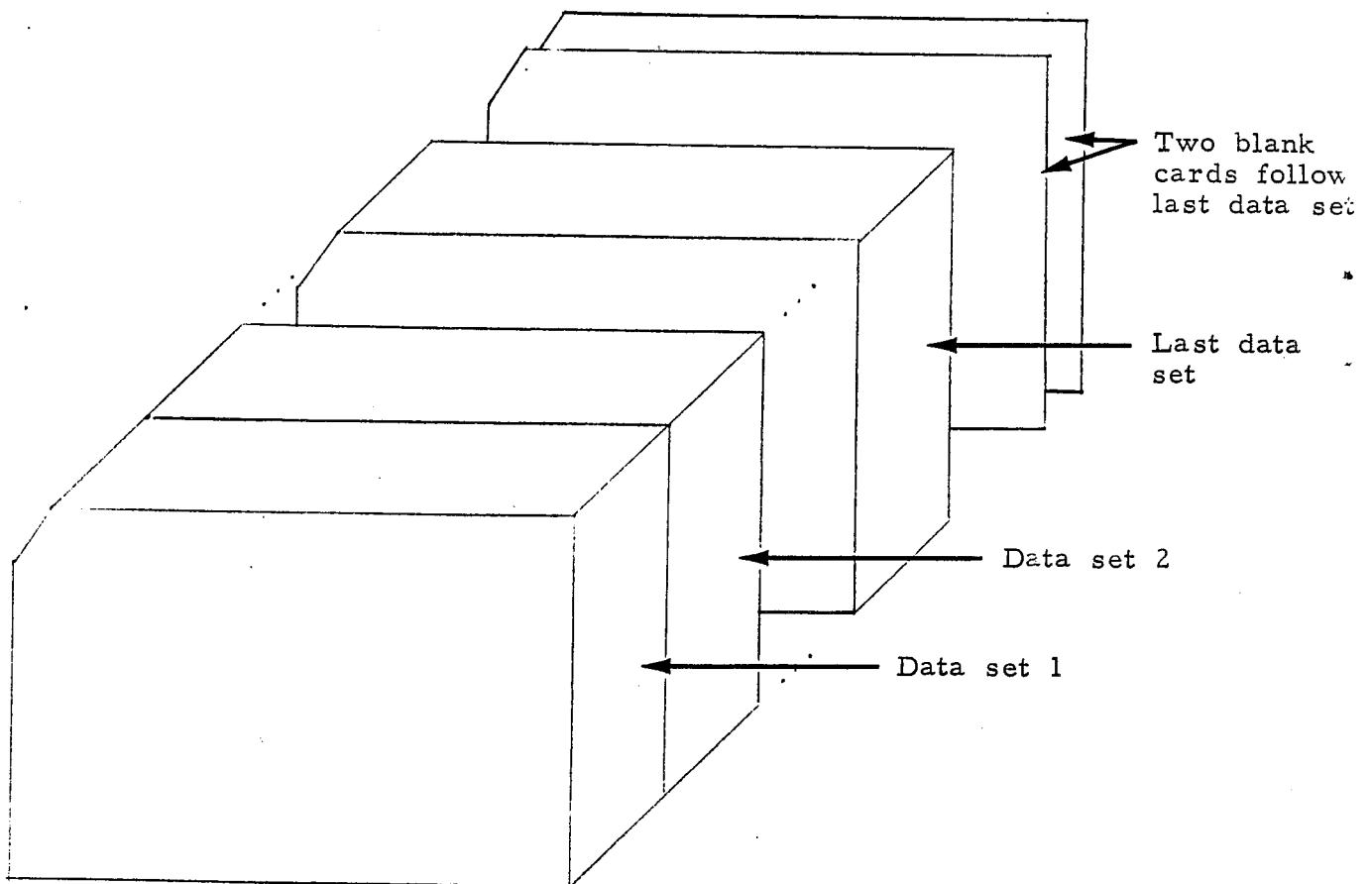
Flowchart (cont'd)



Flowchart (cont'd)



Data Set-Up



Individual Data Cards

1.0 TITLE CARD (20A4)

Note	Columns	Variable	Description
(1)	1-80	TITLE(20)	Job title for heading the output

Notes/

- (1) Each data set begins with a title card. Two blank cards follow the last data set.

2.0 CONTROL CARD (6I5)

Note	Columns	Variable	Description
	1-5	NUMNP	Number of nodes; EQ. 0, program stops
(1)	6-10	NSD	Number of spatial dimensions GE.1 and LE.3
(2)	11-15	NDOF	Number of degrees of freedom at each node; GE.1 and LE.6
(3)	16-20	NUMEG	Number of element groups
(4)	21-25	ICODE	Analysis code EQ. 0, static analysis EQ. 1, stiffness rank check
(5)	26-30	MODE	Execution mode EQ. 0, data check only EQ. 1, execution

Notes/

(1) Consult individual element routines in Section 6 for requirements.

(2) In problems for which there are different numbers of degrees of freedom at different nodes, NDOF should be set equal to the maximum number. Consult individual element routines in Section 6 for requirements. Superfluous degrees of freedom can then be eliminated by employing boundary condition codes (see Section 4.0).

(3) The elements are read in groups.

(4) A static analysis consists of solution of the equilibrium equations, displacement and stress output.

A rank check consists of factorization of the stiffness matrix, $K = U^T D U$, and printing the elements of the diagonal matrix D . The number of zero-energy modes can be deduced from the number of zeroes along the diagonal of D . This option allows for the checking of elements and meshes suspected of containing spurious zero-energy modes, such as those arising in so-called "underintegrated finite elements."

2.0 CONTROL CARD, Page 2

Notes Continued:

- (5) In the data check mode input data is printed out and storage requirements are indicated. This mode should be employed prior to making expensive executions.

3.0 COORDINATE DATA (2I5, NSD*F10.0)

Note	Columns	Variable	Description
(1)	1-5	N	Node number; GE.1 and LE.NUMNP
(2)	6-10	NG	Generation increment
(3)	11-20	X(1,N)	x_1 -coordinate of node N
	21-30	X(2,N)	x_2 -coordinate of node N
	31-40	X(3,N)	x_3 -coordinate of node N

Notes/

- (1) The coordinates of each node must be defined, but need not be read in order. Terminate with a blank card.
- (2) Coordinate data can be generated along straight lines by employing a two card sequence as follows:

Card 1: L, LG, X(1, L), ..., X(NSD, L)

Card 2: N, NG, X(1, N), ..., X(NSD, N)

N-L must be an integral multiple of LG. The coordinates of all nodes

$L+LG, L+2*LG, \dots, N-LG$

are generated by linearly interpolating the coordinates of nodes L and N. If LG is blank or zero no generation takes place between L and N.

- (3) If more than one coordinate data card for node N is input, the last one read takes priority.

4.0 BOUNDARY CONDITION DATA ((2+NDOF)*I5)

Note	Columns	Variable	Description
(1)	1-5	N	Node number; GE.1 and LE.NUMNP
(2)	6-10	NG	Generation increment
(3)	11-15	ID(1, N)	Degree of freedom 1 boundary code
	16-20	ID(2, N)	Degree of freedom 2 boundary code
	etc.	$\vdots$	$\vdots$
		ID(NDOF, N)	Degree of freedom NDOF boundary code

Notes/

(1) Boundary condition data must be input for each node which has one or more specified displacements. Cards need not be input in order. Terminate with a blank card.

(2) Boundary condition data can be generated by employing a two card sequence as follows:

Card 1: L, LG, ID(1, L), ..., ID(NDOF, L)

Card 2: N, NG, ID(1, N), ..., ID(NDOF, N)

The boundary codes of all nodes

$L+LG, L+2*LG, \dots, N-MOD(N-L, LG)$

(i.e., less than N) are set equal to those of node L. If LG is blank or zero no generation takes place between L and N.

(3) Boundary condition codes may be assigned the following values:

ID(I, N) = 0 , unspecified displacement

ID(I, N) = 1 , specified displacement

4.0 BOUNDARY CONDITION DATA, Page 2

Notes Continued:

where $I = 1, 2, \dots, \text{NDOF}$. Specified displacements are assumed to be fixed (i.e., have the value 0.0) unless assigned a nonzero value as described in Section 5.0. If more than one boundary condition data card for node N is input, the last one read takes priority.

5.0 APPLIED NODAL FORCES AND PRESCRIBED DISPLACEMENTS (2I5, NDOF*F10.0)

Note	Columns	Variable	Description
(1)	1-5	N	Node number; GE.1 and LE.NUMNP
(2)	6-10	NG	Generation increment
(3)	11-20	F(1,N)	Degree of freedom 1 force or displacement
	21-30	F(2,N)	Degree of freedom 2 force or displacement
	etc.	$\vdots$	$\vdots$
		F(NDOF,N)	Degree of freedom NDOF force or displacement

Notes/

(1) Applied nodal force/prescribed displacement data must be included for each node subjected to a nonzero applied force or nonzero prescribed displacement. Cards need not be read in order. Terminate with a blank card.

(2) Applied nodal force and prescribed displacement data may be generated by employing a two card sequence as follows (generation is the same as for coordinate data):

Card 1: L, LG, F(1, L), ..., F(NDOF, L)

Card 2: N, NG, F(1, N), ..., F(NDOF, N)

N-L must be an integral multiple of LG. The applied nodal force and prescribed displacement data of all nodes

$L+LG, L+2*LG, \dots, N-LG$

are generated by linearly interpolating the data of nodes L and N. If LG is blank or zero no generation takes place between L and N.

(3) The elements of the array F(NDOF, NUMNP) are initialized to zero. If more than one applied nodal force/prescribed displacement data card for node N is input, the last one read takes priority.

6.0 ELEMENT DATA

6.1.0 Three-Dimensional Truss Element

Truss elements connect two points in space and transmit axial force only. There are three degrees of freedom at each of the two nodes, i.e. the x_1 , x_2 and x_3 translations. When employing truss elements, NSD must be EQ.3 and NDOF must be GE.3 on the master control card (see Section 1.0). The following sequence of cards is used to describe truss elements:

6.1.1 Element Group Control Card (3I5)

Note	Columns	Variable	Description
	1-5	NPAR(1) (=NTYPE)	The number 1
	6-10	NPAR(2) (=NUMEL)	Number of elements in this group; GE.1
(1)	11-15	NPAR(3) (=NUMAT)	Number of geometric/ material sets in this group; GE.0; if EQ. 0, set internally to 1

Notes/

- (1) Defines number of geometric/material properties cards to be input in Section 6.1.2.

6.1.2 Geometric/Material Properties Cards (I5,5X,3F10.0)

Note	Columns	Variable	Description
(1)	1-5	N	Material identification number
	11-20	E(N)	Young's modulus
	21-30	AREA(N)	Cross-section area
(2)	31-40	WT(N)	Weight per unit length

Notes/

- (1) One card is required for each set of elements possessing the same cross-section and material properties.
- (2) The product of weight per unit length and length defines gravity loads for the truss element.

6.1.3 Gravity Load Multiplier Card (3F10.0)

Note	Columns	Variable	Description
(1)	1-10	GRAV(1)	Multiplier of gravity load in the x_1 direction
	11-20	GRAV(2)	Multiplier of gravity load in the x_2 direction
	21-30	GRAV(3)	Multiplier of gravity load in the x_3 direction

Notes/

- (1) Gravity load multipliers may be used to define the components of the gravity unit vector with respect to the global x_1, x_2, x_3 system. For example, if the gravity load multipliers are

0.0, 0.0, -1.0,

then the entire gravitational load acts in the $-x_3$ direction.
(One card for element group.)

6.1.4 Element Data Cards (5I5)

Note	Columns	Variable	Description
(1)	1-5	N	Element number
	6-10	MAT(N)	Geometric/material properties set number
	11-15	IEN(1,N)	Number of 1st node
	16-20	IEN(2,N)	Number of 2nd node
(2)	21-25	NG	Generation increment

Notes/

(1) All elements must be input on an element data card or generated.
Terminate with a blank card.

(2) Element data cards may be generated by employing a two card sequence as follows:

Card 1: L, MAT(L), IEN(1, L), IEN(2, L), LG

Card 2: N, MAT(N), IEN(1, N), IEN(2, L), NG

N must be greater than L. The geometric/material set number for all elements

$L+1, L+2, \dots, N-1$

is set to MAT(L), and the node numbers are incremented by LG.

6.2.0 Two-Dimensional Isotropic Elasticity Element

The present element may be used in triangle (3 node) or quadrilateral (4 node) form for plane strain and plane stress. In each case the plane of analysis is assumed to be the x_1, x_2 plane. Two displacement degrees of freedom, in the x_1 and x_2 directions, are assigned to each node. Incompatible modes and selective numerical integration may be employed to improve element behavior in various situations. These options should be activated only by users fully knowledgeable in their use. The nodes of the element must be input in counterclockwise order and NSD must EQ.2 and NDOF must be GE.2 on the master control card (see Section 1.0).

Stresses/strains in the global coordinate system, and principal stresses/strains, maximum shear stress/strain and angle of inclination, in degrees, of principal states (see Figure 6.2.0.1) are output at the element centroid, which is generally the point of optimal accuracy. All shear strains are reported according to the "engineering" convention (i.e. twice the value of the tensor components).

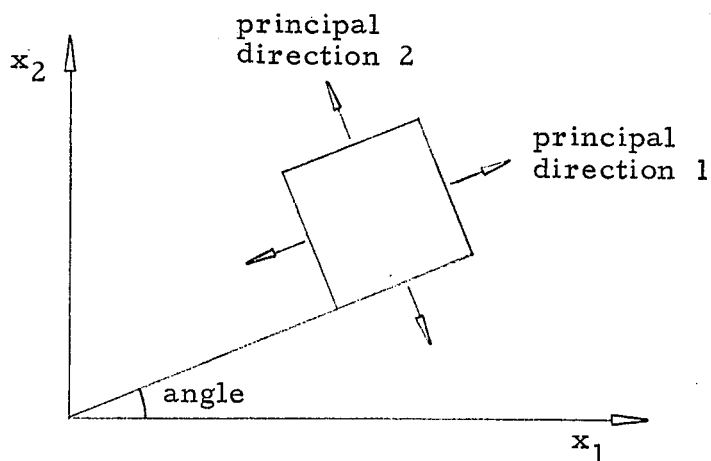


Figure 6.2.0.1

The following sequence of cards is used to describe the elements:

6.2.1 Element Group Control Card (9I5)

Note	Columns	Variable	Description
	1-5	NPAR(1) (=NTYPE)	The number 3
	6-10	NPAR(2) (=NUMEL)	Number of elements in this group; GE. 1
	11-15	NPAR(3) (=NUMAT)	Number of geometric/ material sets in this group; GE. 0; if EQ. 0, set internally to 1
	16-20	NPAR(4) (=IOPT)	Analysis option EQ. 0, plane strain EQ. 1, plane stress
(1)	21-25	NPAR(5) (=IL)	Lambda-term (λ) numerical integration code EQ. 0, two-by-two Gaussian quadrature EQ. 1, one-point Gaussian quadrature
(2)	26-30	NPAR(6) (=IM)	Mhu-term (μ) numerical integration code EQ. 0, two-by-two Gaussian quadrature EQ. 1, one-point Gaussian quadrature
(3)	31-35	NPAR(7) (=INC)	Incompatible modes code EQ. 0, incompatible modes neglected EQ. 1, incompatible modes added
(4)	36-40	NPAR(8) (=IHOURL)	Hourglass stiffness code EQ. 0, hourglass stiffness neglected EQ. 1, hourglass stiffness added
	41-45	NPAR(9) (=NUMPR)	Number of element pressure load cards

6.2.1 Continued

Notes/

- (1) In problems involving nearly incompressible materials an effective approach is to use one-point Gaussian quadrature on the λ term, and two-by-two Gaussian quadrature on the μ term (see T.J.R. Hughes, "Equivalence of Finite Elements for Nearly-Incompressible Elasticity," LBL 5237, University of California, Berkeley, August 1976, or Journal of Applied Mechanics, Vol. 44, Series E, No. 1, p. 181, March 1977).
- (2) The "standard" 4-node quadrilateral employs two-by-two Gaussian quadrature on both the λ and μ terms. However, it is ineffective in application to nearly-incompressible materials (see above note) and also in application to "bending" situations. One-point Gaussian quadrature on both terms produces a more accurate, but dangerous element. The danger is that zero energy modes of deformation, so-called hourglass or keystone modes, may be present, resulting in a singular stiffness matrix. If enough displacement boundary conditions are present these modes may be eliminated. However, one-point Gaussian quadrature on both terms should only be used if you know exactly what you are doing.
- (3) The presence of incompatible modes produces an element effective in bending and in application to incompressible materials. Two-by-two Gaussian quadrature on both the λ and μ terms should be employed when using incompatible modes. Much pro and con has been written about incompatible modes. The present implementation is based upon R.L. Taylor, P.J. Beresford and E.L. Wilson, "A Non-Conforming Element for Stress Analysis," International Journal for Numerical Methods in Engineering, 10 (6) p. 1211, 1976.
- (4) The hourglass stiffness counteracts the zero-energy hourglass modes when one-point Gaussian quadrature is used on both the λ and μ terms. The implementation is based upon D. Kosloff and G. A. Frazier, "Treatment of Hourglass Patterns in Low Order Finite Element Codes," Contribution No. 2895, Division of Geological and Planetary Sciences, California Institute of Technology, Pasadena, California, or Numerical and Analytical Methods in Geomechanics, Vol. 2, 57-72 (1978). The hourglass stiffness should be added only when one-point Gaussian quadrature is employed on both the λ and μ terms.

6.2.2 Geometric/Material Properties Cards (I5,5X,5F10.0)

Note	Columns	Variable	Description
	1-5	N	Material identification number
	11-20	E(N)	Young's modulus
(1)	21-30	POIS(N)	Poisson's ratio
(2)	31-40	WT(N)	Weight density
(3)	41-50	TH(N)	Thickness (if 0.0 defaults to 1.0)

Notes/

- (1) Poisson's ratio cannot be set equal to $1/2$ since it results in division by zero. A value close to $1/2$, say .4999, can be employed for incompressible applications.
- (2) The product of the weight density and element volume is used to compute gravity loads.
- (3) In plane strain the thickness is automatically assumed to be 1.0.

