

REVIEW OF UPW BASICS

Example: \hat{x} -polarized UPW traveling in $+\hat{z}$ direction

$$\bar{\mathbf{E}} = \hat{x}E_o \cos(\omega t - kz)$$

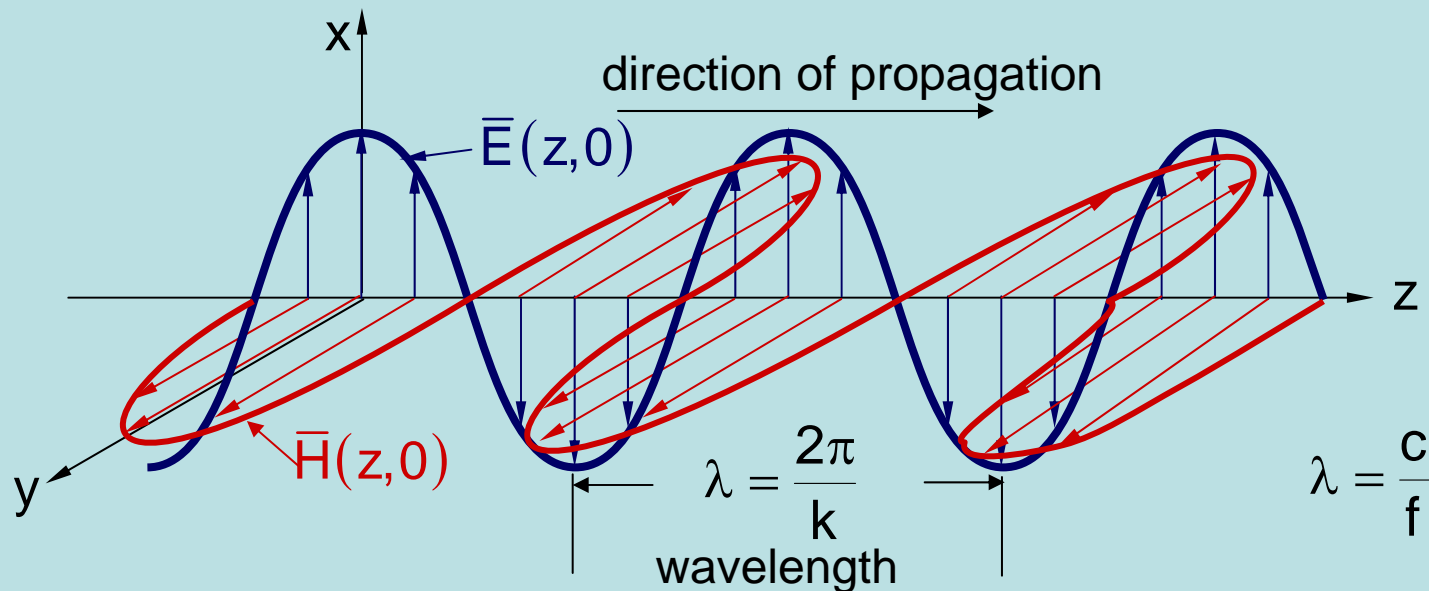
$$\bar{\mathbf{E}}(z) = \hat{x}E_o e^{-jkz}$$

$$\bar{\mathbf{H}} = \hat{y} \frac{E_o}{\eta_o} \cos(\omega t - kz)$$

$$\bar{\mathbf{H}}(z) = \hat{y} \frac{E_o}{\eta_o} e^{-jkz}$$

$\bar{\mathbf{E}} \times \bar{\mathbf{H}}$: \hat{z} direction of propagation

$$c = \frac{1}{\sqrt{\mu_o \epsilon_o}} = 3 \times 10^8 \text{ m/s} \quad \omega (\text{rads/s}) = 2\pi f \quad k (\text{rads/m}) = \frac{2\pi}{\lambda} = \frac{\omega}{c} \quad \eta_o = \sqrt{\frac{\mu_o}{\epsilon_o}}$$



HOW DO WAVES CONVEY POWER, ENERGY?