6.2.3 Gravity Load Multiplier Cards (2F10.0)

Note	Columns	Variables	Description
(1)	1-10	GRAV(1)	Multiplier of gravity load in the x_1 direction
	11-20	GRAV(2)	Multiplier of gravity load in the x_2 direction

Notes/

- (1) Gravity load multipliers are used to define the components of the gravity vector with respect to the global x_1, x_2 system.

6.2.4 Element Data Cards (7I5)

Note	Columns	Variables	Description
(1)	1-5	N	Element number
	6-10	MAT(N)	Geometric/material properties set number
	11-15	IEN(1, N)	Number of 1st node
	16-20	IEN(2, N)	Number of 2nd node
	21-25	IEN(3, N)	Number of 3rd node
(2)	26-30	IEN(4, N)	Number of 4th node
(3)	31-35	NG	Generation increment

Notes/

- (1) All elements must be input on an element data card or generated. Terminate with a blank card.
- (2) Node number IEN(4, N) must equal node number IEN(3, N) for triangular elements. Triangles should not be used for incompressible or nearly-incompressible materials.

6.2.4 Continued

Notes Continued:

- (3) Element data cards may be generated by employing a two card sequence as follows:

Card 1: L, MAT(L), IEN(1, L), IEN(2, L), IEN(3, L), IEN(4, L), LG

Card 2: N, MAT(N), IEN(1, N), IEN(2, N), IEN(3, N), IEN(4, N), NG

N must be greater than L. The geometric/material set number for all

L+1, L+2, ..., N-1

is set to MAT(L) and the node numbers are incremented by LG.

6.2.5 Element Pressure Load Cards (2I5, 10X, F10.0)

Note	Columns	Variables	Description
(1)	1-5	IELNO(I)	Element number 1.LE.IELNO(I).LE.NUMEL
(2)	6-10	ISIDE(I)	Element side number 1.LE.ISIDE(I).LE.4
(3)	21-30	PRES(I)	Pressure (force/unit area)

Notes/

- (1) Each element subjected to a pressure load must be input on an element pressure load card. If more than one side of the element is loaded, one card for each loaded side must be input. The total number of element pressure load cards must equal NUMPR read in on the element group control card (see Section 6.2.1). The index I in the arrays IELNO, ISIDE and PRES corresponds to the order read in; 1.LE.I.LE.NUMPR. The cards need not be read in any particular order; terminate with a blank card.
- (2) The element side number is deduced as follows: Consider the element of Figure 6.2.1.1; I, J, K, L indicate the element node numbers as read in on the element data card (i.e., I is the

6.2.5 Continued

Notes Continued:

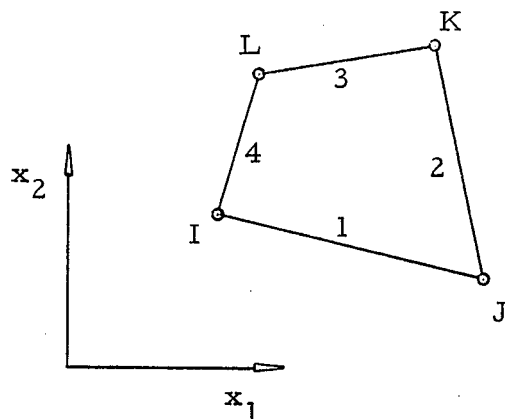


Figure 6.2.1.1

first node, J the second, etc.). Side 1 connects nodes I and J, side 2 connects nodes J and K, etc.; see Figure 6.2.1.1.

- (3) The pressure is assumed to be positive pointing inward, as indicated in Figure 6.2.1.2

note

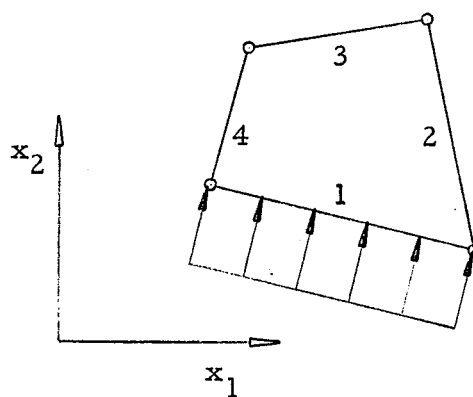
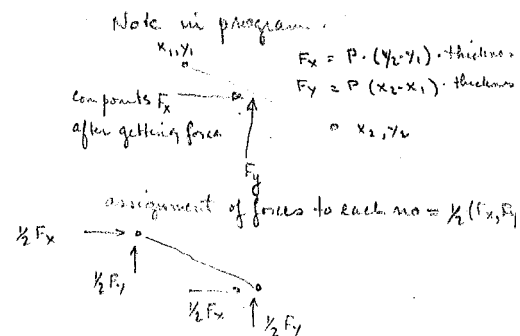


Figure 6.2.1.2



Listing


```

C
C..... MAIN PROGRAM
C
COMMON/ELDATA/IE ,IPREC,LENGTH,MAX,MTOT,NEG,NF,NL,NPAR(16),NUMEG
COMMON/DEVICE/ IEL,IIN,ICUT
COMMON A(5000)
MTOT=5001
IPREC=2
IEL=1
IIN=5 input file
ICUT=6 output file
CALL LEARN
END

C
SUBROUTINE LEARN
C
C *****
C *
C * LEARN
C *
C * A L I N E A R S T A T I C A N A L Y S I S P R O G R A M
C *
C * T . H U G H E S O C T O B E R 1 9 7 6
C *
C * ( D O U B L E P R E C I S I O N V E R S I O N )
C *
C *****
C
COMMON/INFC/ICODL,MODE,NDOF,NSD,NUMNP
COMMON/DEVICE/ IEL,IIN,ICUT
COMMON/POINTS/N1,N2,N3,N4,N5,N6,N7
COMMON/ELDATA/IE ,IPREC,LENGTH,MAX,MTOT,NEG,NF,NL,NPAR(16),NUMEG
COMMON/LABELS/LABELC(3),LABELD(6)
DIMENSION TIME(5),TITLE(20)
COMMON A(1)

C
C DEVICE CODES
C IIN= INPUT
C ICUT=OUTPUT
C IEL= ELEMENT DATA
C
C
C 1 MAX=0
C
C *****
C *
C * I N P U T P H A S E
C *
C *****
C
C..... PROGRAM CONTROL INFORMATION
C
CALL SECOND(TIME(1))
READ(IIN,1000)TITLE,NUMNP,NSD,NDOF,NUMEG,ICODE,MODE
IF(NUMNP.EQ.0) STOP
WRITE(ICUT,2000) TITLE,NUMNP,NSD,NDOF,NUMEG,ICODE,MODE

C
C..... INPUT NODAL DATA
C
N1=1 (N1) = A(N1)
N2=N1+NUMNP*NDOF N2 = A(N2)
N3=N2+NUMNP*NSD*IPREC N3 = A(N3)
IF(N3.GT.MTOT) CALL ERROR(N3-MTOT,1) check size of array for nodal coord & disp data
CALL CCORD(A(N2),NSD,NUMNP) generates nodal points
CALL BC(A(N1),NDOF,NUMNP,NEQ) generates BC

C
C..... INPUT NODAL APPLIED FORCE AND PRESCRIBED DISPLACEMENT DATA
C
N4=N3+NUMNP*NDOF*IPREC N4 = A(N4)
IF(N4.GT.MTOT) CALL ERROR(N4-MTOT,2) check applied forces & displ data

```

```

1 MAIN
2 MAIN
3 MAIN
4 MAIN
5 MAIN
6 MAIN
7 MAIN
8 MAIN
9 MAIN
10 MAIN
11 MAIN
12 MAIN
13 MAIN

1 LEARN
2 LEARN
3 LEARN
4 LEARN
5 LEARN
6 LEARN
7 LEARN
8 LEARN
9 LEARN
10 LEARN
11 LEARN
12 LEARN
13 LEARN
14 LEARN
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55 LEARN
56 LEARN

```

CALL FORCE (A(N3),NDOF,NUMNP)	<i>generate force</i>	
C		57 LEARN
C.... INPUT ELEMENT DATA		58 LEARN
C		59 LEARN
N5=N4+NEQ	A(N4)	60 LEARN
N6=N5		61 LEARN
N7=N5		62 LEARN
IF(N5.GT.MTOT) CALL ERROR(N5-MTOT,3)		63 LEARN
CALL ICLEAR (A(N4),NEQ)		64 LEARN
IE=1		65 LEARN
CALL EL		66 LEARN
CALL SECOND(TIME(2))		67 LEARN
C		68 LEARN
C		69 LEARN
C	*****	70 LEARN
C	*	71 LEARN
C		72 LEARN
C	S O L U T I O N P H A S E	73 LEARN
C	*	74 LEARN
C	*****	75 LEARN
C.... DETERMINE ADDRESSES OF DIAGONALS IN STIFFNESS		76 LEARN
C		77 LEARN
CALL DIAG(A(N4),NEQ,NA)		78 LEARN
N6=N5+NEQ*IPREC	B = A(N5)	79 LEARN
N7=N6+NA *IPREC	A = A(N6)	80 LEARN
N8=N7+MAX-1	E = A(N7)	81 LEARN
IF(N8.GT.MTOT) CALL ERROR(N8-MTOT,5)		82 LEARN
C		83 LEARN
C.... WRITE EQUATION SYSTEM DATA		84 LEARN
C		85 LEARN
MB=NA/NEQ		86 LEARN
WRITE(ICUT,2010) TITLE,NEQ,NA,MB,N8		87 LEARN
C		88 LEARN
C.... IF DATA CHECK MODL (MODE.EQ.0) SUBSEQUENT CALCULATIONS ARE IGNORED		89 LEARN
C		90 LEARN
IF(MODE.EQ.1) GOTO 10		91 LEARN
CALL SECOND (TIME(3))		92 LEARN
CALL SECOND (TIME(4))		93 LEARN
CALL SECOND (TIME(5))		94 LEARN
GOTO 100		95 LEARN
C		96 LEARN
C.... ASSEMBLE STIFFNESS AND LOAD		97 LEARN
C		98 LEARN
10 CALL CLEAR (A(N5),NEQ+NA)		99 LEARN
CALL LOAD(A(N1),A(N3),A(N5),NDOF,NUMNP)		100 LEARN
IE=2		101 LEARN
CALL ELREAD(A(N7))		102 LEARN
CALL SECOND(TIME(3))		103 LEARN
C		104 LEARN
C.... FACTOR STIFFNESS MATRIX		105 LEARN
C		106 LEARN
CALL SOLVE(A(N6),A(N5),A(N4),NEQ,.TRUE.,.FALSE.)		107 LEARN
CALL SECOND (TIME(4))		108 LEARN
C		109 LEARN
C.... IF RANK CHECK ANALYSIS (ICODE.EQ.1) SUBSEQUENT CALCULATIONS ARE IGNORED		110 LEARN
C		111 LEARN
C		112 LEARN
IF(ICODE.EQ.0) GOTO 20		113 LEARN
CALL PRINTP(A(N6),A(N4),NEQ)		114 LEARN
CALL SECOND(TIME(5))		115 LEARN
GOTO 100		116 LEARN
C		117 LEARN
C.... CALCULATION OF DISPLACEMENTS		118 LEARN
C		119 LEARN
20 CALL SOLVE(A(N6),A(N5),A(N4),NEQ,.FALSE.,.TRUE.)		120 LEARN
CALL DISPTR(A(N5),A(N3),A(N1),NDOF,NUMNP)		121 LEARN
CALL PRINTD(A(N3),NDOF,NUMNP)		122 LEARN
C		123 LEARN
C.... CALCULATION OF STRESSES		124 LEARN
C		125 LEARN
IE=3		126 LEARN
CALL ELREAD(A(N7))		127 LEARN
CALL SECOND(TIME(5))		128 LEARN

C		129 LEARN
C	C.... PRINT SOLUTION TIMES	130 LEARN
C	100 TOT=TIME(1)-TIME(5)	131 LEARN
	DO 30 I=1,4	132 LEARN
	30 TIME(I)=TIME(I)-TIME(I+1)	133 LEARN
	WRITE(ICUT,2020) TITLE,(TIME(I),I=1,4),TOT	134 LEARN
C		135 LEARN
C	C.... READ NEXT ANALYSIS CASE	136 LEARN
C		137 LEARN
	GOTO 1	138 LEARN
C		139 LEARN
	1000 FORMAT(20A4/6I5)	140 LEARN
	2000 FORMAT(1H1,20A4///	141 LEARN
	X' C U N T R O L I N F O R M A T I O N	142 LEARN
	X' NUMBER OF NODAL POINTS (NUMNP) =',15//5X,	143 LEARN
	X' NUMBER OF SPACE DIMENSIONS (NSD) =',15//5X,	144 LEARN
	X' NUMBER OF DEGREES OF FREEDOM PER NODE (NDOF) =',15//5X,	145 LEARN
	X' NUMBER OF ELEMENT GROUPS (NUMEG) =',15//5X,	146 LEARN
	X' ANALYSIS CODE (ICODE) =',15/ 5X,	147 LEARN
	X' EQ.0, STATIC ANALYSIS	148 LEARN
	X' EQ.1, RANK CHECK	149 LEARN
	X' SOLUTION MODE (MODE) =',15/ 5X,	150 LEARN
	X' EQ.0, DATA CHECK	151 LEARN
	X' EQ.1, EXECUTION	152 LEARN
	2010 FORMAT(1H1,20A4///	153 LEARN
	X' E Q U A T I O N S Y S T E M D A T A	154 LEARN
	X' NUMBER OF EQUATIONS (NEQ) =',15//5X,	155 LEARN
	X' NUMBER OF TERMS IN STIFFNESS (NA) =',15//5X,	156 LEARN
	X' MEAN HALF BANDWIDTH (MB) =',15//5X,	157 LEARN
	X' TOTAL LENGTH OF BLANK COMMON REQUIRED (N8) =',15	158 LEARN
	2020 FORMAT(1H1,20A4///	159 LEARN
	X' S O L U T I O N T I M E L O G I N S E C O N D S' //5X,	160 LEARN
	X' INPUT PHASE =',F12.2//5X,	161 LEARN
	X' FORMATION OF STIFFNESS AND LOAD =',F12.2//5X,	162 LEARN
	X' FACTOR STIFFNESS =',F12.2//5X,	163 LEARN
	X' DISPLACEMENT AND STRESS ANALYSIS =',F12.2//5X,	164 LEARN
	X' T O T A L S O L U T I O N T I M E =',F12.2)	165 LEARN
C		166 LEARN
	RETURN	167 LEARN
	END	168 LEARN
		169 LEARN
C		
	BLOCK DATA	1 BLOCK
C		2 BLOCK
C	C.... PROGRAM TO DEFINE NODAL COORDINATES AND DEGREE OF FREEDOM OUTPUT	3 BLOCK
C	LABELS	4 BLOCK
C		5 BLOCK
C	LABELC(3) = COORDINATES LABELS	6 BLOCK
C	LABELD(6) = DEGREE OF FREEDOM LABELS	7 BLOCK
C		8 BLOCK
	COMMON/LABELS/LABELC(3),LABELD(6)	9 BLOCK
	DATA LABELC/' X1 ',' X2 ',' X3 ' /	10 BLOCK
X	LABELD/4HDOF1,4HDOF2,4HDOF3,4HDOF4,4HDOF5,4HDOF6/	11 BLOCK
	END	12 BLOCK
		13 BLOCK
C		
C	SUBROUTINE ADDSTF(A,B,S,P,IDIAG,LM,NEE)	1 ADDSTF
C		2 ADDSTF
C	C.... PROGRAM TO ADD ELEMENT STIFFNESS AND FORCE TO GLOBAL ARRAYS	3 ADDSTF
C		4 ADDSTF
	IMPLICIT REAL*8(A-H,O-Z)	5 ADDSTF
C		6 ADDSTF
C	C.... REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	7 ADDSTF
C		8 ADDSTF
	DIMENSION A(1),B(1),S(NEE,1),P(1),IDIAG(1),LM(1)	9 ADDSTF
	DO 10 J=1,NEE	10 ADDSTF
	K=LM(J)	11 ADDSTF
		12 ADDSTF

IF(K.EQ.0) GOTO 10	13 ADDSTF
B(K)=B(K)+P(J)	14 ADDSTF
L=ID(AG(K))-K	15 ADDSTF
DO 20 I=1,NEE	16 ADDSTF
M=LM(I)	17 ADDSTF
IF(M.GT.K.OR.M.EQ.0) GOTO 20	18 ADDSTF
M=L+M	19 ADDSTF
A(M)=A(M)+S(I,J)	20 ADDSTF
20 CONTINUE	21 ADDSTF
10 CONTINUE	22 ADDSTF
C	23 ADDSTF
RETURN	24 ADDSTF
END	25 ADDSTF

C	1 BC
C	2 BC
C	3 BC
C.... PROGRAM FOR READING AND PRINTING BOUNDARY CONDITION DATA	4 BC
C AND ESTABLISHING EQUATION NUMBERS	5 BC
C	6 BC
C ID(NDOF,NUMNP) = BOUNCARY CONDITION CODES	7 BC
C EQ.0, FREE	8 BC
C EQ.1, DISPLACEMENT	9 BC
C	10 BC
C COMMON/LABELS/LABELC(3),LABELD(6)	11 BC
C COMMON/DEVICE/IEL,IIN,IOUT	12 BC
C DIMENSION ID(NDOF,1)	13 BC
C	14 BC
C CALL ICLEAR(ID,NDOF*NUMNP) <i>add 30.0</i>	15 BC
C WRITE(IOUT,2000) (LABELD(I),I=1,NDOF)	16 BC
C CALL IGEN(ID,NDOF) <i>generate model bc, order</i>	17 BC
C NN=0	18 BC
C DO 10 N=1,NUMNP	19 BC
C IF(NN.LE.50) GOTO 20 <i>prints 50 lines / page of data</i>	20 BC
C NN=1	21 BC
C WRITE(IOUT,2000) (LABELD(I),I=1,NDOF)	22 BC
20 DO 30 I=1,NDOF	23 BC
C IF(ID(I,N)) 40,30,40	24 BC
30 CONTINUE	25 BC
C GOTO 10	26 BC
40 NN=NN+1	27 BC
C WRITE(IOUT,2010) N,(ID(I,N),I=1,NDOF)	28 BC
10 CONTINUE	29 BC
C	30 BC
C.... ESTABLISH EQUATION NUMBERS	31 BC
C	32 BC
C NEQ=0	33 BC
C DO 100 N=1,NUMNP	34 BC
C DO 100 I=1,NDOF	35 BC
C IF(ID(I,N).EQ.0) GOTO 110	36 BC
C ID(I,N)=0	37 BC
C GOTO 100	38 BC
110 NEQ=NEQ+1	39 BC
C ID(I,N)=NEQ	40 BC
100 CONTINUE	41 BC
C RETURN	42 BC
C	43 BC
2000 FORMAT(1H1,' N O D A L B O U N D A R Y C O N D I T I O N C O	44 BC
X D E S'///	45 BC
X 4X,' N U D E N U M B E R ',4X,6(6X,A4)/)	46 BC
2010 FORMAT(6X,I5,10X,6(5X,I5))	47 BC
C END	48 BC

C	1 COLHT
C	2 COLHT
C.... PROGRAM TO COMPUTE COLUMN HEIGHTS	3 COLHT
C	4 COLHT
C	5 COLHT
C	6 COLHT

C	DO 10 K=1,NUMEL	7 COLHT
	MIN=100000	8 COLHT
	DO 20 I=1,NED	9 COLHT
	DO 20 J=1,NEN	10 COLHT
	IF(LM(I,J,K).EQ.0) GOTU 20	11 COLHT
	MIN=MINO(MIN,LM(I,J,K))	12 COLHT
20	CONTINUE	13 COLHT
	DO 30 I=1,NED	14 COLHT
	DO 30 J=1,NEN	15 COLHT
	II=LM(I,J,K)	16 COLHT
	IF(II.EQ.0) GOTU 30	17 COLHT
	M=II-MIN	18 COLHT
	IF(M.GT.IDIAG(II)) IDIAG(II)=M	19 COLHT
30	CONTINUE	20 COLHT
10	CONTINUE	21 COLHT
C	RETURN	22 COLHT
	END	23 COLHT
		24 COLHT
		25 COLHT

for each element find min global eqn no.

this then finds the length of the column

this assigns to the IDIAG(I) the length of column I

C	SUBROUTINE CLEAR(A,M)	1 CLEAR
C		2 CLEAR
C....	PROGRAM TO CLEAR A FLOATING POINT ARRAY	3 CLEAR
C		4 CLEAR
	IMPLICIT REAL*8(A-H,O-Z)	5 CLEAR
C		6 CLEAR
C....	REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION	7 CLEAR
C		8 CLEAR
	DIMENSION A(1)	9 CLEAR
	DO 10 I=1,M	10 CLEAR
10	A(I)=0.D0	11 CLEAR
C		12 CLEAR
	RETURN	13 CLEAR
	END	14 CLEAR
		15 CLEAR

C	SUBROUTINE CCOORD(X,NSD,NUMNP)	1 CCOORD
C		2 CCOORD
C....	PROGRAM FOR READING AND PRINTING COORDINATE DATA	3 CCOORD
C		4 CCOORD
	X(NSD,NUMNP) = COORDINATE ARRAY	5 CCOORD
C		6 CCOORD
	IMPLICIT REAL*8(A-H,O-Z)	7 CCOORD
C		8 CCOORD
C....	REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	9 CCOORD
C		10 CCOORD
	COMMON/LABELS/LABELC(3),LABELD(6)	11 CCOORD
	COMMON/DEVICE/IEL,IIN,IOUT	12 CCOORD
	DIMENSION X(NSD,1)	13 CCOORD
C		14 CCOORD
	WRITE(IOUT,2000) (LABELC(I),I=1,NSD)	15 CCOORD
	CALL FGEN(X,NSD)	16 CCOORD
	NN=0	17 CCOORD
	DO 10 N=1,NUMNP	18 CCOORD
	NN=NN+1	19 CCOORD
	IF(NN.LE.50) GOTU 10	20 CCOORD
	NN=1	21 CCOORD
	WRITE(IOUT,2000) (LABELC(I),I=1,NSD)	22 CCOORD
10	WRITE(IOUT,2010) N,(X(I,N),I=1,NSD)	23 CCOORD
	RETURN	24 CCOORD
C		25 CCOORD
	2000 FORMAT(1H1,' N O D A L C O O R D I N A T E D A T A '///	26 CCOORD
	X 4X,' NODE NUMBER ',10X,3(A4,13X)//)	27 CCOORD
	2010 FORMAT(6X,15,10X,3(1PE15.8,2X))	28 CCOORD
	END	29 CCOORD
		30 CCOORD

OK as is, don't change!

generate X(NSD, NUMNP)

print 50 lines / page of data.

C	SUBROUTINE DIAG(IDIAG,NEQ,NA) ✓	1 DIAG
C		2 DIAG
C....	PROGRAM TO COMPUTE DIAGONAL ADDRESSES	3 DIAG
C	DIMENSION IDIAG(1)	4 DIAG
C	NA=1	5 DIAG
	IDIAG(1)=1	6 DIAG
	IF(NEQ.EQ.1) RETURN	7 DIAG
	DO 10 I=2,NEQ	8 DIAG
10	IDIAG(I)=IDIAG(I)+IDIAG(I-1)+1	9 DIAG
	NA=IDIAG(NEQ)	10 DIAG
C	RETURN	11 DIAG
	END	12 DIAG
		13 DIAG
		14 DIAG
		15 DIAG
		16 DIAG
C	SUBROUTINE DISBC(DL,P,S,NEE)	1 DISBC
C		2 DISBC
C....	PROGRAM TO ADJUST LOAD VECTOR FOR PRESCRIBED DISPLACEMENT	3 DISBC
C	BOUNDARY CONDITIONS	4 DISBC
C	IMPLICIT REAL*8 (A-H,O-Z)	5 DISBC
C		6 DISBC
C....	REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	7 DISBC
C	DIMENSION DL(NEE),P(NEE),S(NEE,1)	8 DISBC
C		9 DISBC
	DO 20 J=1,NEE	10 DISBC
	IF(DL(J).EQ.0.) GOTO 20	11 DISBC
	DO 10 I=1,NEE	12 DISBC
10	P(I)=P(I)-S(I,J)*DL(J)	13 DISBC
20	CONTINUE	14 DISBC
C	RETURN	15 DISBC
	END	16 DISBC
		17 DISBC
		18 DISBC
		19 DISBC
		20 DISBC
C	SUBROUTINE DISPTR(B,D,ID,NDOF,NUMNP)	1 DISPTR
C		2 DISPTR
C....	PROGRAM TO TRANSFER COMPUTED NODAL DISPLACEMENTS INTO D ARRAY	3 DISPTR
C	IMPLICIT REAL*8(A-H,O-Z)	4 DISPTR
C		5 DISPTR
C....	REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	6 DISPTR
C	DIMENSION B(1),J(NDOF,1),ID(NDOF,1)	7 DISPTR
C		8 DISPTR
	DO 10 I=1,NDOF	9 DISPTR
	DO 10 J=1,NUMNP	10 DISPTR
	K=ID(I,J)	11 DISPTR
	IF(K.GT.0) D(I,J)=B(K)	12 DISPTR
10	CONTINUE	13 DISPTR
C	RETURN	14 DISPTR
	END	15 DISPTR
		16 DISPTR
		17 DISPTR
		18 DISPTR
		19 DISPTR
C	FUNCTION DOT(A,B,N)	1 DOT
C		2 DOT
C....	PROGRAM TO PERFORM THE DOT PRODUCT OF TWO VECTORS	3 DOT
C	IMPLICIT REAL*8(A-H,O-Z)	4 DOT
C		5 DOT
C....	REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	6 DOT
		7 DOT
		8 DOT

C	DIMENSION A(1),B(1)	9 DOT
C	DOT=0.00	10 DOT
	DO 10 I=1,N	11 DOT
10	DOT=DOT+A(I)*B(I)	12 DOT
C	RETURN	13 DOT
	END	14 DOT
		15 DOT
		16 DOT
		17 DOT
C	SUBROUTINE EL	1 EL
C		2 EL
C....	PROGRAM TO INPUT AND WRITE ELEMENT GROUP DATA ON DISK	3 EL
C		4 EL
	COMMON/DEVICE/ IEL,IIN,IOUT	5 EL
	COMMON/ELDATA/ IE ,IPREC,LENGTH,MAX,MTCT,NEG,NF,NL,NPAR(16),NUMEG	6 EL
	COMMON A(1)	7 EL
C		8 EL
	REWIND IEL	9 EL
	WRITE(IOUT,2000)	10 EL
	DC 10 NEG=1,NUMEG	11 EL
	IF(NEG.GT.1) WRITE(IOUT,2010)	12 EL
	WRITE(IOUT,2020) NEG	13 EL
	READ(IIN,1000) NPAR	14 EL
	CALL ELMLIB	15 EL
	LENGTH=NL-NF+1	16 EL
	IF(LENGTH.GT.MAX) MAX=LENGTH	17 EL
10	WRITE(1EL) LENGTH,NPAR, (A(I),I=NF,NL)	18 EL
	RETURN	19 EL
C		20 EL
	1000 FORMAT(10I5)	21 EL
	2000 FORMAT(1H1,' ELEMENT GROUP DATA '///)	22 EL
	2010 FORMAT(1H1)	23 EL
	2020 FORMAT(1H0,' ELEMENT GROUP NUMBER . . . (NEG) = ',15//)	24 EL
	END	25 EL
		26 EL
C	SUBROUTINE ELGEN (IEN,MAT,NEN)	1 ELGEN
C		2 ELGEN
C....	PROGRAM TO READ AND GENERATE ELEMENT NODAL DATA AND MATERIAL	3 ELGEN
C	SET NUMBERS	4 ELGEN
C		5 ELGEN
C	IEN(NEN,NUMEL) = ELEMENT NODE NUMBERS	6 ELGEN
C	MAT(NUMEL) = ELEMENT MATERIAL SET	7 ELGEN
C	NEN = NUMBER OF NODES PER ELEMENT	8 ELGEN
C	N = ELEMENT NUMBER	9 ELGEN
C	NG = GENERATION PARAMETER	10 ELGEN
C		11 ELGEN
	COMMON/DEVICE/ IEL,IIN,IOUT	12 ELGEN
	DIMENSION IEN(NEN,1),MAT(1)	13 ELGEN
	DIMENSION IA(32)	14 ELGEN
C		15 ELGEN
	L=0	16 ELGEN
	LG=0	17 ELGEN
10	READ(IIN,1000) N,MAT1,(IA(I),I=1,NEN),NG	18 ELGEN
	IF(N.EQ.0) RETURN	19 ELGEN
	MAT(N)=MAT1	20 ELGEN
	DO 30 I=1,NEN	21 ELGEN
30	IEN(I,N)=IA(I)	22 ELGEN
	IF(LG.EQ.0) GO TO 100	23 ELGEN
C		24 ELGEN
C....	GENERATE DATA	25 ELGEN
C		26 ELGEN
	N1=L+1	27 ELGEN
	N2=N-1	28 ELGEN
	DL 20 J=N1,N2	29 ELGEN
	MAT(J)=MAT(L)	30 ELGEN
	JJ=J-1	31 ELGEN
		32 ELGEN

```

      DO 20 I=1,NEN
      20 IEN(I,J)=IEN(I,JJ)+LG
C
      100 L=N
      LG=NG
      GO TO 10
C
      1000 FORMAT(16I5)
      END

```

```

33 ELGEN
34 ELGEN
35 ELGEN
36 ELGEN
37 ELGEN
38 ELGEN
39 ELGEN
40 ELGEN
41 ELGEN

```

```

C
      SUBROUTINE ELMLIB ✓
C
C..... PROGRAM TO CALL THE ELEMENT ROUTINES
C
      COMMON/ELDATA/IE ,IPREC,LENGTH,MAX,MTCT,NEG,NF,NL,NPAR(16),NUMEG
C
      I=NPAP(1)
      GOTO(1,2,3
1      CALL TRUSS
      RETURN
2      CALL BEAM
      RETURN
3      CALL PLANE
C
      RETURN
      END

```

```

1 ELMLIB
2 ELMLIB
3 ELMLIB
4 ELMLIB
5 ELMLIB
6 ELMLIB
7 ELMLIB
8 ELMLIB
9 ELMLIB
10 ELMLIB
11 ELMLIB
12 ELMLIB
13 ELMLIB
14 ELMLIB
15 ELMLIB
16 ELMLIB
17 ELMLIB

```

```

C
      SUBROUTINE ELREAD(E) ✓
C
C..... PROGRAM TO READ ELEMENT GROUP DATA FROM AUXILIARY STORAGE DEVICE
C
      NG. IEL
C
      COMMON/DEVICE/ IEL,IIN,IOUT
      COMMON/ELDATA/IE ,IPREC,LENGTH,MAX,MTCT,NEG,NF,NL,NPAR(16),NUMEG
      DIMENSION E(1)
C
      REWIND IEL
C
      DO 10 NEG=1,NUMEG
      READ(IEL) J,NPAR,(E(I),I=1,J)
10 CALL ELMLIB
C
      RETURN
      END

```

```

1 ELREAD
2 ELREAD
3 ELREAD
4 ELREAD
5 ELREAD
6 ELREAD
7 ELREAD
8 ELREAD
9 ELREAD
10 ELREAD
11 ELREAD
12 ELREAD
13 ELREAD
14 ELREAD
15 ELREAD
16 ELREAD
17 ELREAD
18 ELREAD

```

```

C
      SUBROUTINE ERROR (I,N) ✓
C
C..... PROGRAM TO PRINT ERROR MESSAGES WHEN HIGH-SPEED STORAGE IS EXCEEDED
C
      COMMON/DEVICE/ IEL,IIN,IOUT
C
      WRITE(IOUT,2000) I
      GOTO (1,2,3,4,5
1 WRITE(IOUT,2010)
      STOP
2 WRITE(IOUT,2020)
      STOP
3 WRITE(IOUT,2030)
      STOP
4 WRITE(IOUT,2040)
      STOP
5 WRITE(IOUT,2050)
      STOP

```

```

1 ERROR
2 ERROR
3 ERROR
4 ERROR
5 ERROR
6 ERROR
7 ERROR
8 ERROR
9 ERROR
10 ERROR
11 ERROR
12 ERROR
13 ERROR
14 ERROR
15 ERROR
16 ERROR
17 ERROR
18 ERROR
19 ERROR

```


C	2000 FORMAT(// ' AVAILABLE STORAGE INSUFFICIENT BY ',I10, X' SINGLE-PRECISION WORDS')	20 ERROR
	2010 FORMAT(// ' FOR INPUT OF NODAL COORDINATE AND BOUNDARY DATA')	21 ERROR
	2020 FORMAT(// ' FOR INPUT OF APPLIED FORCE AND PRESCRIBED DISPLACEMENT D XATA')	22 ERROR
	2030 FORMAT(// ' FOR STORAGE IDIAG ARRAY')	23 ERROR
	2040 FORMAT(// ' FOR INPUT OF ELEMENT DATA')	24 ERROR
	2050 FORMAT(// ' FOR ASSEMBLY OF STIFFNESS AND LOAD')	25 ERROR
C	RETURN	26 ERROR
	END	27 ERROR
		28 ERROR
		29 ERROR
		30 ERROR
		31 ERROR

C	SUBROUTINE FGEN(A,M)	1 FGEN
C		2 FGEN
C	 PROGRAM TO READ AND GENERATE FLOATING POINT NODAL DATA	3 FGEN
C		4 FGEN
C	A = INPUT ARRAY	5 FGEN
C	M = NUMBER OF ROWS IN A	6 FGEN
C	N = NODE NUMBER	7 FGEN
C	NG = GENERATION INCREMENT	8 FGEN
C		9 FGEN
	IMPLICIT REAL*8(A-H,O-Z)	10 FGEN
C		11 FGEN
C	 REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	12 FGEN
C		13 FGEN
	COMMON /DEVICE/ IEL,IIN,IOUT	14 FGEN
	DIMENSION A(M,1)	15 FGEN
	DIMENSION F(7)	16 FGEN
C		17 FGEN
	L=0	18 FGEN
	LG=0	19 FGEN
	10 READ(IIN,1000) N,NG,(F(I),I=1,M)	20 FGEN
	IF(N.EQ.0) RETURN	21 FGEN
	DO 30 I=1,M	22 FGEN
	30 A(I,N)=F(I)	23 FGEN
	IF(LG.EQ.0) GOTO 100	24 FGEN
C		25 FGEN
C	 GENERATE DATA	26 FGEN
C		27 FGEN
	NI=(N-L)/LG	28 FGEN
	N1=L+LG	29 FGEN
	N2=N-LG	30 FGEN
	DO 20 I=1,M	31 FGEN
	DA=(A(I,N)-A(I,L))/NI	32 FGEN
	DO 20 J=N1,N2,LG	33 FGEN
	JJ=J-LG	34 FGEN
	20 A(I,J)=A(I,JJ)+DA	35 FGEN
C		36 FGEN
	100 L=N	37 FGEN
	LG=NG	38 FGEN
	GO TO 10	39 FGEN
C		40 FGEN
	1000 FORMAT(2I5,7(F10.0))	41 FGEN
	END	42 FGEN
		43 FGEN

C	SUBROUTINE FORCE(F,NDOF,NUMNP)	1 FORCE
C		2 FORCE
C	 PROGRAM FOR READING AND PRINTING APPLIED FORCE AND PRESCRIBED DISPLACEMENT DATA	3 FORCE
C		4 FORCE
C	F(NDOF,NUMNP) = APPLIED FORCES AND PRESCRIBED DISPLACEMENTS	5 FORCE
C		6 FORCE
	IMPLICIT REAL*8(A-H,O-Z)	7 FORCE
C		8 FORCE
C	 REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	9 FORCE
C		10 FORCE
		11 FORCE