Recall: \bar{E} [V/m] \cdot \bar{J} [A/m²] = P_d [W/m³] But $\bar{E} \perp \bar{H}$

Manipulate Ampere's law to get $\bar{E} \cdot \bar{J}$

$$\bar{E} \cdot (\nabla \times \bar{H} = \bar{J} + \frac{\partial \bar{D}}{\partial t}) \quad \text{For symmetry, compute } \bar{H} \cdot (\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t})$$

$$\underbrace{\bar{H} \cdot (\nabla \times \bar{E}) - \bar{E} \cdot (\nabla \times \bar{H})}_{\text{Vector Identity}} = -\bar{H} \cdot \left(\frac{\partial \bar{B}}{\partial t}\right) - \bar{E} \cdot \left(\bar{J} + \frac{\partial \bar{D}}{\partial t}\right)$$

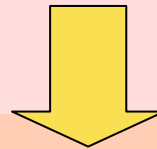
$$\nabla \cdot (\bar{E} \times \bar{H}) = -\bar{H} \cdot \frac{\partial \bar{B}}{\partial t} - \bar{E} \cdot \bar{J} - \bar{E} \cdot \frac{\partial \bar{D}}{\partial t} \quad (\text{W/m}^3)$$

This is Poynting's Theorem

What does it mean?

POYNTING THEOREM

Poynting's Theorem: $\nabla \cdot (\bar{\mathbf{E}} \times \bar{\mathbf{H}}) = -\bar{\mathbf{H}} \cdot \frac{\partial \bar{\mathbf{B}}}{\partial t} - \bar{\mathbf{E}} \cdot \frac{\partial \bar{\mathbf{D}}}{\partial t} - \bar{\mathbf{E}} \cdot \bar{\mathbf{J}}$



$$\bar{\mathbf{B}} = \mu \bar{\mathbf{H}} \quad \bar{\mathbf{D}} = \epsilon \bar{\mathbf{E}}$$

$$\underbrace{\nabla \cdot (\bar{\mathbf{E}} \times \bar{\mathbf{H}})}_{\substack{\text{Poynting} \\ \text{vector,} \\ \bar{\mathbf{S}} \text{ [W/m}^2\text{]}}} = - \frac{d}{dt} \underbrace{\left(\frac{1}{2} \mu |\bar{\mathbf{H}}|^2 \right)}_{\substack{\text{Stored magnetic} \\ \text{energy density,} \\ W_m}} - \frac{d}{dt} \underbrace{\left(\frac{1}{2} \epsilon |\bar{\mathbf{E}}|^2 \right)}_{\substack{\text{Stored electric} \\ \text{energy density,} \\ W_e}} - \underbrace{(\bar{\mathbf{E}} \cdot \bar{\mathbf{J}})}_{\substack{\text{Power density} \\ \text{dissipated/m}^3, \\ W_d}} \text{ (W/m}^3\text{)}$$

Energy is conserved! Note: $\frac{d}{dt} \left(\frac{1}{2} \mu |\bar{\mathbf{H}}|^2 \right) = \mu |\bar{\mathbf{H}}| \frac{d|\bar{\mathbf{H}}|}{dt} = \bar{\mathbf{H}} \cdot \frac{d\bar{\mathbf{B}}}{dt}$

Poynting Vector $\bar{\mathbf{S}} = \bar{\mathbf{E}} \times \bar{\mathbf{H}}$ (W/m²) $\left(\frac{\text{volts}}{\text{m}} \cdot \frac{\text{amps}}{\text{m}} = \frac{\text{watts}}{\text{m}^2} \right)$

INTEGRAL POYNTING THEOREM

Use: $\oint_S \bar{A} \cdot \hat{n} da = \int_V \nabla \cdot \bar{A} dv$

Gauss's Theorem (not Gauss's Law)

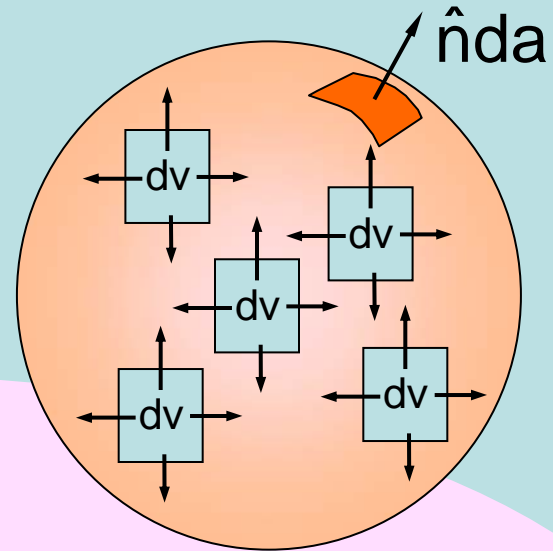
Therefore:

$$\begin{aligned} \oint_S (\bar{E} \times \bar{H}) \cdot \hat{n} da &= \int_V \nabla \cdot (\bar{E} \times \bar{H}) dv \\ &= \int_V \left[-\frac{d}{dt} \left(\frac{1}{2} \mu |\bar{H}|^2 \right) - \frac{d}{dt} \left(\frac{1}{2} \varepsilon |\bar{E}|^2 \right) - (\bar{E} \cdot \bar{J}) \right] dv \end{aligned}$$

$$\oint_S (\bar{E} \times \bar{H}) \cdot \hat{n} da = - \int_V \frac{d}{dt} \left(\frac{1}{2} \varepsilon |\bar{E}|^2 + \frac{1}{2} \mu |\bar{H}|^2 \right) dv - \int_V \bar{E} \cdot \bar{J} dv$$

Power emerging = released stored energy - dissipation [W]

The Poynting vector $\triangleq \bar{S} = \bar{E} \times \bar{H}$ gives both the magnitude of the power density (intensity) and the direction of its flow.



UNIFORM PLANE WAVE EXAMPLE

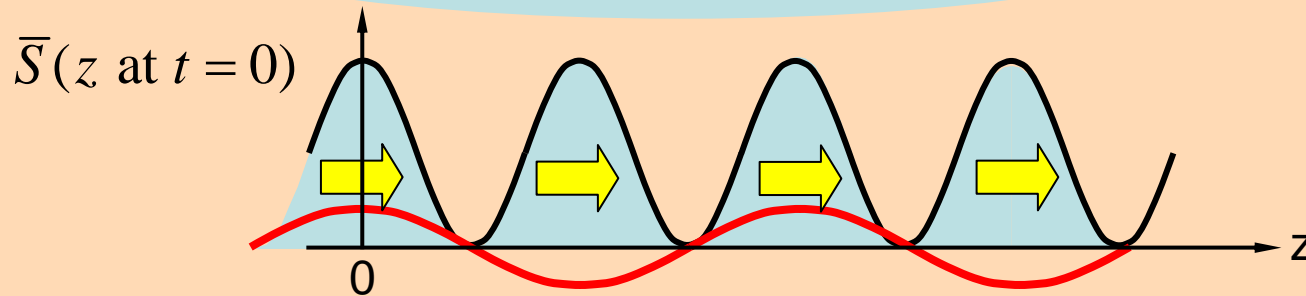
$$\bar{\mathbf{E}} = \hat{x}E_0 \cos(\omega t - kz)$$

$$W_e = \frac{1}{2} \epsilon_0 E_0^2 \cos^2(\omega t - kz)$$

$$\bar{\mathbf{H}} = \hat{y} \left(\frac{E_0}{\eta_0} \right) \cos(\omega t - kz)$$

$$W_m = \frac{1}{2} \frac{\mu_0}{\eta_0^2} E_0^2 \cos^2(\omega t - kz)$$

$$\bar{\mathbf{S}}(t) = \bar{\mathbf{E}} \times \bar{\mathbf{H}} = \hat{z} \left(\frac{E_0^2}{\eta_0} \right) \cos^2(\omega t - kz) \quad (\text{W/m}^2)$$



$$\bar{\mathbf{S}} = \hat{z} \frac{E_0^2}{\eta_0} \cos^2(\omega t - kz) \Rightarrow \langle \bar{\mathbf{S}} \rangle = \hat{z} \frac{1}{2} \frac{E_0^2}{\eta_0} = I(\theta, \phi, r) \quad [\text{W/m}^2]$$

The time average $\langle \bar{\mathbf{S}}(r, \theta, \phi) \rangle$ is “intensity” [W/m²]

COMPLEX NOTATION – POYNTING VECTOR

Defining a meaningful $\underline{\bar{S}}$ and relating it to \bar{S} is not obvious.
Let's work backwards to find the time average $\langle \bar{S} \rangle$ and then $\underline{\bar{S}}$

$$\begin{aligned}\bar{S}(t) &= \bar{E} \times \bar{H} = \text{Re} \left[\underline{\bar{E}} \cdot e^{j\omega t} \right] \times \text{Re} \left[\underline{\bar{H}} \cdot e^{j\omega t} \right] \\ &= [\bar{E}_r \cos(\omega t) - \bar{E}_i \sin(\omega t)] \times [\bar{H}_r \cos(\omega t) - \bar{H}_i \sin(\omega t)] \\ \Rightarrow \langle \bar{S}(t) \rangle &= \frac{1}{2} [(\bar{E}_r \times \bar{H}_r) + (\bar{E}_i \times \bar{H}_i)] \\ &= \frac{1}{2} \text{Re} \left(\underline{\bar{E}} \times \underline{\bar{H}}^* \right) \quad [= \frac{1}{2} \text{Re} \{ (\bar{E}_r + j\bar{E}_i) \times (\bar{H}_r - j\bar{H}_i) \}] \\ &\quad \underline{\bar{S}} \text{ (by definition)}\end{aligned}$$