COMMON/LABELS/LABELC(3),LABELD(6)	12 FORCE
COMMON/DEVICE/IEL,IIN,IOUT	13 FORCE
DIMENSION F(NDOF,1)	14 FORCE
CALL CLEAR(F,NDOF*NUMNP)	15 FORCE
C	16 FORCE
WRITE(IOUT,2000) (LABELD(I),I=1,NDOF)	17 FORCE
CALL FGEN(F,NDOF)	18 FORCE
NN=0	19 FORCE
DO 10 N=1,NUMNP	20 FORCE
IF(NN.LE.50) GOTO 20	21 FORCE
NN=1	22 FORCE
WRITE(IOUT,2000) (LABELD(I),I=1,NDOF)	23 FORCE
20 DO 30 I=1,NDOF	24 FORCE
IF(F(I,N)) 40,30,40	25 FORCE
30 CONTINUE	26 FORCE
GLTC 10	27 FORCE
40 NN=NN+1	28 FORCE
WRITE(IOUT,2010) N, (F(I,N),I=1,NDOF)	29 FORCE
10 CONTINUE	30 FORCE
RETURN	31 FORCE
C	32 FORCE
2000 FORMAT(1H1,' NODAL FORCES AND DISPLACEMENTS '///	33 FORCE
X 4X,' NODE NUMBER',10X,6(A4,13X)/)	34 FORCE
2010 FORMAT(6X,15,10X,6(1PE15.8,2X))	35 FORCE
END	36 FORCE
	37 FORCE
C	
SUBROUTINE ICLEAR(IA,M)	1 ICLEAR
C	2 ICLEAR
C.... PROGRAM TO CLEAR AN INTERGER ARRAY	3 ICLEAR
C	4 ICLEAR
DIMENSION IA(1)	5 ICLEAR
DO 10 I=1,M	6 ICLEAR
10 IA(I)=0	7 ICLEAR
C	8 ICLEAR
RETURN	9 ICLEAR
END	10 ICLEAR
	11 ICLEAR
C	
SUBROUTINE IGEN(IA,M)	1 IGEN
C	2 IGEN
C.... PROGRAM TO READ AND GENERATE INTEGER NODAL DATA	3 IGEN
C	4 IGEN
IA = INPUT ARRAY	5 IGEN
M = NUMBER OF ROWS IN IA	6 IGEN
N = NODE NUMBER	7 IGEN
NG = GENERATION INCREMENT	8 IGEN
C	9 IGEN
COMMON/DEVICE/IEL,IIN,IOUT	10 IGEN
DIMENSION IA(M,1)	11 IGEN
DIMENSION IB(32)	12 IGEN
C	13 IGEN
L=0	14 IGEN
LG=0	15 IGEN
10 READ(IIN,1000) N,NG, (IB(I),I=1,M)	16 IGEN
IF(N.EQ.0) RETURN	17 IGEN
DO 30 I=1,M	18 IGEN
30 IA(I,N)=IB(I)	19 IGEN
IF(LG.EQ.0) GOTO 100	20 IGEN
C	21 IGEN
C.... GENERATE DATA	22 IGEN
C	23 IGEN
N1=L+LG	24 IGEN
N2=N-MOD(N-L,LG)	25 IGEN
IF(N2.EQ.N) N2=N-LG	26 IGEN
DO 20 I=1,M	27 IGEN
DO 20 J=N1,N2,LG	28 IGEN
20 IA(I,J)=IA(I,L)	29 IGEN
	30 IGEN

C	100 L=N	31 IGEN
	LG=NG	32 IGEN
	GC TO 10	33 IGEN
C		34 IGEN
	1000 FORMAT(16I5)	35 IGEN
	END	36 IGEN
		37 IGEN
C		
	SUBROUTINE LCCORD(IEN,X,XL,NEC,NEN,NSD)	1 LCOORD
C		2 LCOORD
C....	PROGRAM TO LOCALIZE COORDINATE DATA	3 LCOORD
C		4 LCOORD
C	XL(NSD,NEN) = ELEMENT NODAL COORDINATES	5 LCOORD
C		6 LCOORD
	IMPLICIT REAL*8 (A-H,C-Z)	7 LCOORD
C		8 LCOORD
C....	REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	9 LCOORD
C		10 LCOORD
	DIMENSION IEN(1),X(NSD,1),XL(NSD,1)	11 LCOORD
C		12 LCOORD
	CALL CLEAR(XL,NEC)	13 LCOORD
	DO 10 J=1,NEN	14 LCOORD
	K=IEN(J)	15 LCOORD
	DO 10 I=1,NSD	16 LCOORD
10	XL(I,J)=X(I,K)	17 LCOORD
C		18 LCOORD
	RETURN	19 LCOORD
	END	20 LCOORD
		21 LCOORD
C		
	SUBROUTINE LDISBC(DL,F,IEN,LM,NDOF,NEE,NEU,NEN)	1 LDISBC
C		2 LDISBC
C....	PROGRAM TO LOCALIZE PRESCRIBED DISPLACEMENT BOUNDARY CONDITIONS	3 LDISBC
C		4 LDISBC
C	DL(NED,NEN)=ELEMENT PRESCRIBED DISPLACEMENT BOUNDARY CONDITION	5 LDISBC
C	ARRAY	6 LDISBC
C		7 LDISBC
	IMPLICIT REAL *8 (A-H,C-Z)	8 LDISBC
C		9 LDISBC
C....	REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	10 LDISBC
C		11 LDISBC
	DIMENSION DL(NED,1),F(NDOF,1),IEN(1),LM(NED,1)	12 LDISBC
C		13 LDISBC
	CALL CLEAR(DL,NEE)	14 LDISBC
C		15 LDISBC
	DO 10 I=1,NED	16 LDISBC
	DO 10 J=1,NEN	17 LDISBC
	IF(LM(I,J).GT.0) GOTO 10	18 LDISBC
	K=IEN(J)	19 LDISBC
	DL(I,J)=F(I,K)	20 LDISBC
10	CONTINUE	21 LDISBC
C		22 LDISBC
	RETURN	23 LDISBC
	END	24 LDISBC
		25 LDISBC
C		
	SUBROUTINE LDISP(DL,D,IEN,NDOF,NED,NEN)	1 LDISP
C		2 LDISP
C....	PROGRAM TO LOCALIZE DISPLACEMENTS	3 LDISP
C		4 LDISP
C	DL(NED,NEN) = ELEMENT DISPLACEMENT ARRAY	5 LDISP
C		6 LDISP
	IMPLICIT REAL*8(A-H,C-Z)	7 LDISP
C		8 LDISP
C....	REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	9 LDISP
		10 LDISP

if global eqn no. $\neq 0$ the node is free & value is that of free

if global eqn no. $= 0$ the node is fixed & value is that of node

find global node no. $K = IEN(J)$

transfer to local displ vector $DL(I)$ with $F(K)$

C	DIMENSION IEN(1),DL(NED,1),D(NDOF,1)	11 LDISP
C	DO 10 J=1,NEN	12 LDISP
	K=IEN(J)	13 LDISP
	DO 10 I=1,NED	14 LDISP
	10 DL(I,J)=D(I,K)	15 LDISP
C	RETURN	16 LDISP
	END	17 LDISP
		18 LDISP
		19 LDISP
		20 LDISP
C	SUBROUTINE LCAU(ID,F,B,NDOF,NUMNP)	1 LOAD
C	C.... PROGRAM TO TRANSFER NODAL APPLIED FORCES INTO RIGHT HAND SIDE VECTOR	2 LOAD
C	B(NEQ) = RIGHT-HAND SIDE VECTOR	3 LOAD
C	IMPLICIT REAL*8 (A-H,C-Z)	4 LOAD
C	C.... REMOVE ABOVE CARD FOR SINGLE-PRECISION OPERATION	5 LOAD
C	DIMENSION ID(NDOF,1),F(NDOF,1),B(1)	6 LOAD
C	DO 10 I=1,NDOF	7 LOAD
	DO 10 J=1,NUMNP	8 LOAD
	K=ID(I,J)	9 LOAD
	10 IF(K.GT.0) B(K)=F(I,J)	10 LOAD
C	RETURN	11 LOAD
	END	12 LOAD
		13 LOAD
		14 LOAD
		15 LOAD
		16 LCAU
		17 LCAU
		18 LCAU
		19 LCAU
		20 LCAU
C	SUBROUTINE LOCAL(ID,IEN,LM,NDOF,NED,NEN,NUMEL)	1 LOCAL
C	C.... SUBROUTINE TO LOCALIZE ID ARRAY	2 LOCAL
C	DIMENSION ID(NDOF,1),IEN(NEN,1),LM(NED,NEN,1)	3 LOCAL
C	DO 10 K=1,NUMEL	4 LOCAL
	DO 10 J=1,NEN	5 LOCAL
	NA=IEN(J,K)	6 LOCAL
	DO 10 I=1,NED	7 LOCAL
	10 LM(I,J,K)=ID(I,NA)	8 LOCAL
C	RETURN	9 LOCAL
	END	10 LOCAL
		11 LOCAL
		12 LOCAL
		13 LOCAL
		14 LOCAL
		15 LOCAL
C	SUBROUTINE PACK(S,NP,NU)	1 PACK
C	C.... PROGRAM TO PACK AN NU*NU MATRIX INTO AN NP*NP MATRIX	2 PACK
C	IMPLICIT REAL*8(A-H,U-Z)	3 PACK
C	C.... REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION	4 PACK
C	DIMENSION S(1)	5 PACK
	DO 10 J=1,NP	6 PACK
	K=(J-1)*NP	7 PACK
	L=(J-1)*NU	8 PACK
	DO 10 I=1,NP	9 PACK
	K=K+1	10 PACK
	L=L+1	11 PACK
	10 S(K)=S(L)	12 PACK
C	RETURN	13 PACK
	END	14 PACK
		15 PACK
		16 PACK
		17 PACK
		18 PACK
		19 PACK
		20 PACK

C	SUBROUTINE PREAD(IA,JA,A,M)	1 PREAD
C		2 PREAD
C....	PROGRAM TO READ ELEMENT CONSISTENT LOAD DATA	3 PREAD
C		4 PREAD
C	IMPLICIT REAL*8(A-H,O-Z)	5 PREAD
C		6 PREAD
C....	REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION	7 PREAD
C		8 PREAD
C	A = INPUT ARRAY	9 PREAD
C	IA = ELEMENT NUMBER	10 PREAD
C	JA = ELEMENT FACE/SIDE NUMBER	11 PREAD
C	M = NUMBER OF LOADS PER FACE/SIDE	12 PREAD
C	N = LCAD NUMBER	13 PREAD
C		14 PREAD
C	COMMON/DEVICE/ IEL,IIN,IOUT	15 PREAD
C	DIMENSION A(M,1),IA(1),JA(1),B(6)	16 PREAD
C		17 PREAD
C	N=0	18 PREAD
10	N=N+1	19 PREAD
	READ (IIN,1000) II,JJ,(B(I),I=1,M)	20 PREAD
	IF (II.EQ.0) RETURN	21 PREAD
	IA(N)=II	22 PREAD
	JA(N)=JJ	23 PREAD
	DO 20 I=1,M	24 PREAD
20	A(I,N)=B(I)	25 PREAD
	GO TO 10	26 PREAD
C		27 PREAD
	1000 FORMAT(2I5,10X,6F10.0)	28 PREAD
	END	29 PREAD
		30 PREAD
C		
C	SUBROUTINE PRINC(N,S,P)	1 PRINC
C		2 PRINC
C....	PROGRAM TO COMPUTE PRINCIPAL VALUES OF SYMMETRIC SECOND RANK TENSOR	3 PRINC
C		4 PRINC
C	S = SYMMETRIC SECOND-RANK TENSOR STORED AS A VECTOR	5 PRINC
C	N = NUMBER OF DIMENSIONS (2 OR 3)	6 PRINC
C	P = PRINCIPAL VALUES	7 PRINC
C		8 PRINC
C....	THE COMPONENTS OF S MUST BE STORED IN THE FOLLOWING ORDERS	9 PRINC
C		10 PRINC
C	2-D PROBLEMS, S11,S12,S22	11 PRINC
C	3-D PROBLEMS, S11,S12,S13,S22,S23,S33	12 PRINC
C		13 PRINC
C	IMPLICIT REAL*8(A-H,O-Z)	14 PRINC
C		15 PRINC
C....	REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION	16 PRINC
C		17 PRINC
C	DIMENSION S(1),P(1)	18 PRINC
	DATA RT2/1.41421356237309505,PI23/2.0943951023932100/	19 PRINC
C		20 PRINC
	IF(N.EQ.3) GLTC 10	21 PRINC
C		22 PRINC
C....	2-D PROGRAM	23 PRINC
C		24 PRINC
	A=22.500/DATAN(1.00)	25 PRINC
	X=0.500*(S(1)+S(3))	26 PRINC
	Y=0.500*(S(1)-S(3))	27 PRINC
	R=DSQRT(Y*Y+S(2)*S(2))	28 PRINC
	P(1)=X+R	29 PRINC
	P(2)=X-R	30 PRINC
	P(3)=45.00	31 PRINC
	IF(Y.NE.0.00.OR.S(2).NE.0.00) P(3)=A*DATAN2(S(2),Y)	32 PRINC
	IF (S(2).EQ.0.00) P(3)=0.00	33 PRINC
	P(4)=R	34 PRINC
	RETURN	35 PRINC
C		36 PRINC
C....	3-D PROBLEM	37 PRINC
C		38 PRINC
		39 PRINC

10 R=0.00	40 PRINC
X=(S(1)+S(4)+S(6))/3.00	41 PRINC
Y=S(1)*(S(4)+S(6))+S(4)*S(6)-S(2)*S(2)-S(3)*S(3)-S(5)*S(5)	42 PRINC
Z=S(1)*S(4)*S(6)+2.00*S(2)*S(3)*S(5)-S(1)*S(5)*S(5)	43 PRINC
X -S(4)*S(3)*S(3)-S(6)*S(2)*S(2)	44 PRINC
T=3.00*X*X-Y	45 PRINC
U=0.00	46 PRINC
IF(T.EQ.0.00) GO TO 20	47 PRINC
U=DSQRT(2.00*T/3.00)	48 PRINC
A=(Z+(T-X*X)*X)*RT2/U**3	49 PRINC
R=DSQRT(DABS(1.00-A*A))	50 PRINC
R=ATAN2(R,A)/3.00	51 PRINC
20 P(1)=X+U*RT2*DCOS(R)	52 PRINC
P(2)=X+U*RT2*DCOS(R-PI/3)	53 PRINC
P(3)=X+U*RT2*DCOS(R+PI/3)	54 PRINC
C	55 PRINC
RETURN	56 PRINC
END	57 PRINC
C	1 PRINTD
SUBROUTINE PRINTD(D,NDOF,NUMNP)	2 PRINTD
C	3 PRINTD
C.... PROGRAM TO PRINT DISPLACEMENTS	4 PRINTD
C	5 PRINTD
IMPLICIT REAL*8(A-H,O-Z)	6 PRINTD
C	7 PRINTD
C.... REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION	8 PRINTD
C	9 PRINTD
DIMENSION D(NDOF,1)	10 PRINTD
COMMON/DEVICE/IEL,IIN,IOUT	11 PRINTD
COMMON/LABELS/LABELC(3),LABELD(6)	12 PRINTD
WRITE(IOUT,2000) (LABELD(I),I=1,NDOF)	13 PRINTD
NA=0	14 PRINTD
DO 10 N=1,NUMNP	15 PRINTD
NA=NA+1	16 PRINTD
IF(NN.LE.50) GOTO 10	17 PRINTD
NN=1	18 PRINTD
WRITE(IOUT,2000) (LABELD(I),I=1,NDOF)	19 PRINTD
10 WRITE(IOUT,2010) N,(D(I,N),I=1,NDOF)	20 PRINTD
C	21 PRINTD
RETURN	22 PRINTD
2000 FORMAT(1H1,' D I S P L A C E M E N T S '///	23 PRINTD
X 4X,' NODE NUMBER',10X,6(1A4,13X)/)	24 PRINTD
2010 FORMAT(6X,15,10X,6(1PE15.8,2X))	25 PRINTD
END	26 PRINTD
C	1 PRINTP
SUBROUTINE PRINTP(A,IDIAG,NEQ)	2 PRINTP
C	3 PRINTP
C.... PROGRAM TO PRINT ARRAY D AFTER FACTORIZATION A=(U)T*D*U	4 PRINTP
C	5 PRINTP
IMPLICIT REAL*8(A-H,O-Z)	6 PRINTP
C	7 PRINTP
C.... REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION	8 PRINTP
C	9 PRINTP
COMMON/DEVICE/IEL,IIN,IOUT	10 PRINTP
DIMENSION A(1),IDIAG(1)	11 PRINTP
C	12 PRINTP
WRITE(IOUT,2000)	13 PRINTP
NN=0	14 PRINTP
DO 10 N=1,NEQ	15 PRINTP
NN=NN+1	16 PRINTP
IF(NN.LE.50) GOTO 20	17 PRINTP
WRITE(IOUT,2000)	18 PRINTP
20 I=IDIAG(N)	19 PRINTP
NN=0	20 PRINTP
D=A(I)	21 PRINTP
10 WRITE(IOUT,2010) N,D	22 PRINTP
RETURN	23 PRINTP

```

C      2000 FORMAT(1H1,' ARRAY D(NEQ) OF FACTORIZATION  A=(U)T*D*U'////)
C      2010 FORMAT(1X,15,4X,1PE20.8)
C      END
24 PRINTP
25 PRINTP
26 PRINTP
27 PRINTP
28 PRINTP

C
C      SUBROUTINE SECONDT
C      C.... PROGRAM TO DETERMINE TIME
C      C      SUBROUTINE RTIME IS USED ON CAL TECH IBM 370 SYSTEM
C      C      INTEGER RTIME
C      C      IX=RTIME(ZDUM)
C      C      T=IX/1000.D0
C      C      RETURN
C      C      END
1 SECOND
2 SECOND
3 SECOND
4 SECOND
5 SECOND
6 SECOND
7 SECOND
8 SECOND
9 SECOND
10 SECOND
11 SECOND
12 SECOND
13 SECOND

C
C      SUBROUTINE SHAP20(S,T,X,DET,SH,NEN,INC,IEN,QUAD)
C      C.... PROGRAM TO COMPUTE SHAPE FUNCTIONS FOR QUADRILATERAL
C      C      S,T      = NATURAL COORDINATES
C      C      SH(NSD,I) = DERIVATIVES OF SHAPE FUNCTIONS
C      C      SH(3,I)    = SHAPE FUNCTIONS
C      C      XS(NSD,NSD)= JACOBIAN MATRIX
C      C      DET      = JACOBIAN DETERMINANT
C      C      X(NSD,NEN) = GLOBAL COORDINATES
C      C      IMPLICIT REAL*8(A-H,U-Z)
C      C      C.... REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION
C      C      LOGICAL QUAD
C      C      COMMON /DEVICE/ IEL,IIN,ICUT
C      C      DIMENSION SA(4),TA(4),SH(3,1),X(2,1),XS(2,2)
C      C      DIMENSION IEN(1)
C      C      DATA SA/-0.500,0.500,0.500,-0.500/,TA/-0.500,-0.500,0.500,0.500/
C      C      DO 10 I=1,4
C      C      SH(3,I)=(0.500+SA(I)*S)*(0.500+TA(I)*T)
C      C      SH(1,I)=SA(I)*(0.500+TA(I)*T)
C      C      SH(2,I)=TA(I)*(0.500+SA(I)*S)
C      C      IF(QUAD) GO TO 30
C      C      DO 20 I=1,3
C      C      SH(I,3)=SH(I,3)+SH(I,4)
C      C      SH(I,4)=0.D0
C      C      IF(NEN.GT.4) CALL SHAP21(S,T,SH,NEN,IEN)
C      C      DO 40 I=1,2
C      C      DO 40 J=1,2
C      C      XS(I,J)=0.D0
C      C      DO 40 K=1,NEN
C      C      XS(I,J)=XS(I,J)+X(I,K)*SH(J,K)
C      C      DET=XS(1,1)*XS(2,2)-XS(1,2)*XS(2,1)
C      C      DO 50 I=1,2
C      C      DO 50 J=1,2
C      C      XS(I,J)=XS(I,J)/DET
C      C      TEMP=XS(2,2)*SH(1,I)-XS(2,1)*SH(2,I)
C      C      SH(2,I)=-XS(1,2)*SH(1,I)+XS(1,1)*SH(2,I)
C      C      SH(1,I)=TEMP
C      C      IF(INC.EQ.0) RETURN
C      C      C.... INCOMPATIBLE MODES
C      C

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```

      IF(QUAD) GOTO 80
      DO 70 I=1,3
      DO 70 J=5,6
70    SH(I,J)=0.00
      RETURN
80    SH(1,5)=-S-S      N5
      SH(2,5)=0.00
      SH(3,5)=1.00-S*S  N5
      SH(1,6)=0.00
      SH(2,6)=-T-T      N6
      SH(3,6)=1.00-T*T  N6
      XS(1,1)=0.2500*(-X(1,1)+X(1,2)+X(1,3)-X(1,4))
      XS(1,2)=0.2500*(-X(1,1)-X(1,2)+X(1,3)+X(1,4))
      XS(2,1)=0.2500*(-X(2,1)+X(2,2)+X(2,3)-X(2,4))
      XS(2,2)=0.2500*(-X(2,1)-X(2,2)+X(2,3)+X(2,4))
      DO 90 I=5,6
      TEMP=(XS(2,2)*SH(1,1)-XS(2,1)*SH(2,1))/DET
      SH(2,1)=(-XS(1,2)*SH(1,1)+XS(1,1)*SH(2,1))/DET
90    SH(1,1)=TEMP
      RETURN
      END

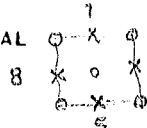
```

48 SHAP20
 49 SHAP20
 50 SHAP20
 51 SHAP20
 52 SHAP20
 53 SHAP20
 54 SHAP20
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 62 SHAP20
 63 SHAP20
 64 SHAP20
 65 SHAP20
 66 SHAP20
 67 SHAP20
 68 SHAP20

```

C      SUBROUTINE SHAP21 (S,T,SH,NEN,IEN)
C
C..... PROGRAM TO COMPUTE SHAPE FUNCTIONS FOR QUADRILATERAL
C      SHAPE FUNCTIONS 5 TO 8
C
C      IMPLICIT REAL*8(A-H,O-Z)
C
C..... REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION
C

```



```

      DIMENSION SH(3,1), IEN(1)
      SS=(1.00-S*S)/2.00
      TT=(1.00-T*T)/2.00
      DO 10 I=5,8
      DO 10 J=1,3
10    SH(J,I)=0.00
      IF(IEN(5).EQ.0) GOTO 20
      SH(1,5)=-S*(1.00-T)
      SH(2,5)=-SS
      SH(3,5)=SS*(1.00-T)
20    IF(NEN.LT.6) GOTO 50
      IF(IEN(6).EQ.0) GOTO 30
      SH(1,6)=TT
      SH(2,6)=-T*(1.00+S)
      SH(3,6)=TT*(1.00+S)
30    IF(NEN.LT.7) GOTO 50
      IF(IEN(7).EQ.0) GOTO 40
      SH(1,7)=-S*(1.00+T)
      SH(2,7)=SS
      SH(3,7)=SS*(1.00+T)
40    IF(NEN.LT.8) GOTO 50
      IF(IEN(8).EQ.0) GOTO 50
      SH(1,8)=-TT
      SH(2,8)=-T*(1.00-S)
      SH(3,8)=TT*(1.00-S)
50    K=8
      DO 70 I=1,4
      L=I+4
      DO 60 J=1,3
60    SH(J,I)=SH(J,I)-0.500*(SH(J,K)+SH(J,L))
70    K=L
      RETURN
      END

```

See 1st quadrant nodes
 addition of nodes using Lagrange

Node 5 & deriv

Node 6 & deriv

Node 7 & deriv

Node 8 & deriv

modified nodes 1-4
 to include addition of N₅ thru N₈

1 SHAP21
 2 SHAP21
 3 SHAP21
 4 SHAP21
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 41 SHAP21
 42 SHAP21
 43 SHAP21
 44 SHAP21

C	SUBROUTINE SOLVE (A,B,IDIAG,NEQ,FACT,BACK)	1 SOLVE
C		2 SOLVE
C	 PROGRAM TO PERFORM A=(U)T*D*(U FACTORIZATION AND/OR FORWARD	3 SOLVE
C	REDUCTION/BACKSUBSTITUTION OF SYMMETRIC POSITIVE DEFINITE	4 SOLVE
C	SYSTEM OF EQUATIONS	5 SOLVE
C		6 SOLVE
C	A(NA) = COEFFICIENT MATRIX STORED IN COMPACTED COLUMN FORM	7 SOLVE
C	(AFTER FACTORIZATION CONTAINS D AND J)	8 SOLVE
C	B(NEQ) = RIGHT SIDE VECTOR (AFTER BACKSUBSTITUTION SOLUTION	9 SOLVE
C	VECTOR)	10 SOLVE
C	IDIAG(NEQ) = ADDRESSES OF DIAGONAL TERMS IN A(NA)	11 SOLVE
C	NEQ = NUMBER OF EQUATIONS	12 SOLVE
C	FACT = .TRUE. , FACTOR A(NA)	13 SOLVE
C	.FALSE. , DO NOT FACTOR A(NA)	14 SOLVE
C	BACK = .TRUE. , FORWARD REDUCE B(NEQ) AND BACKSUBSTITUTE	15 SOLVE
C	.FALSE. , DO NOT FORWARD REDUCE B(NEQ) OR	16 SOLVE
C	BACKSUBSTITUTE	17 SOLVE
C		18 SOLVE
C	IMPLICIT REAL*8(A-H,G-Z)	19 SOLVE
C		20 SOLVE
C	 REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION	21 SOLVE
C		22 SOLVE
C	LOGICAL FACT,BACK	23 SOLVE
C	DIMENSION A(1),B(1),IDIAG(1)	24 SOLVE
C		25 SOLVE
C	 FACTOR A , REDUCE B	26 SOLVE
C	JR = 0	27 SOLVE
C	DO 60 J=1,NEQ	28 SOLVE
C	JD = IDIAG(J)	29 SOLVE
C	JH = JD - JR	30 SOLVE
C	IS = J - JH + 2	31 SOLVE
C	IF(JH-2) 60,30,10	32 SOLVE
C	10 IF(.NOT.FACT) GO TO 50	33 SOLVE
C	IE = J - 1	34 SOLVE
C	K = JR + 2	35 SOLVE
C	ID = IDIAG(IS - 1)	36 SOLVE
C		37 SOLVE
C	 REDUCE ALL EQUATIONS EXCEPT DIAGONAL	38 SOLVE
C		39 SOLVE
C	DO 20 I=IS,IE	40 SOLVE
C	IR = ID	41 SOLVE
C	ID = IDIAG(I)	42 SOLVE
C	IH = MIN0(ID-IR-1,I-IS+1)	43 SOLVE
C	IF(IH.GT.0) A(K) = A(K) - DOT(A(K-IH),A(ID-IH),IH)	44 SOLVE
C	20 K = K + 1	45 SOLVE
C		46 SOLVE
C	 REDUCE DIAGONAL TERM	47 SOLVE
C		48 SOLVE
C	30 IF(.NOT.FACT) GO TO 50	49 SOLVE
C	IR = JR + 1	50 SOLVE
C	IE = JD - 1	51 SOLVE
C	K = J - JD	52 SOLVE
C	DO 40 I=IR,IE	53 SOLVE
C	ID = IDIAG(K+1)	54 SOLVE
C	IF(A(ID).EQ.0.000) GO TO 40	55 SOLVE
C	D = A(I)	56 SOLVE
C	A(I) = A(I)/A(ID)	57 SOLVE
C	A(JD) = A(JD) - D*A(I)	58 SOLVE
C	40 CONTINUE	59 SOLVE
C		60 SOLVE
C	 REDUCE RHS	61 SOLVE
C		62 SOLVE
C	50 IF(BACK) B(J) = B(J) - DOT(A(JR+1),B(IS-1),JH-1)	63 SOLVE
C	60 JR = JD	64 SOLVE
C	IF(.NOT.BACK) RETURN	65 SOLVE
C		66 SOLVE
C	 DIVIDE BY DIAGONAL PIVOTS	67 SOLVE
C		68 SOLVE
C	DO 70 I=1,NEQ	69 SOLVE
C	ID = IDIAG(I)	70 SOLVE
C	IF(A(ID).NE.0.000) B(I) = B(I)/A(ID)	71 SOLVE
C	70 CONTINUE	72 SOLVE
		73 SOLVE

C		74 SOLVE
C.... BACKSUBSTITUTE		75 SOLVE
C	J = NEQ	76 SOLVE
	JD = IDIAG(J)	77 SOLVE
80	D = B(J)	78 SOLVE
	J = J - 1	79 SOLVE
	IF(J.LE.0) RETURN	80 SOLVE
	JR = IDIAG(J)	81 SOLVE
	IF(JD-JR.LE.1) GO TO 100	82 SOLVE
	IS = J - JD + JR + 2	83 SOLVE
	K = JR - IS + 1	84 SOLVE
	DO 90 I=IS,J	85 SOLVE
90	B(I) = B(I) - A(I+K)*D	86 SOLVE
100	JD = JR	87 SOLVE
	GO TO 80	88 SOLVE
C		89 SOLVE
	END	90 SOLVE
		91 SOLVE
C		
	SUBROUTINE STACCN (S,P,N,M)	1 STACCN
C		2 STACCN
C.... PROGRAM TO STATICALLY CONDENSE SYMMETRIC POSITIVE DEFINITE MATRIX		3 STACCN
C	S AND VECTOR P	4 STACCN
C		5 STACCN
C	N= ORIGINAL NUMBER ROWS/COLUMNS IN A AND VECTOR P	6 STACCN
C	M= NUMBER OF ROWS/COLUMNS TO BE CONDENSED	7 STACCN
C		8 STACCN
	IMPLICIT REAL*8(A-H,O-Z)	9 STACCN
C		10 STACCN
C.... REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION		11 STACCN
C		12 STACCN
	DIMENSION S(N,1),P(1)	13 STACCN
C		14 STACCN
	DO 10 I=1,M	15 STACCN
	II=N-I	16 STACCN
	JJ=II+1	17 STACCN
	SS=S(JJ,JJ)	18 STACCN
	DO 10 J=1,II	19 STACCN
	TT=S(J,JJ)/SS	20 STACCN
	P(J)=P(J)-TT*P(JJ)	21 STACCN
	DO 10 K=1,J	22 STACCN
	S(K,J)=S(K,J)-S(K,JJ)*TT	23 STACCN
10	S(J,K)=S(K,J)	24 STACCN
C		25 STACCN
	RETURN	26 STACCN
	END	27 STACCN
		28 STACCN
C		
	SUBROUTINE TRUSS	1 TRUSS
C		2 TRUSS
C.... PROGRAM TO SET UP STORAGE FOR TRUSS ELEMENTS		3 TRUSS
C		4 TRUSS
	COMMON/ELDATA/IE ,IPREC,LENGTH,MAX,MTCT,NEG,NF,NL,NPAR(16),NUMEG	5 TRUSS
	COMMON/POINTS/N1,N2,N3,N4,N5,N6,N7	6 TRUSS
	COMMON A(1)	7 TRUSS
C		8 TRUSS
	EQUIVALENCE (NPAR(2),NUMEL),(NPAR(3),NUMAT)	9 TRUSS
C		10 TRUSS
	IF(NUMAT.EQ.0) NUMAT=1	11 TRUSS
	NF=N5	12 TRUSS
	IF(IE.GT.1) NF=N7	13 TRUSS
	N101=NF	14 TRUSS
	N=NUMAT*IPREC	15 TRUSS
	N102=N101+N	16 TRUSS
	N103=N102+N	17 TRUSS
	N104=N103+N	18 TRUSS
	N105=N104+3*IPREC	19 TRUSS
		20 TRUSS

N106=N105+NUMEL	21 TRUSS
N107=N106+2*NUMEL	22 TRUSS
N108=N107+6*NUMEL	23 TRUSS
NL=N108	24 TRUSS
C IF((IF.EQ.1.AND.NL.GT.MTCT) CALL ERROR(NL-MTCT,4)	25 TRUSS
CALL TRUSS1(A(N1),A(N2),A(N3),A(N4),A(N5),A(N6),	26 TRUSS
XA(N101),A(N102),A(N103),A(N104),A(N105),A(N106),A(N107))	27 TRUSS
C	28 TRUSS
RETURN	29 TRUSS
END	30 TRUSS
	31 TRUSS
C	
SUBROUTINE TRUSS1 (ID,X,C,IDIAG,B,A,E,AREA,WT,GRAV,	1 TRUSS1
X MAT,IEN,LM)	2 TRUSS1
C	3 TRUSS1
C.... TRUSS ELEMENT PROGRAM	4 TRUSS1
C	5 TRUSS1
IMPLICIT REAL*8 (A-H,O-Z)	6 TRUSS1
C	7 TRUSS1
C.... REMOVE ABOVE CARD FLR SINGLE PRECISION OPERATION	8 TRUSS1
C	9 TRUSS1
COMMON/INFO/ICGDE,MODE,NDCF,NSD,NUMNP	10 TRUSS1
COMMON/ELDATA/IE ,IPREC,LENGTH,MAX,MTCT,NEG,NF,NL,NPAR(16),NUMEG	11 TRUSS1
COMMON/DEVICE/IEL,IIN,IOUT	12 TRUSS1
DIMENSION ID(NDCF,1),X(NSC,1),D(NDCF,1),IDIAG(1),B(1),A(1)	13 TRUSS1
DIMENSION E(1),AREA(1),WT(1),GRAV(3),MAT(1),IEN(2,1),LM(3,2,1)	14 TRUSS1
DIMENSION ST(6),S(6,6),P(6),DL(3,2),XL(3,2)	15 TRUSS1
EQUIVALENCE (NPAR(1),NTYPE),(NPAR(2),NUMEL),(NPAR(3),NUMAT)	16 TRUSS1
C	17 TRUSS1
NED=3	18 TRUSS1
NEN=2	19 TRUSS1
NEE=6	20 TRUSS1
NEC=6	21 TRUSS1
C	22 TRUSS1
GCTO (1,2,3	23 TRUSS1
),IE	24 TRUSS1
C	25 TRUSS1
C.... READ AND PRINT ELEMENT DATA	26 TRUSS1
C	27 TRUSS1
1 WRITE(IOUT,2000) NTYPE,NUMEL	28 TRUSS1
WRITE(IOUT,2010) NUMAT	29 TRUSS1
WRITE(IOUT,2020)	30 TRUSS1
DO 10 N=1,NUMAT	31 TRUSS1
READ(IIN,1000) N,E(N),AREA(N),WT(N)	32 TRUSS1
10 WRITE(IOUT,2030) N,E(N),AREA(N),WT(N)	33 TRUSS1
READ(IIN,1010) GRAV	34 TRUSS1
WRITE(IOUT,2040) GRAV	35 TRUSS1
C	36 TRUSS1
CALL ELGEN(IEN,MAT,NEN)	37 TRUSS1
WRITE(IOUT,2050)	38 TRUSS1
NN=0	39 TRUSS1
DO 20 N=1,NUMEL	40 TRUSS1
NN=NN+1	41 TRUSS1
IF(NN.LE.50) GCTO 20	42 TRUSS1
NN=1	43 TRUSS1
WRITE(IOUT,2050)	44 TRUSS1
20 WRITE(IOUT,2060) N,MAT(N),(IEN(I,N),I=1,NEN)	45 TRUSS1
C	46 TRUSS1
CALL LCCAL(ID,IEN,LM,NDCF,NED,NEN,NUMEL)	47 TRUSS1
CALL COLHT(IDIAG,LM,NED,NEN,NUMEL)	48 TRUSS1
C	49 TRUSS1
RETURN	50 TRUSS1
C	51 TRUSS1
C.... FORM STIFFNESS MATRIX AND CONSISTENT LOAD VECTOR	52 TRUSS1
C	53 TRUSS1
2 DO 100 N=1,NUMEL	54 TRUSS1
M=MAT(N)	55 TRUSS1
CALL LCLORD(IEN(1,N),X,XL,NEC,NEN,NSD)	56 TRUSS1

```

      Y=0.00
      DO 110 I=1,3
        ST(I)=XL(I,1)-XL(I,2)
110    Y=Y+ST(I)**2
        Y=DSQRT(Y)
        DO 120 I=1,3
          ST(I)=ST(I)/Y
120    ST(I+3)=-ST(I)
        FOR=WT(M)*Y/2.00
        DO 130 I=1,3
          P(I)=FOR*GRAV(I)
130    P(I+3)=P(I)
        Y=E(M)*AREA(M)/Y
        DO 140 I=1,6
          Z=ST(I)*Y
          DO 140 J=1,I
            S(I,J)=S(J,I)*Z
140    S(J,I)=S(I,J)
        CALL LDISBC(DL,D,IEN(1,N),LM(1,1,N),NDOF,NEE,NED,NEN)
        CALL DISBC(DL,P,S,NEE)
100    CALL ADUSTF(A,B,S,P,IDIAG,LM(1,1,N),NEE)
C
      RETURN
C
C..... STRESS COMPUTATION
C
      3  NN=0
        WRITE(IGOUT,2070) NEG
        DO 200 N=1,NUMEL
          M=MAT(N)
          CALL LCOORD(IEN(1,N),X,XL,NEC,NEN,NSD)
          Y=0.00
          DO 210 I=1,3
            ST(I)=XL(I,1)-XL(I,2)
C
210    Y=Y+ST(I)**2
          Y=DSQRT(Y)
          CALL LDISP(DL,D,IEN(1,N),NDOF,NED,NEN)
          Z=0.00
          DO 220 I=1,3
            ST(I)=ST(I)+DL(I,1)-DL(I,2)
220    Z=Z+ST(I)**2
          Z=DSQRT(Z)
          STR=E(M)*(Z-Y)/Y
          FOR=STR*AREA(M)
          NN=NN+1
          IF(NN.LE.50) GOTO 200
          NN=1
          WRITE(IGOUT,2070) NEG
200    WRITE(IGOUT,2080) N,STR,FOR
          RETURN
C
1000  FORMAT(15,5X,3F10.0)
1010  FORMAT(3F10.0)
2000  FORMAT(1H0,
        X' T R U S S   E L E M E N T S           ',//5X,
        X' ELEMENT TYPE NUMBER . . . . . (NTYPE) =',15//5X,
        X' NUMBER OF ELEMENTS . . . . . (NUMEL) =',15//)
2010  FORMAT(1H0,
        X' M A T E R I A L   S E T   D A T A           ',//5X,
        X' NUMBER OF MATERIAL SETS . . . . . (NUMAT) =',15//)
2020  FORMAT(1H0,
        X 5X,'SET',10X,'YOUNGS',10X,'CROSS-SECTION',8X,'WEIGHT/UNIT',/,
        X 4X,'NUMBER',8X,'MODULUS',14X,'AREA',16X,'LENGTH'/)
2030  FORMAT(3X,15,2X,3(1PE20.8))
2040  FORMAT(1H0,
        X' G R A V I T Y   L G A D   M U L T I P L I E R S           ',//5X,
        X' X1-DIRECTION . . . . . =',
        X1PE20.8//5X,
        X' X2-DIRECTION . . . . . =',
        X1PE20.8//5X,
        X' X3-DIRECTION . . . . . =',
        X1PE20.8//)