Thus, we can define $\langle \bar{S} \rangle = \frac{1}{2} \text{Re} \left(\underline{\bar{E}} \times \underline{\bar{H}}^* \right)$ and $\underline{\bar{S}} = \underline{\bar{E}} \times \underline{\bar{H}}^*$

Recall: $\underline{\bar{E}} = \bar{E}_r + j\bar{E}_i$ $\underline{\bar{H}} = \bar{H}_r + j\bar{H}_i$ $e^{j\omega t} = \cos \omega t + j \sin \omega t$

UPW REFLECTED BY PERFECT CONDUCTOR

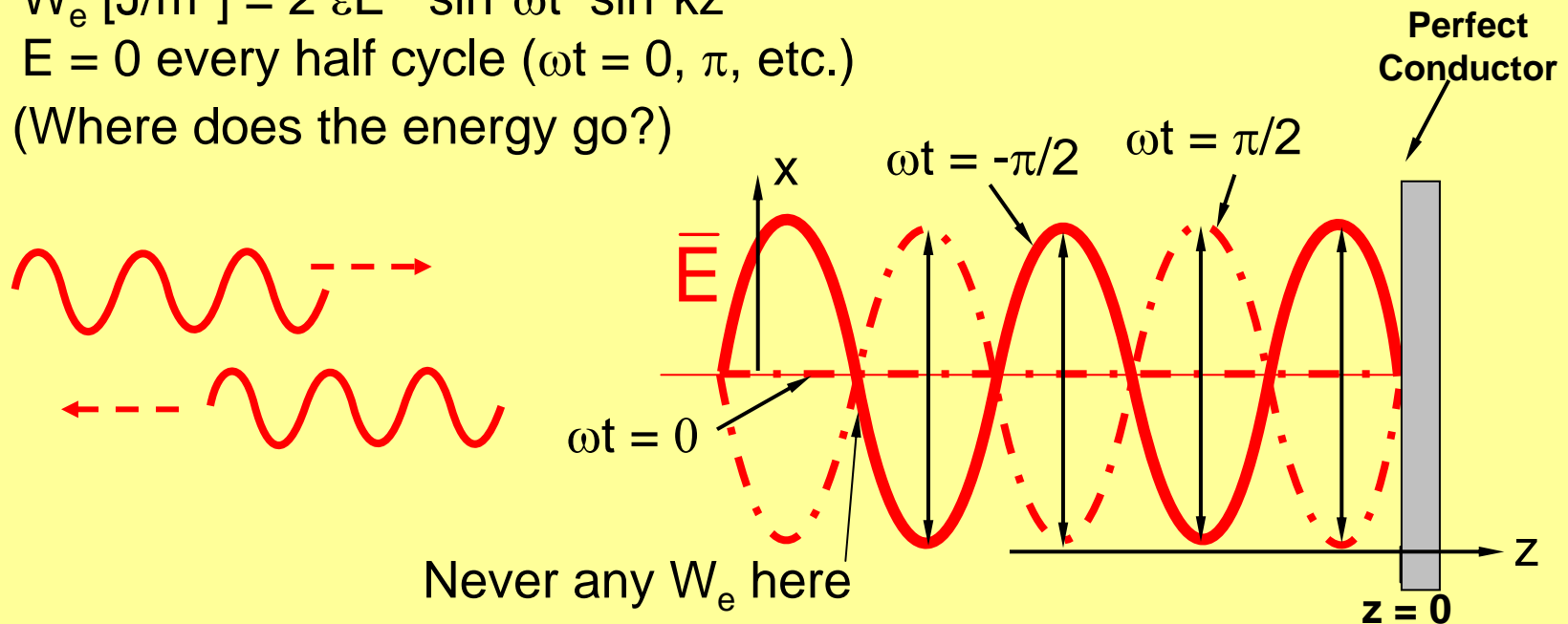
$$\left. \begin{aligned} \bar{E} &= \hat{x}E_+ \cos(\omega t - kz) + \hat{x}E_- \cos(\omega t + kz) \\ &= 0 \text{ at } z = 0 \text{ (perfect conductor)} \end{aligned} \right\} \begin{array}{l} \text{Forward plus} \\ \text{reflected wave} \end{array}$$