```

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57 TRUSS1
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129 TRUSS1

```

2050 FORMAT(1H1,	130 TRUSS1
X' E L E M E N T D A T A',//5X,	131 TRUSS1
X' ELEMENT MATERIAL NGDE 1 NODE 2'//5X,	132 TRUSS1
X' NUMBER NUMBER'//)	133 TRUSS1
2060 FORMAT(6X,15,3(5X,15))	134 TRUSS1
2070 FORMAT(1H1,	135 TRUSS1
X' T R U S S E L E M E N T S T R E S S E S ', //5X,	136 TRUSS1
X' ELEMENT GROUP NUMBER (NEG) ='//5X,	137 TRUSS1
X' ELEMENT',7X,'STRESS',9X,'FORCE',/5X,	138 TRUSS1
X' NUMBER'//)	139 TRUSS1
2080 FORMAT(6X,15,2(1PE20.8))	140 TRUSS1
END	141 TRUSS1
C	1 BEAM
SUBROUTINE BEAM	2 BEAM
C	3 BEAM
IMPLICIT REAL*8 (A-H,O-Z)	4 BEAM
C	5 BEAM
RETURN	6 BEAM
END	7 BEAM
C	1 PLANE
SUBROUTINE PLANE	2 PLANE
C	3 PLANE
C.... PROGRAM TO SET UP STORAGE FOR 4-NODE QUADRILATERAL ISOTROPIC	4 PLANE
C PLANE STRESS/STRAIN ELEMENTS	5 PLANE
C	6 PLANE
COMMON/POINTS/N1,N2,N3,N4,N5,N6,N7	7 PLANE
COMMON/ELDATA/IE ,IPREC,LENGTH,MAX,MTOT,NEG,NF,NL,NPAR(16),NUMEG	8 PLANE
COMMON A(1)	9 PLANE
C	10 PLANE
EQUIVALENCE (NPAR(2),NUMEL),(NPAR(3),NUMAT),(NPAR(9),NUMPR)	11 PLANE
C	12 PLANE
IF (NUMAT.EQ.0) NUMAT=1	13 PLANE
NF=N5	14 PLANE
IF(IE.GT.1) NF=N7	15 PLANE
N101=NF	16 PLANE
N=NUMAT*IPREC	17 PLANE
N102=N101+N	18 PLANE
N103=N102+N	19 PLANE
N104=N103+N	20 PLANE
N105=N104+N	21 PLANE
N106=N105+2*IPREC	22 PLANE
N107=N106+NUMEL	23 PLANE
N108=N107+4*NUMEL	24 PLANE
N109=N108+8*NUMEL	25 PLANE
N110=N109+NUMPR	26 PLANE
N111=N110+NUMPR	27 PLANE
NL=N111+NUMPR*IPREC	28 PLANE
C	29 PLANE
IF(IE.EQ.1 .AND.NL.GT.MTOT) CALL ERROR(NL-MTOT,4)	30 PLANE
CALL PLANE1(A(N1),A(N2),A(N3),A(N4),A(N5),A(N6),A(N101),	31 PLANE
XA(N102),A(N103),A(N104),A(N105),A(N106) ,A(N107),A(N108),	32 PLANE
XA(N109),A(N110),A(N111))	33 PLANE
C	34 PLANE
RETURN	35 PLANE
END	36 PLANE
C	1 PLANE1
SUBROUTINE PLANE1(ID,X,D,IDIAG,B,A,E,POIS,WT,TH,GRAV,	2 PLANE1
XMAT,IEN,LM,IELNO,ISIDE,PRES)	3 PLANE1
C	4 PLANE1
C.... 4-NODE QUADRILATERAL ISOTROPIC PLANE STRESS/STRAIN ELEMENT	5 PLANE1
C	6 PLANE1
C.... OPTIONS	7 PLANE1

```

C      IOPT      = 0 PLANE STRAIN
C              = 1 PLANE STRESS
C
C      IL        = 0 TWO BY TWO QUADRATURE ON LAMBDA TERM
C              = 1 ONE-POINT QUADRATURE ON LAMBDA TERM
C
C      IM        = 0 TWO BY TWO QUADRATURE ON MU TERM
C              = 1 ONE-POINT QUADRATURE ON MU TERM
C
C      INC       = 0 INCOMPATIBLE MODES NEGLECTED
C              = 1 INCOMPATIBLE MODES ADDED
C
C      IHOURL    = 0 HOURGLASS STIFFNESS NEGLECTED
C              = 1 HOURGLASS STIFFNESS ADDED

```

```

C.... REPITION OF LAST TWO NODS PRODUCES CONSTANT STRESS/STRAIN TRIANGLE
C

```

```

      IMPLICIT REAL*8(A-H,C-Z)

```

```

C.... REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION
C

```

```

      LOGICAL QUAD
      COMMON/INFO/ICODE,MODE,NDOF,NSD,NUMNP
      COMMON/ELDATA/IE,IPREC,LENGTH,MAX,MTOT,NEG,NF,NL,NPAR(16),NUMEG
      COMMON/DEVICE/IEL,IIN,ICUT
      DIMENSION IC(NDOF,1),X(NSC,1),D(NDOF,1),IDIAG(1),B(1),A(1)
      DIMENSION E(1),POIS(1),WT(1),TH(1),GRAV(2),MAT(1),
      XIEN(4,1),LM(2,4,1),IELNO(1),ISIDE(1),PRES(1)
      DIMENSION S(12,12),P(12),DL(2,4),XL(2,4),SG(4),TG(4),
      XSH(3,6),EE(3),ST(3),PS(4),PE(4)
      DIMENSION S1(4,4),S2(2,2)
      DATA SG/-1.00,1.00,1.00,-1.00/
      DATA TG/-1.00,-1.00,1.00,1.00/
      EQUIVALENCE (NPAR(1),NTYPE),(NPAR(2),NUMEL),(NPAR(3),NUMAT),
      X(NPAR(4),IOPT),(NPAR(5),IL),(NPAR(6),IM),(NPAR(7),INC),
      X(NPAR(9),NUMPR),(NPAR(8),IHOURL)

```

```

C      NED=2
C      NEN=4
C      NEE=8
C      NEC=8
C      NEN1=NEN
C      NEE1=NEE
C      IF(11NC.EQ.0) GOTO 5
C      NEN1=6
C      NEE1=12
C      5 CONTINUE

```

changes

```

      GOTO (1,2,3

```

```

),IE

```

```

C      C.... READ AND PRINT ELEMENT DATA
C

```

```

1  WRITE(ICUT,2000) NTYPE,NUMEL,IOPT,IL,IM,INC,IHOURL,NUMPR
   WRITE(ICUT,2010) NUMAT
   WRITE(ICUT,2020)
   DO 10 N=1,NUMAT
     READ(IIN,1030) N,E(N),POIS(N),WT(N),TH(N)
     IF(TH(N).EQ.0.00) TH(N)=1.00
10  WRITE(ICUT,2030) N,E(N),POIS(N),WT(N),TH(N)
     READ(IIN,1010) GRAV
     WRITE(ICUT,2040) GRAV

```

```

C      CALL ELGEN(1CN,MAT,NEN) ~ correct material no. with element
      WRITE(ICUT,2050)
      NN=0
      DO 20 N=1,NUMEL
        NN=NN+1
        IF(NN.LE.50) GOTO 20

```

8 PLANE1
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76 PLANE1

```

NN=1
WRITE(IOUT,2050)
20 WRITE(IOUT,2060) N,MAT(N),(IEN(I,N),I=1,NEN)
C
CALL LOCAL(ID,IEN,LM,NDOF,NED,NEN,NUMEL)
CALL COLHT(IDIAG,LM,NED,NEN,NUMEL)
C
IF(NUMPR.EQ.0) RETURN
CALL PREAD(IELNO,ISIDE,PRES,1)
WRITE(IOUT,2070)
NN=0
DO 30 N=1,NUMPR
NN=NN+1
IF(NN.LE.50) GOTO 30
NN=1
WRITE(IOUT,2070)
30 WRITE(IOUT,2080) IELNO(N),ISIDE(N),PRES(N)
C
RETURN
C
C..... FORM STIFFNESS MATRIX AND GRAVITY LOAD VECTOR
C
2 DO 100 N=1,NUMEL
M=MAT(N)
QUAD=.TRUE.
IF(IEN(3,N).EQ.IEN(4,N)) QUAD=.FALSE.
CALL LUCRO(IEN(1,N),X,XL,NEC,NEN,NSD)
CALL CLEAR(S,144)
CALL CLEAR(P,12)
C
C..... SET MATERIAL CONSTANTS
C
AM=E(M)/(1.00+NUIS(M))
AL=AM*NUIS(M)/(1.00-2.00*NUIS(M))
AV=0.500*AM
IF(IUPT.EQ.1) AL=2.00*AL*AM/(AL+2.00*AM)
C
C..... LAMBDA TERM CONTRIBUTION TO STIFFNESS (GRAVITY LOAD COMPUTED IN
THIS SEGMENT)
C
NINT=4
W=1.00
G=1.00/DSQRT(3.00)
IF(IL.EQ.0) GOTO 110
NINT=1
W=4.00
G=0.00
110 DO 120 L=1,NINT
CALL SHAP20(SG(L)*G,TG(L)*G,XL,DET,SH,NEN,INC,IEN(1,N),QUAD)
C = DET*w*AL*TH(M)
GRAV1=DET*w*ht(M)*TH(M)
GRAV2=GRAV1*GRAV(2)
GRAV1=GRAV1*GRAV(1)
DO 120 J=1,NEN1
JB1=C*SH(1,J)
JB12=C*SH(2,J)
IF(J.LE.4) P(2*J-1)=P(2*J-1)+GRAV1*SH(3,J)
IF(J.LE.4) P(2*J)=P(2*J)+GRAV2*SH(3,J)
DO 120 I=1,J
S(2*I-1,2*J-1)=S(2*I-1,2*J-1)+SH(1,I)*JB1
S(2*I-1,2*J)=S(2*I-1,2*J)+SH(1,I)*JB12
S(2*I,2*J-1)=S(2*I,2*J-1)+SH(2,I)*JB1
S(2*I,2*J)=S(2*I,2*J)+SH(2,I)*JB12
120
C
C..... MU CONTRIBUTION TO STIFFNESS
C
NINT=4
W=1.00
G=1.00/DSQRT(3.00)

```

localize the ID array in LM
ie put global eqn no. of each element
node into LM

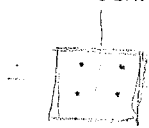
Circle for Δ

localize coord of element N. in X_b (NSD, NEN)

$$\mu = \frac{E}{1+\nu}$$

$$\lambda = \frac{E}{(1+\nu)} \cdot \frac{\nu}{1-2\nu}$$

plane stress $\lambda = \frac{2\lambda\mu}{2\lambda+2\mu}$



$$\bar{X}_a = \frac{mg}{V}$$

$$\int_{vol} N_a \bar{X} dSL = \int_{vol} N_a \bar{g} \sum dSL dt$$

$$N_a \left(\rho \sum_1 \right)$$

$$N_a \left(\rho \sum_2 \right)$$

$$S_{IJ} = S_{IJ} + B_a^T D B_b$$

S is equivalent to K_{ab} matrix

```

IF(IM.EQ.0) GOTO 130
NINT=1
W=4.00
G=0.00
130 DO 140 L=1,NINT
CALL SHAP20 (SG(L)*G,TG(L)*G,XL,DET,SH,NEN,INC,IEN(1,N),QUAD)
C =DET*W*AH*TH(M)
C2=2.00*C
DO 140 J=1,NEN1
DB11= C2*SH(1,J)
DB22= C2*SH(2,J)
DB31= C *SH(2,J)
DB32= C *SH(1,J)
DO 140 I=1,J
S(2*I-1,2*J-1)=S(2*I-1,2*J-1)+SH(1,I)*DB11+SH(2,I)*DB31
S(2*I-1,2*J )=S(2*I-1,2*J )+SH(2,I)*DB32
S(2*I ,2*J-1)=S(2*I ,2*J-1)+SH(1,I)*DB31
140 S(2*I ,2*J )=S(2*I ,2*J )+SH(2,I)*DB32+SH(1,I)*DB32

```

```

C
IF(THICK.EQ.0.OR..NOT.QUAD) GO TO 149 no honeycomb / no triangle go to 149
CALL FOFMS1(XL,S1) ↓ honeycomb or triangle
CALL FOFMS2(XL,DET,S2)
IF(IUPT.EQ.0) C=4.000*AM*(AM+AL)/(AL+2.000*AM)
IF(IUPT.EQ.1) C=C(M)
C=C/(12.00*DET)
S2(1,1)=C*S2(1,1)
S2(1,2)=C*S2(1,2)
S2(2,2)=C*S2(2,2)
DO 145 J=1,NEN
DO 145 I=1,J
C=S1(I,J)*S2(1,2)
S(2*I-1,2*J-1)=S(2*I-1,2*J-1)+S1(I,J)*S2(1,1)
S(2*I-1,2*J )=S(2*I-1,2*J )+C
S(2*I ,2*J-1)=S(2*I ,2*J-1)+C
145 S(2*I ,2*J )=S(2*I ,2*J )+S1(I,J)*S2(2,2)

```

C.... SYMMETRIZE THE STIFFNESS

```

149 DO 150 I=2,NEN1
DO 150 J=1,I
150 S(I,J)=S(J,I)
IF(INC.EQ.1.AND.QUAD) CALL STACUN(S,P,12,4)
CALL PACK(S,8,12)
CALL DISBC(DL,D,IEN(1,N),LM(1,1,N),NDOF,NEE,NED,NEN)
CALL DISBC(DL,P,S,NEE)
100 CALL ADJUSTF(A,B,S,P,DIAG,LM(1,1,N),NEE)

```

C IF(NUMPR.EQ.C) RETURN

C.... PRESSURE LOADS

```

DO 160 K=1,NUMPR
N=IELNC(K) element no.
M=MAT(N) get material
CALL LCJORD(IEN(1,N),X,XL,NEC,NEN,NSD) localize coord of element N
I=ISIDE(K) IEN holds global node no of node I of element N
J=I+1 I = Side # of element
IF(J.EQ.5) J=1
EE(1)=(XL(2,I)-XL(2,J))*PRES(K)*TH(M)/2.00 1/2 force / side in x dir
EE(2)=(XL(1,J)-XL(1,I))*PRES(K)*TH(M)/2.00 1/2 force / side in y dir
DO 170 L=1,2
II=LM(L,I,N) get eqn no. of def L, node I of element N
JJ=LM(L,J,N) " " of def L, node J of element N
IF(II.NE.0) B(II)=B(II)+EE(L)
IF(JJ.NE.0) B(JJ)=B(JJ)+EE(L) assign 1/2 force to each node
170 CONTINUE
160 CONTINUE
RETURN

```

if node is not fixed



146 PLANE1
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```

C
C..... STRESS COMPUTATIONS (AT CENTROID)
C
3 NN=0
  WRITE(IOUT,2090) NEG
  DO 200 N=1,NUMEL
    M=MAT(N)
    CALL LCOGRO(IEN(1,N),X,XL,NEC,NEN,NSD)
    CALL LDISP(OL,D,IEN(1,N),NDGF,NED,NEN)
    QUAD=.TRUE.
    IF(IEN(3,N).EQ.IEN(4,N)) QUAD=.FALSE.
    CALL SHAP2G(0.00,0.00,XL,DET,SH,NEN,0,IEN(1,N),QUAD)
    EE(1)=0.00
    EE(2)=0.00
    EE(3)=0.00
    XX=0.00
    YY=0.00
    AM=E(M)/(1.00+PCIS(M))
    AL=AM*PCIS(M)/(1.00-2.00*PCIS(M))
    AM=0.500*AM
    IF(IUPT.EQ.1) AL=2.00*AL*AM/(AL+2.00*AM)
    AL2=AL+2.00*AM
    DO 210 I=1,NEN
      EE(1)=EE(1)+SH(1,I)*DL(1,I)
      EE(2)=EE(2)+SH(2,I)*DL(1,I)+SH(1,I)*DL(2,I)
      EE(3)=EE(3)+SH(2,I)*DL(2,I)
      XX=XX+SH(3,I)*XL(1,I)
      YY=YY+SH(3,I)*XL(2,I)
210 ST(1)=AL2*EE(1)+AL*EE(3)
      ST(2)=AM*EE(2)
      ST(3)=AL*EE(1)+AL2*EE(3)
      PE(2)=0.500*EE(2)
      CALL PRINC(NSD,ST,PS)
      CALL PRINC(NSD,EE,PE)
      EE(2)=2.00*EE(2)
      NN=NN+1
      IF(NN.LE.5) GOTO 200
      NN=1
      WRITE(IOUT,2090) NEG
200 WRITE(IOUT,2100) N,M,XX,YY,ST,EE,PS,PE
C
      RETURN
1000 FORMAT(15,5X,4F10.0)
1010 FORMAT(2F10.0)
2000 FORMAT(1H0,
  X' P L A N E   S T R E S S / S T R A I N   E L E M E N T S'//5X,
  X' ELEMENT TYPE NUMBER . . . . . (NTYPE) =',15//5X,
  X' NUMBER OF ELEMENTS . . . . . (NUMEL) =',15//5X,
  X' ANALYSIS OPTION . . . . . (IOPT) =',15//5X,
  X'   EQ.0, PLANE STRAIN           ',15X,
  X'   EQ.1, PLANE STRESS           ',15X,
  X' LAMBDA TERM INTEGRATION CODE . . . . . (IL) =',15//5X,
  X'   EQ.0, 2X2 GAUSSIAN QUADRATURE ',15X,
  X'   EQ.1, 1-POINT GAUSSIAN QUADRATURE ',15X,
  X' MUJ TERM INTEGRATION CODE . . . . . (IM) =',15//5X,
  X'   EQ.0, 2X2 GAUSSIAN QUADRATURE ',15X,
  X'   EQ.1, 1-POINT GAUSSIAN QUADRATURE ',15X,
  X' INCOMPATIBLE MODES CODE . . . . . (INC) =',15//5X,
  X'   EQ.0, INCOMPATIBLE MODES NEGLECTED ',15X,
  X'   EQ.1, INCOMPATIBLE MODES ADDED ',15X,
  X' HOURGLASS STIFFNESS CODE . . . . . (IHOUR) =',15//5X,
  X'   EQ.0, HOURGLASS STIFFNESS NEGLECTED ',15X,
  X'   EQ.1, HOURGLASS STIFFNESS ADDED ',15X,
  X' NUMBER OF ELEMENT PRESSURE LOAD CARDS . . . (NUMPR) =',15//)
2010 FORMAT(1H1,
  X' M A T E R I A L   S E T   D A T A           ',15X,
  X' NUMBER OF MATERIAL SETS . . . . . (NUMAT) =',15//)
2020 FORMAT(1H0,
  X 5X,'SET',11X,'YOUNGS',13X,'POISSONS',10X,'WEIGHT/UNIT',10X,
  X'THICKNESS',11X,/,
  X 4X,'NUMBER',9X,'MODULUS',14X,'RATIO',14X,'VOLUME',33X,/)
2030 FORMAT(3X,15,2X,4(1PE20.8))

```


C	SUBROUTINE FORMS2(X,DET,S2)	1 FORMS2
C		2 FORMS2
C....	PROGRAM TO COMPUTE ARRAY S2 FOR HOURGLASS STIFFNESS FORMULATION	3 FORMS2
C		4 FORMS2
	IMPLICIT REAL*8 (A-H,O-Z)	5 FORMS2
C		6 FORMS2
C....	REMOVE ABOVE CARD FOR SINGLE PRECISION OPERATION	7 FORMS2
C		8 FORMS2
	DIMENSION X(2,1),S2(2,1),XS(2,2),A(2,4)	9 FORMS2
C	DATA A/-0.500,-0.500,0.500,-0.500,0.500,0.500,-0.500,0.500/	10 FORMS2
		11 FORMS2
	DO 10 I=1,2	12 FORMS2
	DO 10 J=1,2	13 FORMS2
	XS(I,J)=0.000	14 FORMS2
	DO 10 K=1,4	15 FORMS2
10	XS(I,J)=XS(I,J)+X(I,K)*A(J,K)	16 FORMS2
C		17 FORMS2
	DET=XS(1,1)*XS(2,2)-XS(1,2)*XS(2,1)	18 FORMS2
C		19 FORMS2
	S2(1,1)=XS(2,1)**2+XS(2,2)**2	20 FORMS2
	S2(1,2)=-XS(1,1)*XS(2,1)-XS(1,2)*XS(2,2)	21 FORMS2
	S2(2,1)=S2(1,2)	22 FORMS2
	S2(2,2)=XS(1,1)**2+XS(1,2)**2	23 FORMS2
C		24 FORMS2
	RETURN	25 FORMS2
	END	26 FORMS2
		27 FORMS2

Solved Problems

TRUSS PROBLEM

CONTROL INFORMATION

NUMBER OF NODAL POINTS (NUMNP) = 27
 NUMBER OF SPACE DIMENSIONS (NSD) = 3
 NUMBER OF DEGREES OF FREEDOM PER NODE . . . (NDOF) = 3
 NUMBER OF ELEMENT GROUPS (NUMEG) = 1
 ANALYSIS CODE (ICODE) = 0
 EQ.0, STATIC ANALYSIS
 EQ.1, RANK CHECK
 SOLUTION MODE (MODE) = 1
 EQ.0, DATA CHECK
 EQ.1, EXECUTION

NODAL COORDINATE DATA

NODE NUMBER	X1	X2	X3
1	0.0	0.0	0.0
2	6.00000000D 01	0.0	0.0
3	1.20000000D 02	0.0	0.0
4	1.80000000D 02	0.0	0.0
5	2.40000000D 02	0.0	0.0
6	3.00000000D 02	0.0	0.0
7	3.60000000D 02	0.0	0.0
8	4.20000000D 02	0.0	0.0
9	4.80000000D 02	0.0	0.0
10	5.40000000D 02	0.0	0.0
11	6.00000000D 02	0.0	0.0
12	6.60000000D 02	0.0	0.0
13	7.20000000D 02	0.0	0.0
14	7.80000000D 02	0.0	0.0
15	3.00000000D 01	6.00000000D 01	0.0
16	9.00000000D 01	6.00000000D 01	0.0
17	1.50000000D 02	6.00000000D 01	0.0
18	2.10000000D 02	6.00000000D 01	0.0
19	2.70000000D 02	6.00000000D 01	0.0
20	3.30000000D 02	6.00000000D 01	0.0
21	3.90000000D 02	6.00000000D 01	0.0
22	4.50000000D 02	6.00000000D 01	0.0
23	5.10000000D 02	6.00000000D 01	0.0
24	5.70000000D 02	6.00000000D 01	0.0
25	6.30000000D 02	6.00000000D 01	0.0
26	6.90000000D 02	6.00000000D 01	0.0
27	7.50000000D 02	6.00000000D 01	0.0

NODAL BOUNDARY CCNDITION CODES

NODE NUMBER	DOF1	DOF2	DOF3
1	1	1	1
2	0	0	1
3	0	0	1
4	0	0	1
5	0	0	1
6	0	0	1
7	0	0	1
8	0	0	1
9	0	0	1
10	0	0	1
11	0	0	1
12	0	0	1
13	0	0	1
14	1	1	1
15	0	0	1
16	0	0	1
17	0	0	1
18	0	0	1
19	0	0	1
20	0	0	1
21	0	0	1
22	0	0	1
23	0	0	1
24	0	0	1
25	0	0	1
26	0	0	1
27	0	0	1

NODAL FORCES AND DISPLACEMENTS

NODE NUMBER	DOF1	DOF2	DOF3
21	0.0	-1.00000000D 03	0.0

ELEMENT GROUP DATA

ELEMENT GROUP NUMBER . . . (NEG) = 1

TRUSS ELEMENTS

ELEMENT TYPE NUMBER (NTYPE) = 1

NUMBER OF ELEMENTS (NUMEL) = 51

M A T E R I A L S E T D A T A

NUMBER OF MATERIAL SETS (NUMAT) = 1

SET NUMBER	YOUNGS MODULUS	CROSS-SECTION AREA	WEIGHT/UNIT LENGTH
1	3.00000000D 07	1.00000000D 01	0.0

G R A V I T Y L O A D M U L T I P L I E R S

X1-DIRECTION	=	0.0
X2-DIRECTION	=	0.0
X3-DIRECTION	=	0.0

E L E M E N T D A T A

ELEMENT NUMBER	MATERIAL NUMBER	NODE 1	NODE 2
1	1	1	2
2	1	2	3
3	1	3	4
4	1	4	5
5	1	5	6
6	1	6	7
7	1	7	8
8	1	8	9
9	1	9	10
10	1	10	11
11	1	11	12
12	1	12	13
13	1	13	14
14	1	15	16
15	1	16	17
16	1	17	18
17	1	18	19
18	1	19	20
19	1	20	21
20	1	21	22
21	1	22	23
22	1	23	24
23	1	24	25
24	1	25	26
25	1	26	27
26	1	1	15
27	1	2	16
28	1	3	17
29	1	4	18
30	1	5	19

31	1	6	20
32	1	7	21
33	1	8	22
34	1	9	23
35	1	10	24
36	1	11	25
37	1	12	26
38	1	13	27
39	1	2	15
40	1	3	16
41	1	4	17
42	1	5	18
43	1	6	19
44	1	7	20
45	1	8	21
46	1	9	22
47	1	10	23
48	1	11	24
49	1	12	25
50	1	13	26

E L E M E N T D A T A

ELEMENT NUMBER	MATERIAL NUMBER	NODE 1	NODE 2
51	1	14	27

TRUSS PROBLEM

E Q U A T I O N S Y S T E M D A T A

NUMBER OF EQUATIONS (NEQ) = 50

NUMBER OF TERMS IN STIFFNESS (NA) = 791

MEAN HALF BANDWIDTH (MB) = 15

TOTAL LENGTH OF BLANK COMMON REQUIRED (NB) = 2609

D I S P L A C E M E N T S

NODE NUMBER	DOF1	DOF2	DOF3
1	0.0	0.0	0.0
2	-2.76923077D-04	-2.51797004D-03	0.0
3	-4.53846154D-04	-5.16286315D-03	0.0
4	-5.30769231D-04	-7.73467934D-03	0.0
5	-5.07692308D-04	-1.00334186D-02	0.0

6	-3.84615385D-04	-1.18590809D-02	0.0
7	-1.61538462D-04	-1.30116664D-02	0.0
8	1.61538462D-04	-1.30116664D-02	0.0
9	3.84615385D-04	-1.18590809D-02	0.0
10	5.07692308D-04	-1.00334186D-02	0.0
11	5.30769231D-04	-7.73467934D-03	0.0
12	4.53846154D-04	-5.16286315D-03	0.0
13	2.76923077D-04	-2.51797004D-03	0.0
14	0.0	0.0	0.0
15	2.10000000D-03	-1.18975425D-03	0.0
16	2.00000000D-03	-3.79618582D-03	0.0
17	1.80000000D-03	-6.42954047D-03	0.0
18	1.50000000D-03	-8.88981820D-03	0.0
19	1.10000000D-03	-1.09770190D-02	0.0
20	6.00000000D-04	-1.24911429D-02	0.0
21	-2.00577402D-18	-1.32321898D-02	0.0
22	-6.00000000D-04	-1.24911429D-02	0.0
23	-1.10000000D-03	-1.09770190D-02	0.0
24	-1.50000000D-03	-8.88981820D-03	0.0
25	-1.80000000D-03	-6.42954047D-03	0.0
26	-2.00000000D-03	-3.79618582D-03	0.0
27	-2.10000000D-03	-1.18975425D-03	0.0

TRUSS ELEMENT STRESSES

ELEMENT GROUP NUMBER (NEG) = 1

ELEMENT NUMBER	STRESS	FORCE
1	-1.38435121D 02	-1.38435121D 03
2	-8.84323906D 01	-8.84323906D 02
3	-3.84339791D 01	-3.84339791D 02
4	1.15604790D 01	1.15604790D 02
5	6.15523492D 01	6.15523492D 02
6	1.11543997D 02	1.11543997D 03
7	1.61538462D 02	1.61538462D 03
8	1.11543997D 02	1.11543997D 03
9	6.15523492D 01	6.15523492D 02
10	1.15604790D 01	1.15604790D 02
11	-3.84339791D 01	-3.84339791D 02
12	-8.84323906D 01	-8.84323906D 02
13	-1.38435121D 02	-1.38435121D 03
14	-4.99716938D 01	-4.99716938D 02
15	-9.99711059D 01	-9.99711059D 02
16	-1.49974779D 02	-1.49974779D 03
17	-1.99981848D 02	-1.99981848D 03
18	-2.49990448D 02	-2.49990448D 03
19	-2.99997712D 02	-2.99997712D 03
20	-2.99997712D 02	-2.99997712D 03
21	-2.49990448D 02	-2.49990448D 03
22	-1.99981848D 02	-1.99981848D 03
23	-1.49974779D 02	-1.49974779D 03

24	-9.99711059D 01	-9.99711059D 02
25	-4.99716938D 01	-4.99716938D 02
26	-5.58823331D 01	-5.58823331D 02
27	-5.58790241D 01	-5.58790241D 02
28	-5.58794705D 01	-5.58794705D 02
29	-5.58835569D 01	-5.58835569D 02
30	-5.58901680D 01	-5.58901680D 02
31	-5.58971883D 01	-5.58971883D 02
32	-5.59015024D 01	-5.59015024D 02
33	5.59044836D 01	5.59044836D 02
34	5.59115877D 01	5.59115877D 02
35	5.59194428D 01	5.59194428D 02
36	5.59254335D 01	5.59254335D 02
37	5.59279445D 01	5.59279445D 02
38	5.59263604D 01	5.59263604D 02
39	5.59263604D 01	5.59263604D 02
40	5.59279445D 01	5.59279445D 02
41	5.59254335D 01	5.59254335D 02
42	5.59194428D 01	5.59194428D 02
43	5.59115877D 01	5.59115877D 02
44	5.59044836D 01	5.59044836D 02
45	-5.59015024D 01	-5.59015024D 02
46	-5.58971883D 01	-5.58971883D 02
47	-5.58901680D 01	-5.58901680D 02
48	-5.58835569D 01	-5.58835569D 02
49	-5.58794705D 01	-5.58794705D 02
50	-5.58790241D 01	-5.58790241D 02

TRUSS ELEMENT STRESSES

ELEMENT GROUP NUMBER (NEG) = 1

ELEMENT NUMBER	STRESS	FORCE
51	-5.58823331D 01	-5.58823331D 02

TRUSS PROBLEM

SOLUTION TIME LOG IN SECONDS

INPUT PHASE	=	0.43
FORMATION OF STIFFNESS AND LOAD	=	0.28
FACTOR STIFFNESS	=	0.23
DISPLACEMENT AND STRESS ANALYSIS	=	0.47
TOTAL SOLUTION TIME	=	1.42

PLANE STRAIN P.R.=0.300 1X1 MU TERM , 1X1 LAMBDA TERM , HOURGLASS ADDED

C O N T R O L I N F O R M A T I O N

NUMBER OF NODAL POINTS (NUMNP) = 45
 NUMBER OF SPACE DIMENSIONS (NSD) = 2
 NUMBER OF DEGREES OF FREEDOM PER NODE (NDOF) = 2
 NUMBER OF ELEMENT GROUPS (NUMEG) = 1
 ANALYSIS CODE (ICODE) = 0
 EQ.0, STATIC ANALYSIS
 EQ.1, RANK CHECK
 SOLUTION MODE (MODE) = 1
 EQ.0, DATA CHECK
 EQ.1, EXECUTION