$$\Rightarrow E_- = -E_+ \text{ (Solving for unknown reflection)}$$

$$\Rightarrow \bar{E} = \hat{x}2E_+ \sin\omega t \cdot \sin kz \quad \text{Standing waves, oscillate without moving}$$

(recall: $\cos\alpha - \cos\beta = -2 \sin \frac{\alpha + \beta}{2} \cdot \sin \frac{\alpha - \beta}{2}$)

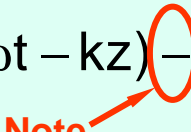
$W_e \text{ [J/m}^3\text{]} = 2 \epsilon E^2 \sin^2\omega t \sin^2kz$
 $E = 0$ every half cycle ($\omega t = 0, \pi$, etc.)
 (Where does the energy go?)



STANDING WAVE EXAMPLE - CONTINUED

(It's in the H field!) $\bar{\mathbf{E}} = \hat{x} [E_+ \cos(\omega t - kz) + E_- \cos(\omega t + kz)]$

$$\bar{\mathbf{H}} = \hat{y} \left[\frac{E_+}{\eta_0} \cos(\omega t - kz) - \frac{E_-}{\eta_0} \cos(\omega t + kz) \right]$$

Note 

$$E_- = -E_+ \Rightarrow \bar{\mathbf{H}} = \hat{y} \frac{2E_+}{\eta_0} \cos \omega t \cdot \cos kz$$

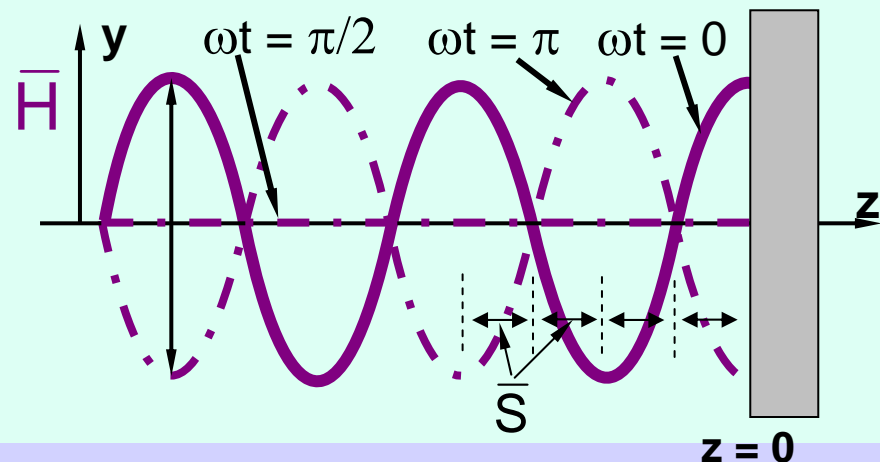
$$W_m [\text{J/m}^3] = \frac{1}{2} \mu_0 \left(\frac{2E_+}{\eta_0} \right)^2 \cos^2 \omega t \cos^2 kz = 2\epsilon E_+^2 \cos^2 \omega t \cos^2 kz$$

= 0 when $\omega t = \pi/2, 3\pi/2, \text{ etc.}$)

$$\bar{\mathbf{S}} = \bar{\mathbf{E}} \times \bar{\mathbf{H}} = \hat{z} 4 \frac{E_+^2}{\eta_0} \cos \omega t \sin \omega t \cdot \cos kz \sin kz$$

$$= \hat{z} \frac{E_+^2}{\eta_0} \cdot \sin 2kz \cdot \sin 2\omega t$$

$$\Rightarrow \langle \bar{\mathbf{S}} \rangle = 0 = \frac{1}{2} \text{Re}(\bar{\mathbf{E}} \times \bar{\mathbf{H}}^*)$$



MIT OpenCourseWare
<http://ocw.mit.edu>

6.013 Electromagnetics and Applications
Spring 2009

For information about citing these materials or our Terms of Use, visit: <http://ocw.mit.edu/terms>.