N O D A L C O O R D I N A T E D A T A

NODE NUMBER	X1	X2
1	0.0	0.0
2	2.00000000 00	0.0
3	4.00000000 00	0.0
4	6.00000000 00	0.0
5	8.00000000 00	0.0
6	1.00000000 01	0.0
7	1.20000000 01	0.0
8	1.40000000 01	0.0
9	1.60000000 01	0.0
10	0.0	5.00000000-01
11	2.00000000 00	5.00000000-01
12	4.00000000 00	5.00000000-01
13	6.00000000 00	5.00000000-01
14	8.00000000 00	5.00000000-01
15	1.00000000 01	5.00000000-01
16	1.20000000 01	5.00000000-01
17	1.40000000 01	5.00000000-01
18	1.60000000 01	5.00000000-01
19	0.0	1.00000000 00
20	2.00000000 00	1.00000000 00
21	4.00000000 00	1.00000000 00
22	6.00000000 00	1.00000000 00
23	8.00000000 00	1.00000000 00
24	1.00000000 01	1.00000000 00
25	1.20000000 01	1.00000000 00
26	1.40000000 01	1.00000000 00
27	1.60000000 01	1.00000000 00
28	0.0	1.50000000 00
29	2.00000000 00	1.50000000 00
30	4.00000000 00	1.50000000 00
31	6.00000000 00	1.50000000 00
32	8.00000000 00	1.50000000 00
33	1.00000000 01	1.50000000 00
34	1.20000000 01	1.50000000 00
35	1.40000000 01	1.50000000 00
36	1.60000000 01	1.50000000 00
37	0.0	2.00000000 00
38	2.00000000 00	2.00000000 00
39	4.00000000 00	2.00000000 00
40	6.00000000 00	2.00000000 00
41	8.00000000 00	2.00000000 00
42	1.00000000 01	2.00000000 00
43	1.20000000 01	2.00000000 00
44	1.40000000 01	2.00000000 00
45	1.60000000 01	2.00000000 00

NODAL BOUNDARY CCNDITION CODES

NODE NUMBER	DOF1	DOF2
1	1	1
2	1	0
3	1	0
4	1	0
5	1	0
6	1	0
7	1	0
8	1	0
9	1	0
37	1	0

NODAL FORCES AND DISPLACEMENTS

NODE NUMBER	DOF1	DOF2
9	0.0	-9.280000000-02
10	-7.500000000-01	1.738300000-01
18	0.0	-1.736300000-01
19	-1.500000000 00	1.386700000-01
27	0.0	-1.386700000-01
28	-2.250000000 00	8.008000000-02
36	0.0	-8.008000000-02
37	0.0	1.465000000-02
45	0.0	-1.465000000-02

ELEMENT GROUP DATA

ELEMENT GROUP NUMBER . . . (NEG) = 1

PLANE STRESS / STRAIN ELEMENTS

ELEMENT TYPE NUMBER (NTYPE) = 3
NUMBER OF ELEMENTS (NUMEL) = 32
ANALYSIS OPTION (IOPT) = 0

EQ.0, PLANE STRAIN
EQ.1, PLANE STRESS

LAMBDA TERM INTEGRATION CODE (IL) = 1
EQ.0, 2X2 GAUSSIAN QUADRATURE
EQ.1, 1-POINT GAUSSIAN QUADRATURE

MU TERM INTEGRATION CODE (IM) = 1
EQ.0, 2X2 GAUSSIAN QUADRATURE
EQ.1, 1-POINT GAUSSIAN QUADRATURE

INCOMPATIBLE MODES CODE (INC) = 0
EQ.0, INCOMPATIBLE MODES NEGLECTED
EQ.1, INCOMPATIBLE MODES ADDED

HOURGLASS STIFFNESS CODE (IHOUR) = 1
EQ.0, HOURGLASS STIFFNESS NEGLECTED
EQ.1, HOURGLASS STIFFNESS ADDED

NUMBER OF ELEMENT PRESSURE LOAD CARDS . . . (NUMPR) = 0

M A T E R I A L S E T D A T A

NUMBER OF MATERIAL SETS (NUMAT) = 1

SET NUMBER	YOUNGS MODULUS	POISSONS RATIO	WEIGHT/UNIT VOLUME	THICKNESS
1	1.000000000 00	3.000000000-01	0.0	1.000000000 00

G R A V I T Y L C A D M U L T I P L I E R S

X1-DIRECTION = 0.0

X2-DIRECTION = 0.0

E L E M E N T D A T A

ELEMENT NUMBER	MATERIAL NUMBER	NODE 1	NODE 2	NODE 3	NODE 4
1	1	1	2	11	10
2	1	2	3	12	11
3	1	3	4	13	12
4	1	4	5	14	13
5	1	5	6	15	14
6	1	6	7	16	15
7	1	7	8	17	16
8	1	8	9	18	17
9	1	10	11	20	19
10	1	11	12	21	20
11	1	12	13	22	21
12	1	13	14	23	22
13	1	14	15	24	23
14	1	15	16	25	24
15	1	16	17	26	25
16	1	17	18	27	26
17	1	19	20	29	28
18	1	20	21	30	29
19	1	21	22	31	30
20	1	22	23	32	31
21	1	23	24	33	32
22	1	24	25	34	33
23	1	25	26	35	34
24	1	26	27	36	35
25	1	28	29	38	37
26	1	29	30	39	38
27	1	30	31	40	39
28	1	31	32	41	40
29	1	32	33	42	41
30	1	33	34	43	42
31	1	34	35	44	43
32	1	35	36	45	44

PLANE STRAIN P.R.=0.300 1X1 MU TERM , 1X1 LAMBDA TERM , HOURGLASS ADDED

E Q U A T I O N S Y S T E M D A T A

NUMBER OF EQUATIONS (NEQ) = 79

NUMBER OF TERMS IN STIFFNESS (NA) = 1401

MEAN HALF BANDWIDTH (MB) = 17

TOTAL LENGTH OF BLANK COMMON REQUIRED (NB) = 3918

DISPLACEMENTS

NODE NUMBER	DOF1	DOF2
1	0.0	0.0
2	0.0	-6.474653820 00
3	0.0	-2.252313030 01
4	0.0	-4.675671700 01
5	0.0	-7.781646640 01
6	0.0	-1.143365390 02
7	0.0	-1.549511050 02
8	0.0	-1.983012320 02
9	0.0	-2.429982530 02
10	-1.369940190-01	-1.423141900-01
11	2.419415360 00	-6.603561650 00
12	4.641170700 00	-2.263235260 01
13	6.517505300 00	-4.684804140 01
14	8.052915360 00	-7.788959550 01
15	9.247664650 00	-1.143914730 02
16	1.010151150 01	-1.549576450 02
17	1.060989380 01	-1.983135700 02
18	1.078355480 01	-2.430022060 02
19	-2.203007990-01	-5.685737110-01
20	4.891882760 00	-6.990453230 00
21	9.334055470 00	-2.296181790 01
22	1.308656920 01	-4.712209060 01
23	1.615785270 01	-7.810898260 01
24	1.854688330 01	-1.145562010 02
25	2.025473610 01	-1.550974670 02
26	2.127285910 01	-1.983704190 02
27	2.162085190 01	-2.430143460 02
28	-1.947076100-01	-1.279149380 00
29	7.470012620 00	-7.635726700 00
30	1.413018260 01	-2.350942850 01
31	1.975846000 01	-4.757915320 01
32	2.436740780 01	-7.847463370 01
33	2.794897140 01	-1.148304200 02
34	3.051111730 01	-1.552811740 02
35	3.204158010 01	-1.984563950 02
36	3.256708700 01	-2.430354850 02
37	0.0	-2.271419700 00
38	1.020510420 01	-8.539198150 00
39	1.907007640 01	-2.427535330 01
40	2.653879940 01	-4.821952140 01
41	3.272970020 01	-7.858654130 01
42	3.750943270 01	-1.152138590 02
43	4.091945080 01	-1.555390670 02
44	4.296688690 01	-1.985767190 02
45	4.368313060 01	-2.430670870 02

PLANE ELEMENT STRESSES

ELEMENT GROUP NUMBER

1

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
1	1	1.0000 00	2.5000-01	7.0390-01	-3.0600-01	3.6060-03	6.3910-01	-9.5160-01	-2.7120-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		8.6020-01	-1.5280-01	-2.3130 01	5.0650-01	8.4240-01	-4.7450-01	-2.3130 01	6.5840-01
2	1	3.0000 00	2.5000-01	6.1000-01	-3.6880-01	-7.8730-04	5.5540-01	-9.5890-01	-2.3860-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		7.8350-01	-1.7420-01	-2.5190 01	4.7880-01	7.8090-01	-4.6410-01	-2.5190 01	6.2250-01
3	1	5.0000 00	2.5000-01	5.1550-01	-3.6670-01	-1.4500-05	4.6910-01	-9.5350-01	-2.0100-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		7.0600-01	-1.9050-01	-2.7450 01	4.4820-01	7.1670-01	-4.4870-01	-2.7450 01	5.8270-01
4	1	7.0000 00	2.5000-01	4.2190-01	-3.6730-01	7.3690-05	3.8390-01	-9.5490-01	-1.6450-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		6.3450-01	-2.1260-01	-3.0070 01	4.2350-01	6.6030-01	-4.4090-01	-3.0070 01	5.5080-01
5	1	9.0000 00	2.5000-01	3.2820-01	-3.6730-01	-7.3720-05	2.9870-01	-9.5490-01	-1.2810-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		5.6630-01	-2.3820-01	-3.2960 01	4.0230-01	6.0830-01	-4.3770-01	-3.2960 01	5.2300-01

PLANE ELEMENT STRESSES

ELEMENT GROUP NUMBER

1

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
6	1	1.1000 01	2.5000-01	2.3460-01	-3.6670-01	1.2480-05	2.1350-01	-9.5350-01	-9.1470-02
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		5.0230-01	-2.6770-01	-3.6130 01	3.8500-01	5.6150-01	-4.3950-01	-3.6130 01	5.0050-01
7	1	1.3000 01	2.5000-01	1.4000-01	-3.6880-01	7.9690-04	1.2710-01	-9.5890-01	-5.3880-02
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		4.4570-01	-3.0490-01	-3.9660 01	3.7530-01	5.2450-01	-4.5130-01	-3.9660 01	4.8790-01
8	1	1.5000 01	2.5000-01	4.6160-02	-3.6600-01	-3.6140-03	4.3410-02	-9.5170-01	-2.1290-02
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		3.8820-01	-3.4560-01	-4.3060 01	3.6690-01	4.8800-01	-4.6590-01	-4.3060 01	4.7700-01

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
9	1	1.0000 00	7.5000-01	2.1110 00	-3.1980-01	1.0880-02	1.9170 00	-8.3150-01	-8.1360-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		2.1590 00	-3.6730-02	-8.4680 00	1.0980 00	1.9790 00	-8.7540-01	-8.4680 00	1.4270 00

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
10	1	3.0000 00	7.5000-01	1.8300 00	-3.2110-01	-2.5100-03	1.6660 00	-8.3480-01	-7.1590-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		1.8840 00	-5.7150-02	-9.6580 00	9.7070-01	1.7370 00	-7.8690-01	-9.6580 00	1.2620 00

PLANE ELEMENT STRESSES
ELEMENT GROUP NUMBER

1

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
11	1	5.0000 00	7.5000-01	1.5460 00	-3.2000-01	1.0270-04	1.4070 00	-8.3190-01	-6.0300-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		1.6100 00	-6.3490-02	-1.1240 01	8.3680-01	1.4900 00	-6.8570-01	-1.1240 01	1.0880 00

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
12	1	7.0000 00	7.5000-01	1.2660 00	-3.2040-01	1.8590-04	1.1520 00	-8.3310-01	-4.9340-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		1.3420 00	-7.6320-02	-1.3430 01	7.0920-01	1.2510 00	-5.9290-01	-1.3430 01	9.2200-01

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
13	1	9.0000 00	7.5000-01	9.8450-01	-3.2040-01	-1.8560-04	8.9590-01	-8.3310-01	-3.8410-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		1.0800 00	-9.5280-02	-1.6530 01	5.8740-01	1.0200 00	-5.0770-01	-1.6530 01	7.6360-01

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
14	1	1.1000 01	7.5000-01	7.0370-01	-3.2000-01	-1.0900-04	6.4040-01	-8.3190-01	-2.7450-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		0.2740-01	-1.2380-01	-2.1140 01	4.7560-01	8.0130-01	-4.3540-01	-2.1140 01	6.1830-01

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
15	1	1.3000 01	7.5000-01	4.2050-01	-3.2110-01	2.5360-03	3.8160-01	-8.3480-01	-1.6170-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		5.9460-01	-1.7160-01	-2.8470 01	3.8310-01	6.0800-01	-3.8800-01	-2.8470 01	4.9800-01

PLANE ELEMENT STRESSES
ELEMENT GROUP NUMBER

1

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
16	1	1.5000 01	7.5000-01	1.3860-01	-3.1990-01	-1.0900-02	1.3040-01	-8.3160-01	-6.3990-02
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		3.9230-01	-2.6460-01	-3.8420 01	3.2850-01	4.6020-01	-3.9360-01	-3.8420 01	4.2700-01

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
17	1	1.0000 00	1.2500 00	3.5180 00	-2.2720-01	1.8180-02	3.1940 00	-5.9080-01	-1.3550 00
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		3.5330 00	3.4380-03	-3.6990 00	1.7650 00	3.2130 00	-1.3750 00	-3.6990 00	2.2940 00
18	1	3.0000 00	1.2500 00	3.0490 00	-2.2580-01	-4.5060-03	2.7760 00	-5.8700-01	-1.1930 00
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		3.0650 00	-2.1110-02	-4.2070 00	1.5430 00	2.7970 00	-1.2140 00	-4.2070 00	2.0060 00
19	1	5.0000 00	1.2500 00	2.5770 00	-2.2670-01	5.5410-04	2.3450 00	-5.8950-01	-1.0050 00
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		2.5970 00	-1.9240-02	-4.9900 00	1.3080 00	2.3710 00	-1.0300 00	-4.9900 00	1.7010 00
20	1	7.0000 00	1.2500 00	2.1100 00	-2.2660-01	2.2640-04	1.9200 00	-5.8910-01	-8.2270-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		2.1340 00	-2.3840-02	-6.0610 00	1.0790 00	1.9510 00	-8.5400-01	-6.0610 00	1.4030 00

PLANE ELEMENT STRESSES
ELEMENT GROUP NUMBER

1

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
21	1	9.0000 00	1.2500 00	1.6400 00	-2.2660-01	-2.2090-04	1.4930 00	-5.8910-01	-6.3990-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		1.6710 00	-3.0940-02	-7.7210 00	8.5090-01	1.5330 00	-6.7980-01	-7.7210 00	1.1060 00
22	1	1.1000 01	1.2500 00	1.1730 00	-2.2670-01	-5.7500-04	1.0670 00	-5.8950-01	-4.5790-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		1.2150 00	-4.2870-02	-1.0570 01	6.2900-01	1.1220 00	-5.1290-01	-1.0570 01	8.1770-01
23	1	1.3000 01	1.2500 00	7.0210-01	-2.2570-01	4.5490-03	6.3710-01	-5.8690-01	-2.6970-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		7.6880-01	-6.2130-02	-1.6460 01	4.1550-01	7.2380-01	-3.5640-01	-1.6460 01	5.4010-01
24	1	1.5000 01	1.2500 00	2.3220-01	-2.2720-01	-1.8210-02	2.1840-01	-5.9080-01	-1.0710-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		3.6640-01	-1.5250-01	-3.0570 01	2.5940-01	3.9290-01	-2.8160-01	-3.0570 01	3.3730-01
25	1	1.0000 00	1.7500 00	4.9200 00	-8.7030-02	2.5420-02	4.4670 00	-2.2630-01	-1.8960 00
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		4.9220 00	2.3870-02	-1.0180 00	2.4490 00	4.4690 00	-1.8980 00	-1.0180 00	3.1840 00

PLANE ELEMENT STRESSES
ELEMENT GROUP NUMBER

1

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
26	1	3.0000 00	1.7500 00	4.2640 00	-8.4410-02	-6.9070-03	3.8830 00	-2.1950-01	-1.6690 00
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		4.2660 00	-8.5750-03	-1.1320 00	2.1370 00	3.8850 00	-1.6720 00	-1.1320 00	2.7790 00
27	1	5.0000 00	1.7500 00	3.6100 00	-8.6630-02	1.9610-03	3.2850 00	-2.2520-01	-1.4060 00
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		3.6130 00	-1.1720-04	-1.3740 00	1.8060 00	3.2870 00	-1.4090 00	-1.3740 00	2.3480 00
28	1	7.0000 00	1.7500 00	2.9530 00	-8.5770-02	-6.8160-04	2.6870 00	-2.2300-01	-1.1520 00
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		2.9550 00	-3.1700-03	-1.5620 00	1.4790 00	2.6910 00	-1.1560 00	-1.5620 00	1.9230 00
29	1	9.0000 00	1.7500 00	2.2970 00	-8.5780-02	6.8180-04	2.0900 00	-2.2300-01	-8.9530-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		2.3010 00	-2.5170-03	-2.1360 00	1.1520 00	2.0940 00	-8.9950-01	-2.1360 00	1.4970 00
30	1	1.1000 01	1.7500 00	1.6400 00	-8.6610-02	-1.9610-03	1.4930 00	-2.2520-01	-6.4130-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		1.6440 00	-6.5180-03	-3.0110 00	8.2550-01	1.4990 00	-6.4730-01	-3.0110 00	1.0730 00

PLANE ELEMENT STRESSES
ELEMENT GROUP NUMBER

1

ELEMENT NUMBER	MATERIAL NUMBER	COORDINATE X1	COORDINATE X2	STRESS 11	STRESS 12	STRESS 22	STRAIN 11	STRAIN 12	STRAIN 22
31	1	1.3000 01	1.7500 00	9.8590-01	-8.4450-02	6.9050-03	8.9450-01	-2.1960-01	-3.7820-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		9.9310-01	-3.2600-04	-4.8940 00	4.9670-01	9.0390-01	-3.8760-01	-4.8940 00	6.4570-01
32	1	1.5000 01	1.7500 00	3.3020-01	-8.6930-02	-2.5420-02	3.1040-01	-2.2600-01	-1.5190-01
		STRESS 1	STRESS 2	ANGLE	SHEAR	STRAIN 1	STRAIN 2	ANGLE	SHEAR
		3.5040-01	-4.5530-02	-1.3030 01	1.9790-01	3.3660-01	-1.7810-01	-1.3030 01	2.5730-01

PLANE STRAIN P.R.=0.300 1X1 MU TERM , 1X1 LAMBDA TERM , HOURGLASS ADDED

SOLUTION TIME LOG IN SECONDS

INPUT PHASE	=	0.55
FORMATION OF STIFFNESS AND LOAD	=	0.50
FACTOR STIFFNESS	=	0.40
DISPLACEMENT AND STRESS ANALYSIS	=	0.82
TOTAL SOLUTION TIME	=	2.26